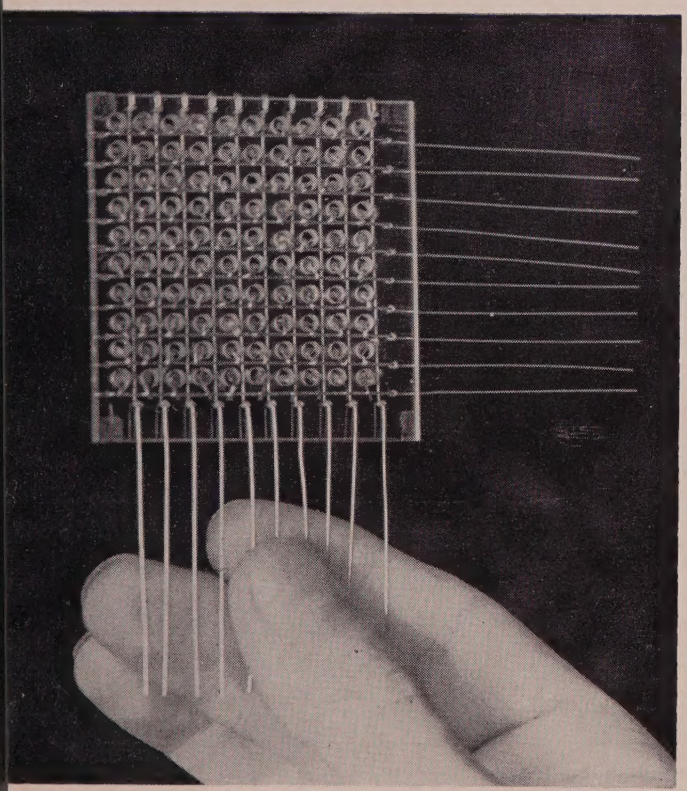
MAY · 1954

Proceedings of the I·R·E

A Journal of Communications and Electronic Engineering

FERROELECTRIC MEMORY MATRIX



Catholic University, Air Force Computer Laboratory

all ferroelectric storage condensers, individually selected for uniform city, coercive field, and sensitivity, are mounted in a cross-barlike structure oduce a simple, electrostatic memory matrix.

Volume 42 Number 5

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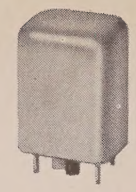
The IRE Standards on Volume Measurements of Electrical Speech appears in this issue.

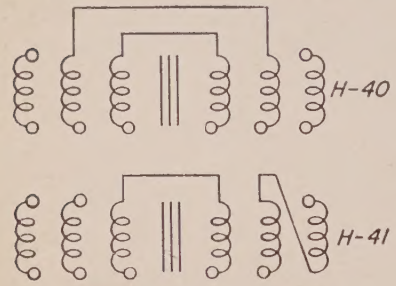
The Institute of Radio Engineers

HERMETICALLY SEALED PULSE TRANSFORMERS

RCOF CASE
1 25/64 X 61/64"



SM CASE
11/16 X 1/2"



Because of the wide variety of blocking oscillator, interstage, and modulator pulse applications, the bulk of UTC pulse transformers are designed to customer's specifications. Through versatile design, however, the stock hermetic MIL-T-27 pulse transformers listed below take care of most low level applications. Wide ranges of pulse duration, loading, and level are obtainable by variations in the manner of connecting the balanced coil structure windings as shown in the engineering sheet accompanying each unit.

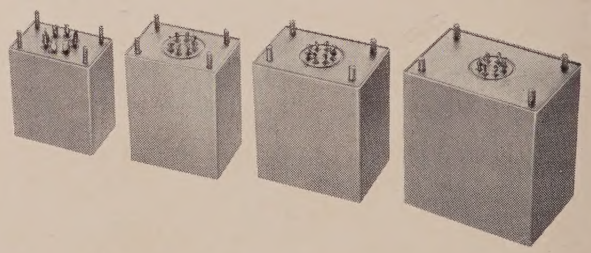
The H-40 and H-41 units employ identical windings suitable for different applications because of the manner in which the windings are brought out to the terminals. Pulse widths from .1 to 5 microseconds are realized with excellent fidelity. H-42 and H-43 are highly miniaturized units. They incorporate three equal windings capable of being inter-connected for wide versatility in blocking oscillator, interstage, and impedence matching service.

Type No.	Description*	Pulse Width Microsec.	Ins. Test Volts RMS	Case
H-40	Two 250 ohm windings . . . two 1000 ohm windings	.1 to 5	1000	RCOF
H-41	Two 250 ohm windings . . . two 1000 ohm windings	.1 to 5	1000	RCOF
H-42	Three 250 ohm windings	.15 to .75	500	SM**
H-43	Three 500 ohm windings	.5 to 2	500	SM**

* Impedances shown are nominal, subject to wide variation with application.
Mtg. screw is centered on large **side of case.

for SPECIAL APPLICATIONS

SERVO MOTOR MAGNETIC AMPLIFIERS



he MAT 1-4 Magnetic Amplifiers are exceptionally stable units designed for the control of 2 phase 400 cycle servo motors. They are ompact . . . hermetically sealed . . . magnetically shielded . . . meet IL-T-27 . . . high input impedance . . . high damping . . . high gain. The tput is sinusoidal, amplitude variable, and phase reversible. Control provided by a dual triode such as a 12AU7 operating with a plate ltage of 115 volts, 400 cycles, or higher. The signal to the triode ids can be polarity reversible DC or phase reversible 400 cycles. ower gain of the MAGNETIC STRUCTURE is approximately 40 . . . reponse time approximately 7.5 milliseconds . . . maximum null voltage V. RMS.

For AC signal control, the circuit of Figure 1 is employed. For C signal control, Figure 2 applies. Figure 3 shows the use of a ower transformer (MAT-5) which provides higher plate voltages and iminates the input transformer (MAT-6). The typical response curve Figure 4 applies to all units, the larger units feeding heavier loads.

TYPE NO.	MAT-1	MAT-2	MAT-3	MAT-4
230 Volt Supply				
Power output	4 W.	8 W.	11 W.	18 W.
RL, ohms	3300	1600	1200	720
CL, mfd.	2	.3	.5	.7
115 Volt Supply				
Power output	2 W.	4 W.	6 W.	9 W.
RL, ohms	6500	3300	2200	1450
CL, mfd.	.13	.2	.3	.45
Reson. Freq.	40 cyc.	35 cyc.	35 cyc.	20 cyc.
Log-Decr.	.18	.23	.03	.65
Cont. Wdg. Res.	6200 ohms	8450 ohms	4750 ohms	5650 ohms
Case				
Length, In.	1¼	1½	1¾	2⅛
Width, In.	1 5/16	2⅛	2½	3⅛
Height, In.	2 5/16	2¾	2 5/16	3⅜
Unit Weight, lbs.	.67	1.1	1.7	2.75

MAT-5 115V.-400 cyc. to 460 VCT; provides 230V. 48 MA DC or 460V. 24 MA DC. RC-37 Case . . . ⅜ x ⅜ x ⅝ . . . ⅛ mtg. holes 1⅛ x 1⅛ . . . 6 oz.

MAT-6 Input . . . 10,000 ohms pri . . . 1:15 C.T. ratio . . . phase shift under 1° . . . RCOF case.

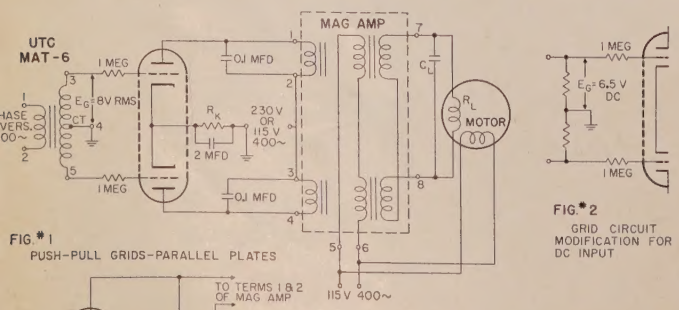


FIG.#1
PUSH-PULL GRIDS-PARALLEL PLATES

FIG.#2
GRID CIRCUIT MODIFICATION FOR DC INPUT

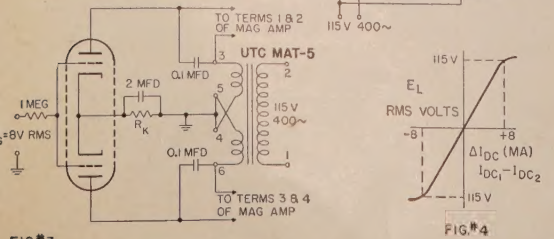


FIG.#3
PARALLEL GRIDS-PUSH-PULL PLATES

FIG.#4

United Transformer Co.
150 VARICK STREET • NEW YORK 13, N. Y.
EXPORT DIVISION: 13 EAST 40th STREET, NEW YORK 16, N. Y. CABLES: "ARLAB"

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The characteristics, limitations, applications, development, performance and reliability of component parts.

Mr. Floyd A. Paul, Chairman, Supervisor, Electronic Reliability Sec., Dept. 3230, Northrop Aircraft, Inc., Hawthorne, Calif.

Fee $2. Transaction* PGCP*1.

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Electron devices, including particularly electron tubes and solid state devices.

Dr. Leon S. Nergaard, Chairman, RCA Laboratories, Princeton, N.J.

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Gen. Tom C. Rives, U.S.A. (Ret'd.), Chairman, General Electric Company, Electronics Park, Syracuse, N.Y.

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Information theory and its application in radio circuitry and systems.

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Measurements and instrumentation utilizing electronic techniques.

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Mr. L. H. Montgomery, Jr., Chairman, 612 Craighead Street, Nashville, Tenn.

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Microwave theory, microwave circuitry and techniques, microwave measurements and the generation and amplification of microwaves.

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Communications problems in the field of land and mobile radio services, such as public safety, public utilities, railroads, commercial land transportation, etc.

Mr. W. A. Shipman, Chairman, Columbia Gas Sys. Ser. Corp., 120 East 41 Street, New York 17, N.Y.

Fee $2. *2, *3.

Ultrasonics Engineering

Ultrasonic measurements and communications, including underwater sound, ultrasonic delay lines, and various chemical and industrial ultrasonic devices.

Mr. A. L. Lane, Chairman, 706 Chillum Road, Apt. 101, Hyattsville, Md.

Fee $2.

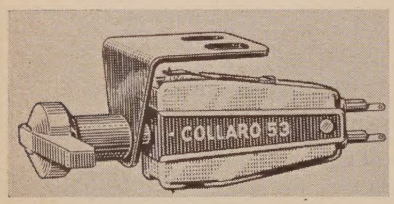
ENGINEERS **1 East 79th Street, New York 21, N. Y.**

NEWS and NEW PRODUCTS

May 1954

High Fidelity Crystal Cartridges

Rockbar Corp., 215 E. 37 St., New York 16, N. Y., has announced two new types of high fidelity crystal cartridges, manufactured in England by Collaro. It is claimed that these cartridges reproduce a smoother response than has previously been attained in crystal cartridges.

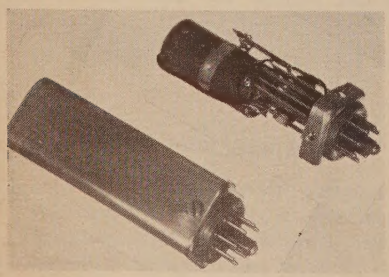


The Collaro cartridges are nonhygroscopic. They employ metal-sealed crystals impervious to normal humidity and temperature changes. They are designed to operate into a 1 megohm impedance and require no low frequency equalization.

The two types are identified as "O" and "P." These differ only in over-all frequency response and voltage output. Type "O" has a frequency response from 50 to 10,000 cycles 4-db, and a 0.5-volt output; Both cartridge types are of the "turnover" variety and are supplied with two interchangeable sapphire styli: 0.003 inch for standard, and 0.001 inch for microgroove. Standard $\frac{1}{2}$ inch mounting permits ready replacement of existing types.

Small Heater Rectifier

Vector Electronic Co., 3352 San Fernando Rd., Los Angeles 65, Calif., has a small rectifier-filter circuit which was developed especially for reducing hum by supplying dc heater current to the first stage of low level amplifiers. It is provided with an octal plug for convenient mounting and occupies a space above the chassis measuring only $1\frac{3}{8}\times1\frac{3}{8}\times4$ inches.

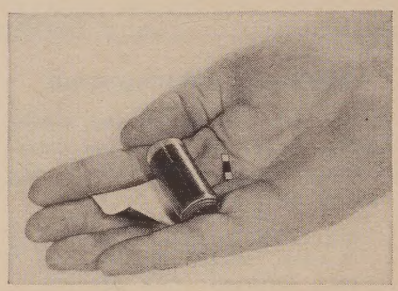


Since its rated input voltage is 6.3 to 7.5 volts ac, the device can be connected directly to the available ac filament supply. It provides up to 0.3 amperes dc, the dc output voltage being about 85 per cent of the ac input voltage at 0.3 amperes load, 95 per cent at 0.15 amperes load. Filtering is adequate to reduce the ac component in the load over 20 db.

Because of its small size the device can frequently be applied to existing amplifiers.

These manufacturers have invited PROCEEDINGS readers to write for literature and further technical information. Please mention your I.R.E. affiliation.

Tape Resistors for Printed Circuits



Sanders Associates, Inc., 137 Canal St., Nashua, N. H., now offers stable tape resistors for a wide range of printed circuit applications. They are available either as cured, ready-to-use resistors $\frac{1}{2}$ inch long, $\frac{1}{8}$ inch wide, and 1/100 inch thick, or as uncut, uncured tape rolls. Both types have a resistance range of 100 ohms to 10 megohms. These components conform to all JAN-R-11 specifications. These resistors are suitable for semi-automatic applications in which a single operation, requiring less than one second, fastens them permanently to the chassis and connects them into the circuit, without soldering, bending or punching holes in chassis. Characteristics are as follows: Power rating: $\frac{1}{4}$ watt at 150° C, resistance tolerance: ±10 per cent. Operating temperature range $-55°$ C to $+200°$ C, humidity: 95 per cent at 40° C for 250 hours, temperature coefficient: within requirements of JAN-R-11, shelf life: 18 months minimum, load life: 500 hours minimum $\frac{1}{4}$ watt at 150° C.

Remote Control Equipment For Broadcast Transmitters

The Hammarlund Mfg Co., 460 W. 34 St., New York 1, N. Y., has a brochure available describing its remote control equipment for broadcast transmitters. This type of operation by broadcasting companies was made possible recently through a new directive of the Federal Communications Commission.

The control equipment makes use of audio tones for complete control and metering of the remote transmitter including nine possible control functions and nine telemetering functions.

It makes use of recognized principles of telemetering and requires only a single circuit which may be wire, radio or microwave. No dc line is needed.

All the requirements of the Federal Communications Commission rules and regulations are fulfilled by the Hammarlund equipment.

Low-Frequency Noise Generator

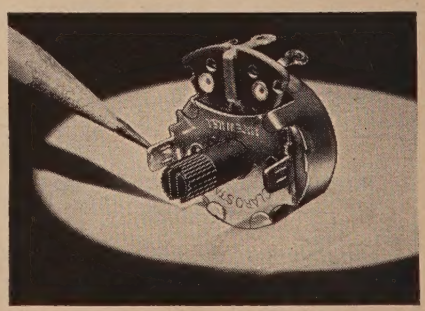
The model RUG-1-10 low-frequency noise generator manufactured by **Statistical Instrument Co.,** P.O. Box 552, Church St. Station, New York 8, N. Y., makes available for simulation studies and test purposes a random voltage source of controlled frequency spectrum and probability distribution.

Examples of the use of the RUG-1-10 are: a study of random airload effects in airframe design, noise problems in missile guidance, study of the statistical properties of ground electromagnetic reflection, low frequency phenomena including chemical and thermal processes and certain bioelectrical effects.

The basic noise source is a gas tube producing an approximate Gaussian output. This provides the signal source to the three controlled distribution channels generating the Gaussian, Rayleigh, and Uniform distributions. These distributions are accurate to 2 per cent. The frequency coverage is from 0–10 cps in three steps, 0–2, 0–5, 0–10 cps for all three distributions. Wider bandwidths are available on request. The RUG-1-10 will deliver approximately 5 volts RMS at the wide bandwidth and 1 volt at the narrow setting, and is continuously variable to 0.1-volt accuracy. Either cabinet or relay rack mounted is supplied. Size $7\frac{1}{2}\times19\times19$ inches.

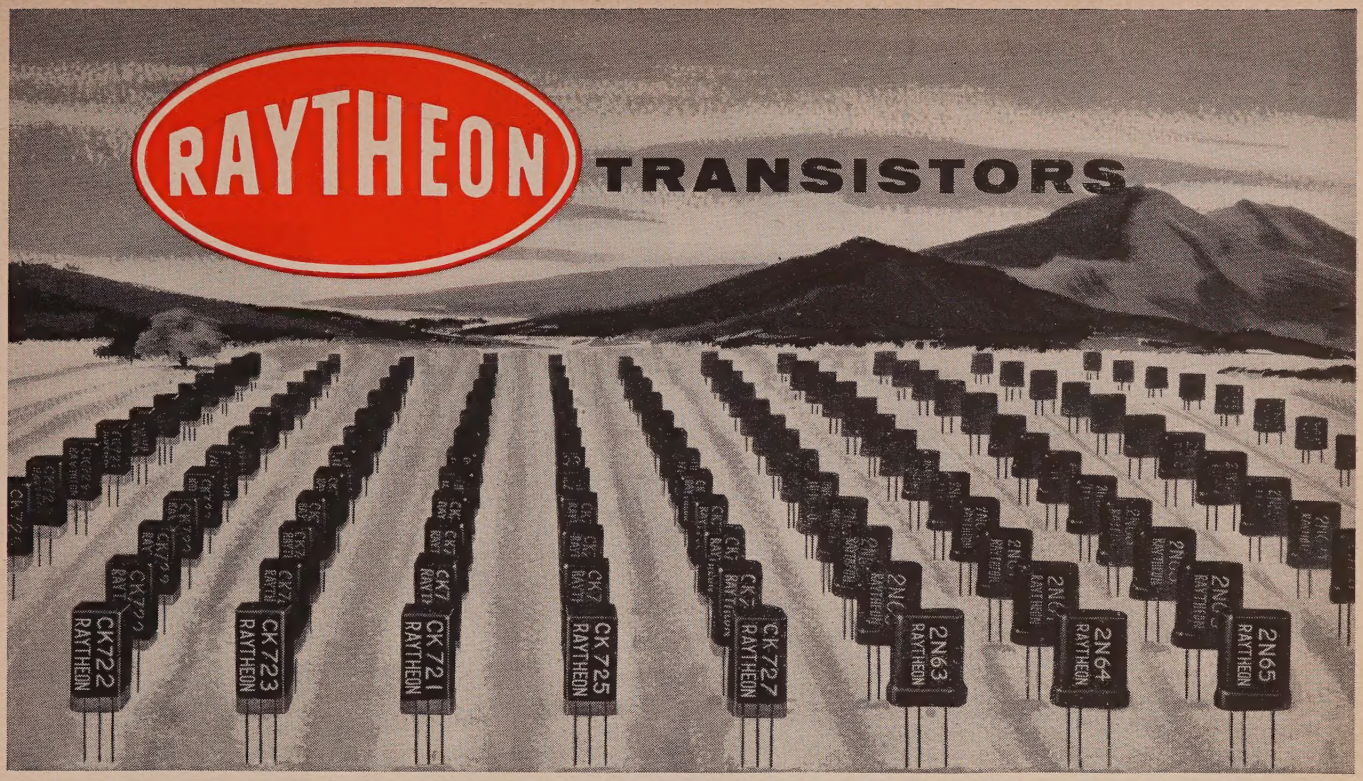
Low Cost Resistor

Reduced cost, states **Clarostat Mfg. Co., Inc.,** Dover, N. H., distinguishes its Series 47 twisted tab control. This twisted-tab mounting eliminates the usual bushing, lockwasher and nut, effecting the economy. The unit is mounted by inserting the tabs through slots in panel or chassis, and twisting them to secure the control in place.



Electrically, the control is the msae as the Clarostat Series 47, 15/16 inch diameter units. Available with or without switch. In resistance values from 500 ohms to 5 megohms; 0.5 watt rating; choice of tapers and taps; all types of metal or plastic shafts, including, if desired, a rear protruding slotted shaft. For more information write for drawing number 251914.

(Continued on page 24A)



Foremost in the field

RAYTHEON TRANSISTORS are FOREMOST IN THE FIELD with PROVEN RELIABILITY Over 1,000,000,000 OPERATING HOURS of actual field performance in commercial equipment with only a FRACTION OF 1% FIELD RETURNS proves their reliability to be superior to the reliability of vacuum tubes.

RAYTHEON TRANSISTORS are foremost in number of units in use in commercial equipment. Raytheon successfully made transistors in "experimental," "pilot" and now MASS PRODUCTION phases. The latest continuous, mass production, and inspection techniques are employed in the making of Raytheon Transistors. HUNDREDS OF THOUSANDS are IN ACTUAL COMMERCIAL USE — MANY TIMES MORE THAN ALL OTHER MAKES COMBINED. No other manufacturers can make these statements.

RAYTHEON GERMANIUM DIFFUSED JUNCTION PNP TRANSISTORS

RATINGS: — ABSOLUTE MAXIMUM VALUES:	CK722	CK723	CK721	CK725	CK727	2N63*	2N64*	2N65*
Collector Voltage (volts)	−22	−22	−22	−22	−6	−22	−22	−22
Collector Current (ma)	10	10	10	10	10	10	10	10
Collector Dissipation (30°C) (mw)	33	33	33	33	30	33	33	33
Emitter Current (ma)	10	10	10	10	10	10	10	10
Ambient Temperature (°C)	50	50	50	50	50	50	50	50
AVERAGE CHARACTERISTICS (27° C)								
Collector Voltage (volts)	−6	−6	−6	−6	−1.5	−6	−6	−6
Emitter Current (ma)	1	1	1	1	0.5	1	1	1
Collector Resistance (meg)	2.0	2.0	2.0	2.0	2.0	2.0	2.0	2.0
Emitter Resistance (ohms)	25	25	25	25	50	25	25	25
Base Resistance (ohms)	250	350	700	1500	500	350	700	1500
Base Current Amplification Factor	12	22	45	90	35	22	45	90
Cutoff Current (approx.) (ua)	6	6	6	6	6	6	6	6
Noise Factor (max) (db)**	30†	25†	22†	20†	12††	25†	22†	20†

*Hermetically sealed in metal package
**In a one cycle band width at 1000 cycles

†Measured at $V_c = -2.5$ volts in common emitter circuit
††Measured at $V_c = -1.5$ volts; $I_c = 0.5$ ma in common emitter circuit

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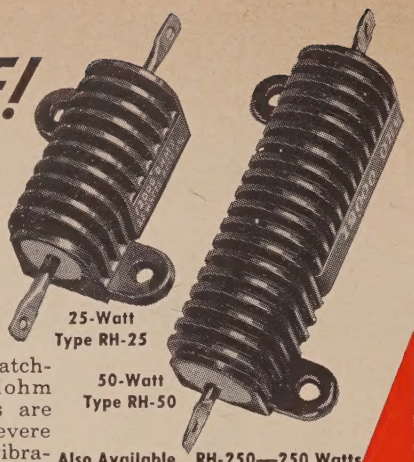
Miniature POWER Resistors



25-Watt Type RH-25

50-Watt Type RH-50

Carefully crafted for matchless performance, Dalohm miniature power resistors are made to survive the most severe environmental, shock, and vibration conditions.

Also Available RH-250—250 Watts
RS Types—2, 5, and 10 Watt

Dalohm RH type resistors are completely welded from terminal to terminal. They are silicone sealed in a die-cast, black anodized radiator finned housing and mount on subpanel for maximum heat dissipation.

- Temperature coefficient 0.00002/Deg. C.
- Ranges from 0.1 Ohms to 55,000 Ohms, depending on type.
- Tolerance 0.05%, 0.1%, 0.25%, 0.5%, 1%, 3%, 5%.
- Manufactured in accordance to applicable JAN and MIL specifications.

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SINCE 1901

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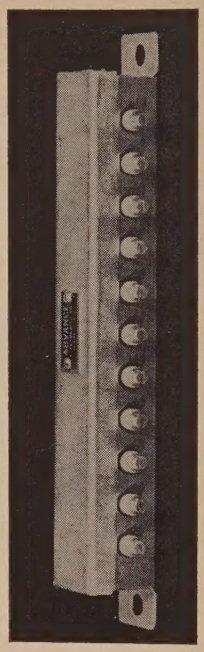
121 So. Columbus Avenue • Mount Vernon, N.Y.

News—New Products

These manufacturers have invited PROCEEDINGS readers to write for literature and further technical information. Please mention your I.R.E. affiliation.

(Continued from page 14A)

Tapped Delay Lines



Advanced Electronics Co., Inc., P.O. Box 394, Passaic, N. J., has six different types of tapped delay lines available: Type 6T1, Type 6T2, Type 6T3, Type 6T4, Type 6T5, and Type 6T6. Type 6T1 has a total time delay of 0.1 microsecond, 0.01 delay between taps, 75 ohms impedance, 32 mc cut off frequency and 0.01 microsecond maximum rise time. Type 6T2 has a total time delay of 0.25 μs, 0.025 μs delay between taps, 95 ohms impedance, 15 mc cutoff frequency, and 0.025 maximum rise time. Type 6T3 has a total time delay of 0.5 μs, 0.05 μs delay between taps, 190 ohms impedance, 6.2 mc cutoff frequency, and 0.05 μs rise time. Type 6T4 has a total time delay of 1 μs, 0.1 μs between taps, 300 ohms impedance, 3.2 mc cutoff frequency, and 0.1 μs rise time. Type 6T5 has between taps, 750 ohms impedance, 1.6 mc cutoff frequency, and 0.2 μs maximum rise time. Type 6T6 has a total time delay of 5 μs, 0.5 μs delay between taps, 1000 ohms impedance, 620 kc cutoff frequency, and 0.5 μs maximum rise time.

All the above instruments have ten taps each and their physical size is $6 \times \frac{7}{8} \times \frac{7}{8}$ inches.

Magnetic Deflection Yokes Catalog

Three new deflection yokes for military and oscilloscope applications are described in a new page of a catalogue now available from **Syntronic Instruments Inc.,** 100 Industrial Road, Addison, Ill. Type Y15-5 is illustrated. Complete technical information is given, including dimensional drawing, electrical and mechanical data, and complete tables of push-pull and single-ended deflection coil data.

(Continued on page 26A)

ENTIRELY NEW

POWERSTAT

Variable Transformers

TYPE 136 AND 236

provide new higher ratings . . . smaller size . . . easier installation and servicing . . . smoother operation and longer life . . . greater overload characteristics.

POWERSTAT types 136 and 236 are all new . . . new design . . . new performance . . . new ratings. Incorporated into each unit are all the features essential for the ultimate in variable transformers.

Types 136 and 236 are offered for manually-operated and motor driven duty in 120, 240, 480 volt ratings. Here are the reasons they are superior:

- **Higher Ratings:** Type 136 is rated 120 volts, 50/60 cycles input; 0-120/140 volts, 20.0 amperes output.
- **Smaller Size:** "Pancake" coil design provides compactness for bench or panel mounting.
- **Easier Installation:** Three sets of mounting holes suit all needs. Binding post type terminals provide all connection methods.
- **Easier Service:** Simply remove plate block for access to brush assembly. Brush easily removed and replaced.
- **New Rhodium Plated Commutator:** The one best answer to smooth performance and long life. The contact surface remains forever free of oxides. Corrosion is reduced. Uniform contact drop maintained and greater overload characteristics allowed.

There is a complete standard line of POWERSTAT variable transformers type 136 and 236 to suit individual requirements. Write for Bulletin P354.

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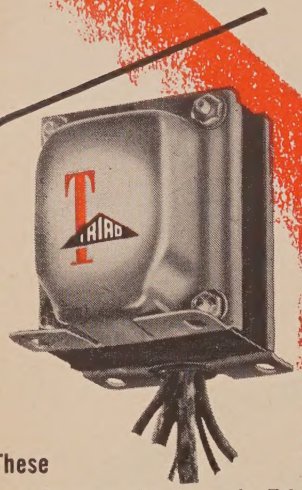
Power Supply
TRANSFORMERS



These transformers, developed by Triad especially for voltage regulated power supplies, cathode ray tube supplies, preamplifiers, and VT voltmeters, are among the more than

50 NEW Items in Triad's NEW 1954 Catalog

Write for your copy of Catalog TR-54 —it completely describes the finest line of transformers made.

TRIAD TRANSFORMER CORP.

4055 Redwood Ave. • Venice, Calif.

News—New Products

These manufacturers have invited PROCEEDINGS readers to write for literature and further technical information. Please mention your I.R.E. affiliation.

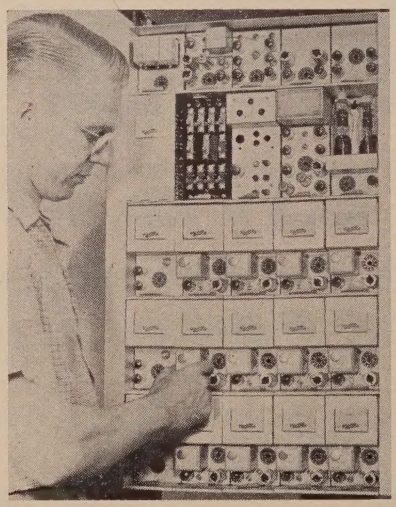
(Continued from page 24A)

Impulse Counter Brochure

Kellogg Switchboard and Supply Co., 79 W. Monroe St., Chicago 3, Ill., has a new brochure which describes their new selector switch utilizing relay-type contacts. Applications include the fields of industrial control, computer design, vhf radio, microwave, selective signaling, and telemetering. The switches are designed to handle electrical impulses that are to be registered, stored, and released at a rate of 20 per second (under proper conditions).

Miniaturized 12-Channel Carrier System

Deliveries of new type 45A carrier telephone systems are now being made by **Lenkurt Electric Co.**, San Carlos, Calif.



Providing up to 12 carrier-derived voice channels on an open wire line, Type 45A systems co-ordinate with systems such as Western Electric J and Lenkurt 42C. They can be installed on lines already equipped with carrier systems using frequencies up to 35 kc. Four staggered frequency allocations are available to permit installation of several systems on a single pole line.

All units for a complete 12-channel system plug into a prewired shelf which occupies 31½ inches on a standard 19-inch equipment rack. The shelf mounts through the rack, extending 5 inches to the front and 2 inches to the rear. Filters, hybrids, relays, and other interchangeable components plug into each equipment unit.

Sufficient gain and regulation are provided to permit repeater spacing of 170 miles in nonsleet areas. The system broadband regulator provides up to 48 db of flat loss correction at 99 kc and up to 24 db of slope correction between 99 and 150 kc. In addition, individual regulators in each channel will automatically compensate for up to 7 db of attenuation due to line absorption peaks or other irregularities.

(Continued on page 40A)

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Over forty years of precision engineering experience is now available to you through new Special Experimental & Research Service facilities of Kahle Engineering Company. Kahle's cumulative knowledge springs from its pioneering successes in electronics and allied industries coupled with vast experience with the actual end products which its machines produce.

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- Any special equipment for manufacture or research for tubes or lamps

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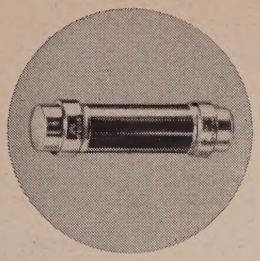
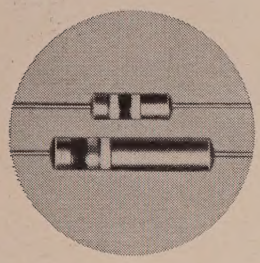


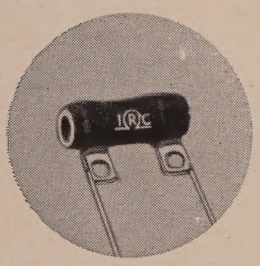
JAN-R-29 *specification*

For all requirements of JAN-R-29 Specification, Amendment 4, IRC sealed precision Voltmeter Multipliers function efficiently even when exposed to the most severe humidity. Used with 1-milliampere DC instructments, they enable voltage measurements to be made up to 6000 volts. Send for Bulletin.



JAN-R-184 *specification*

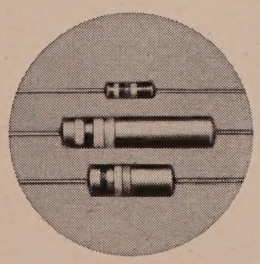
Unusually stable and inexpensive, IRC BW Wire Wounds meet JAN-R-184 Specification, Amendment 5, at ½ and 1 watt. Resistance element is uniformly and tightly wound on insulated core. Molded housing provides full insulation. Widely used in meters, analyzers, high stability attenuators, low-power ignition circuits, etc. Send for Bulletin.



MIL-R-26B *specification*

For high power dissipation, IRC Power Wire Wounds meet every commercial requirement of MIL-R-26B Specification, Characteristic G. Tubular, flat, fixed, adjustable, inductive, non-inductive, lead, lug and ferrule types provide resistors for virtually any circuit. From 5 to 225 watts. Send for Bulletin.

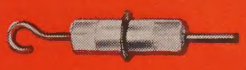
MIL TYPE RESISTORS



MIL-R-11A *specification*

IRC Advanced BT Resistors meet and beat MIL-R-11A Specification, Amendment 2. Filament-type resistance element and other exclusive features afford extremely low operating temperature and superior power dissipation in a compact, light, fully insulated unit. Available at ¼, ½ and 1 watt to MIL specification and 2 watts to commercial specification. Send for Bulletin.

NEW *product*



HERMETIC *sealing terminal*

Overcomes limitations of other types of hermetic sealing terminals.

Molded KEL-F* body—chemically inert to organic solvents, acids, oils, fumes.

Rugged construction—tough and resilient; withstands constant vibration.

Type HS-1 Feed-Thru Terminals, provide assured hermetic sealing for electrical and electronic components. Exclusive IRC molding Technique bonds Kel-F* to metal in a superior seal. Designed to the sealing requirements of MIL-T-27. Send coupon for full data

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Boron & Deposited Carbon Precistors • Power Resistors • Voltmeter Multipliers • Low Wattage Wire Wounds • Insulated Composition Resistors • Volume Controls •

Wherever the Circuit Says

Precision Wire Wounds • Ultra HF and Hi-Voltage Resistors • Selenium Rectifers • Insulated Chokes • Hermetic Sealing Terminals •

IRC

INTERNATIONAL RESISTANCE CO.

405 N. Broad St., Philadelphia 8, Pa.

In Canada: International Resistance Co., Ltd., Toronto, Licensee

Send me data on ☐ MF Voltmeter Multipliers, ☐ BW Resistors, ☐ Power Wire Wounds, ☐ Advanced BT Resistors, ☐ HS-1 Terminals.

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TYPE 324-A VIDEO PHASE METER

This instrument of laboratory precision makes possible the rapid and accurate measurement of phase angle THROUGH THE VIDEO RANGE. It provides verification of design calculations, a criterion for optimum adjustment of delicate phase relationships, or a maintenance checking tool in the wide-band electronic system in modern use.

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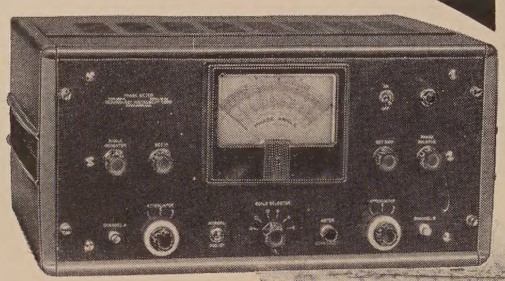

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BROAD BAND NEGATIVE FEEDBACK CIRCUITS
POLYPHASE SELECTIVITY METHODS
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PHASE DELAY LINE MEASUREMENT

SPECIFICATIONS

METER RANGES:	Phase angles from 0° to 360° full scale; and 90° quadrants full scale; no ambiguity.
FREQUENCY RANGE:	20 Kc. to 4.5 Mc. — Range down to 20 cycles may be supplied on special order.
WAVEFORMS ACCEPTED:	Sine waves and any complex waves having not more than one positive-going zero axis crossing per cycle. Phase angle measurement is defined as phase difference between corresponding positive going zero axis crossings of the periodic signals being compared.
AMPLITUDE RANGE:	2 volts to 300 volts peak.
ACCURACY:	± 4° on quadrant scales. Incremental change of 0.25° is easily read.
INPUT IMPEDANCE:	10 megohms shunted by 14 mmf.

FULL DETAILS UPON REQUEST

TECHNOLOGY INSTRUMENT CORP.

535 MAIN ST., ACTON, MASS., ACton 3-7711

News—New Products

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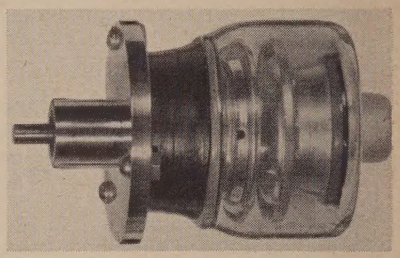
(Continued from page 26A)

Capacitor For Amateur Service

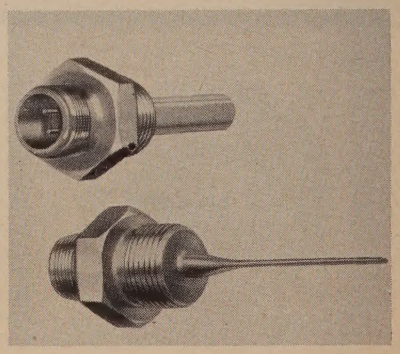
Jennings Radio Manufacturing Corp., P.O. Box 1278, San Jose 8, Calif. has introduced a new vacuum variable capacitor designed for service in the amateur field.



The Type VAC has a capacity range of 4 to 40 $\mu\mu$f with a nonlinear variation that makes tuning possible at the low capacity end. Its rating of 42 amperes, rms at 10 kv peak (and 22 mc) also makes it useful as a neutralizing capacitor in commercial applications. The voltage rating of 10 kv peak is determined at maximum capacity and increases rapidly as the plates are separated at lower capacities. The unit is 5 inches long and $2\frac{5}{8}$ inches in diameter, with a $\frac{1}{4}$ inch-diameter tuning shaft.

Temperature Detector

A new, hermetically sealed, electrical resistance temperature detector with a fast response time has been announced by **Thomas A. Edison, Inc., Instrument Div.,** West Orange, N. J.

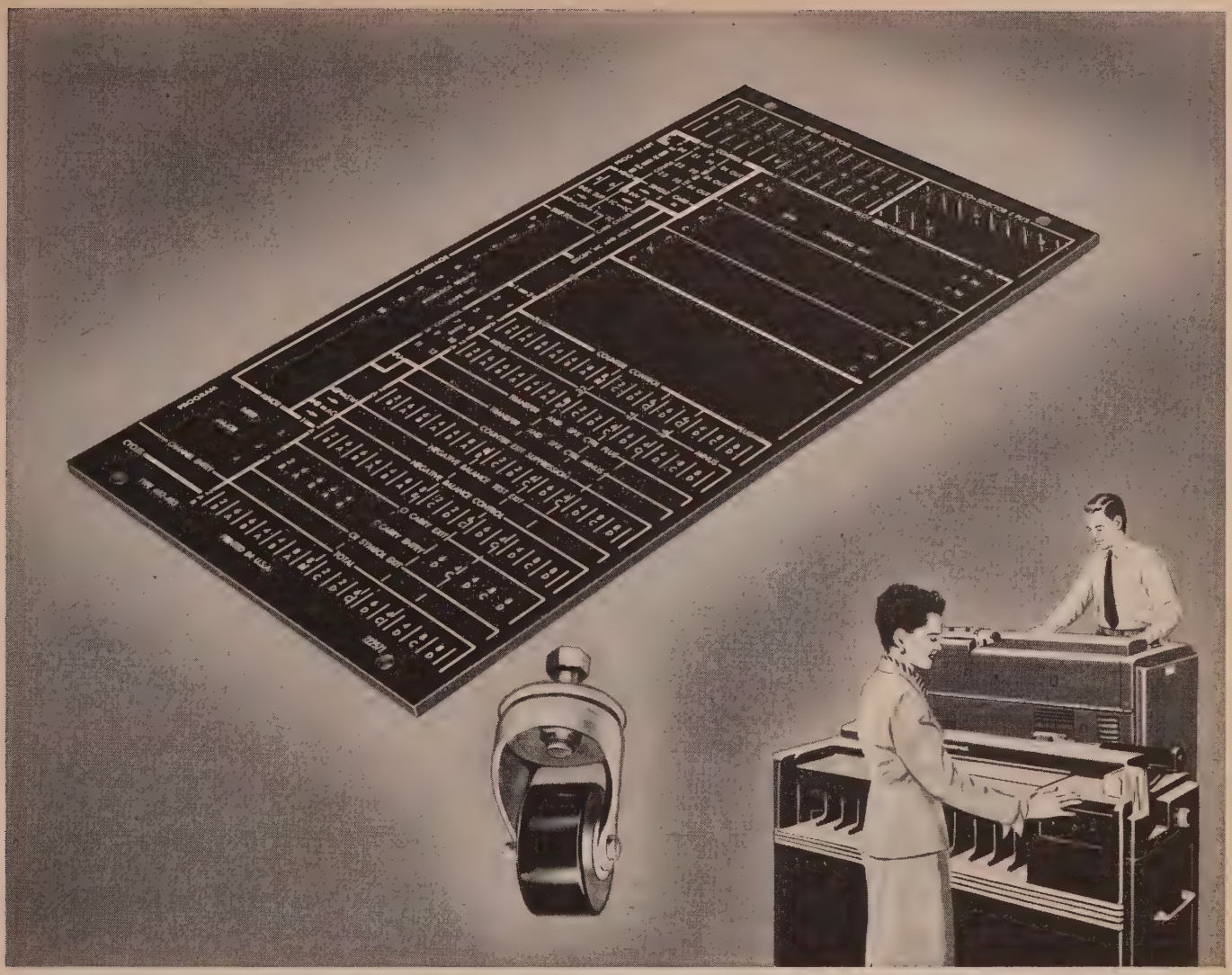


Designed for use with temperature recording, indicating, and control equipment where fast response is a primary requirement, the detector has an exponential time-control of 0.8 second or better in an agitated water bath.

Sealed in a stainless steel housing with a glass-to-metal sealed base through which the electrical connection is made, the new detector offers maximum resistance to corrosion and is not affected by most forms of destructive radiation.

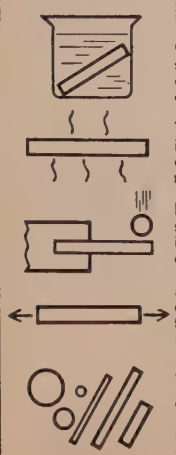
Its useful temperature range is from −70°C to +200°C. Temperature accuracy is ±1.0 per cent or better in mid-scale region. The sensing element is a nickel winding having a basic resistance of 90.38 ohms at 0°C.

(Continued on page 52A)



Head and feet for an office worker who neither errs nor tires

Properties of Synthane



In addition to those mentioned in the text, Synthane has the following important properties:

Chemical resistance. Synthane resists most acids and alkalis in moderate solutions, and corrosive atmospheres.

Temperature resistance. Synthane is thermosetting, will not flow under elevated temperatures. Grades resisting up to 400° F are available.

High impact strength. Synthane stands up well in applications where it is subject to vibration, pounding and shock loads.

Mechanical strength. Synthane exhibits excellent strength under tension, compression, and other loads. It will not delaminate.

Availability. Synthane is supplied in more than 33 grades of sheets, also in rods, tubes and molded-laminated or molded-macerated parts. A complete fabricating service is available.

• The uncanny ability of tabulating machines to do complicated jobs quickly and accurately is famous. One of the materials which helps to make this possible is *Synthane*—a laminated plastic.

Synthane serves as the base for the brains of the machines—the plug boards upon which the control circuits are set up. *Synthane* is excellent for the purpose because of its combination of high dielectric strength, resistance to moisture, dimensional stability and ease of machining. *Synthane* is printable, too—circuit designations are readily printed on its surface.

On tabulating machines, casters that are friendly to office-type flooring are needed. Casters of molded-macerated *Synthane* fill the bill. *Synthane* caster wheels are strong, do not flatten by constant pressure, and do not mar office floors.

Should you require a versatile material—one with many properties in combination—*Synthane* may be your answer. Our catalog tells the full story. To receive yours, drop us a note on your letter-head. Synthane Corporation, 8 River Road, Oaks, Pa.

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computers . . .

for fire control,
bombing,
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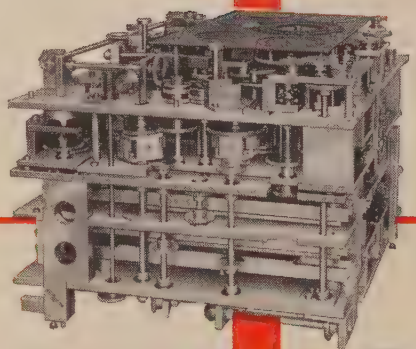

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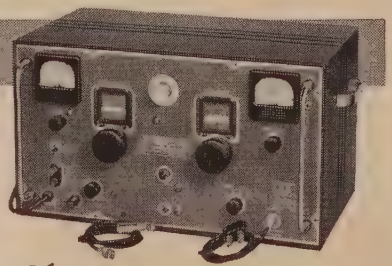
UHF TV SWEEP GENERATOR **MODEL 130**

Features continuous frequency coverage in one band; at least one volt output into 75 ohms; wide sweep; blanked signal on return sweep provides a reference baseline.

SPECIFICATIONS
Freq. Range: 450-900 mc.
Sweep Width: 0-40 mc min.
Sweep: 60 cycle, sine wave.
Output: (1.) 0.1-1.0 volts
(2.) 0.01-0.1 volts approx.

FM SIGNAL GENERATOR **MODEL 100C**



Designed to give precision performance over a single tuning range (27-230 mc). Negligible leakage; low spurious outputs; no auxiliary frequency changer unit required.

Write for specifications and catalog on complete line of measuring equipment.

NEW LONDON INSTRUMENT Company
P. O. BOX 189P
NEW LONDON, CONN.

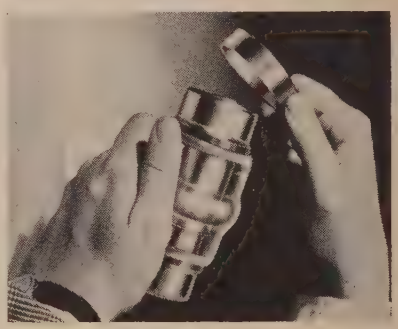
News—New Products

These manufacturers have invited PROCEEDINGS readers to write for literature and further technical information. Please mention your I.R.E. affiliation.

(Continued from page 26A)

Vacuum Capacitors

A new range of high-energy content, vacuum capacitors has been announced by **United Electronics Co.**, 42 Spring St., Newark 2, N. J., different type numbers each with the same over-all physical dimensions; nominal $5\frac{1}{4}$ inch length, $2\frac{3}{4}$ inch diameter. Each type is rated for 35,000 volts peak; 100 amperes EMS.



Exhaust tube seal-off tip is not exposed.

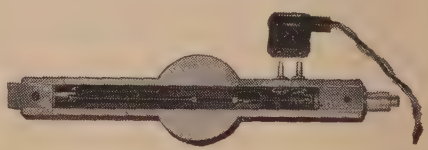
The capacitors are featured by wide circumference low-resistance contact provisions. Copper is used for all internal active area as well as for external terminals. Mounting and contact is accomplished by use of 2-inch diameter phosphor bronze rings supplied with the capacitors.

The 5 types differing as to capacitance values, 250, 200, 150, 100, and 50 $\mu\mu$f, carry the following designations: CAP 250/100/35 EC, CAP 250/100/35 EC, CAP 150/100/35 EC, CAP 100/100/35 EC, CAP 50/100/35 EC.

Temperature coefficient in these new units is low, and the periphery contact area assures efficient connection to circuitry.

New Galvanometer

A new and highly sensitive galvanometer, designed for use with relatively high impedance transducers, has been developed by the **Heiland Research Corp.**, 130 East Fifth Ave., Denver 9, Colo. This improved type is inter-changeable with all other types of Heil and galvanometers.



The new galvanometer, type 40–1000, designed and manufactured to stand up under the most exacting operational and environmental conditions, has a frequency of 40 cps, a sensitivity of 2.8 microamperes per inch (12 inch optical arm) and external damping resistance of 1,000 ohms, and a coil resistance of 109 ohms.

Other features include: Solid case construction completely dustproof, rotational limit stop, improved connector terminal, rapid horizontal and vertical adjustment.

(Continued on page 66A)

IT PAYS TO ASK THE STACKPOLE MAN!

These Stackpole sales engineering representatives can help you in many ways—from making cost-saving suggestions to providing engineering samples and helping you procure dependable, economical components to match your production requirements and scheduling. Write, wire or call the engineer in your territory next time you need any of the materials listed below.

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- **FIXED AND VARIABLE RESISTORS**
- **SPECIAL RESISTORS**
- **LINE, SLIDE OR ROTARY-ACTION SWITCHES**
- **IRON CORES** — All types and shapes
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- **CERAMAG® ferromagnetic CORES**
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H. M. Richardson
9 East 22nd Street
PHONE: GEneva 4078

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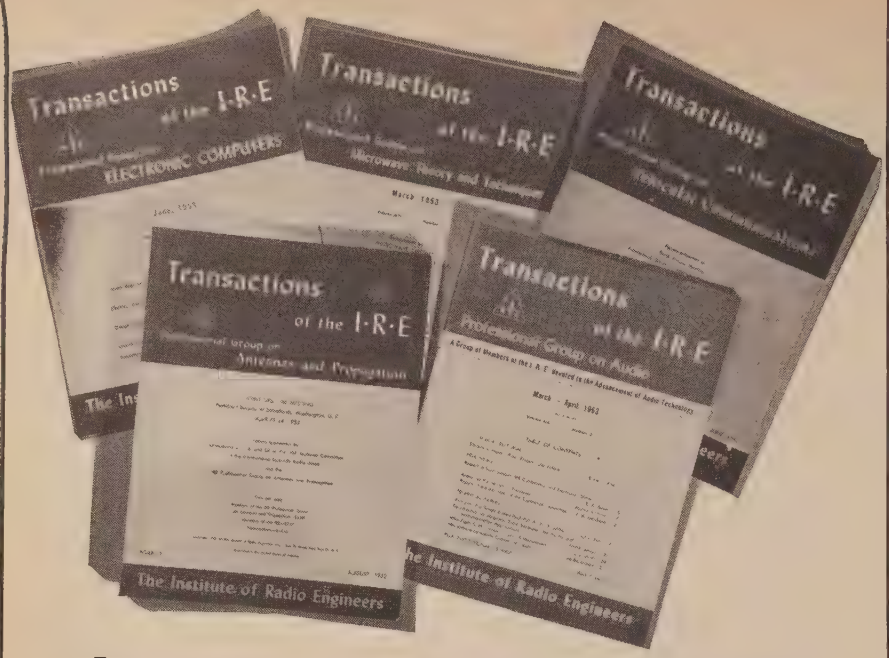

At least one of your interests is now served by one of IRE's 21 Professional Groups

Each group publishes its own specialized papers in its *Transactions*, some annually, and some bi-monthly. The larger groups have organized local Chapters, and they also sponsor technical sessions at IRE Conventions.

Aeronautical and Navigational Electronics (G 11)	Fee $2
Antennas and Propagation (G 3)	Fee $4
Audio (G 1)	Fee $2
Broadcast & Television Receivers (G 8)	Fee $2
Broadcast Transmission Systems (G 2)	Fee $2
Circuit Theory (G 4)	Fee $2
Communication Systems (G 19)	Fee $2
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Vehicular Communications (G 6)	Fee $2

IRE Professional Groups are only open to those who are already members of the IRE. Copies of Professional Group Transactions are available to non-members at three times the cost-price to group members.

The Institute of Radio Engineers
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Professional Group on Broadcast and Television Receivers

The field of broadcast receivers is one which is closely associated with the general public, perhaps more so than any other branch of the radio engineering field. In fact, to the layman the term "radio" is synonymous with "broadcast receiver."

As a result, the receiver engineer has been concerned with an additional factor not generally common to other fields, namely, that of responding to—or endeavoring to create—public demand for a product. This factor has played a prominent role in such developments as FM, car radios, portable receivers, and black-and-white television sets. It is now conspicuously evident in connection with current efforts to produce and market color television receivers.

The IRE Professional Group on Broadcast and Television Receivers is playing a major role in making available vitally needed technical information, not only on color television, but on all aspects of the receiver field. Through this exchange of information, the radio and television industry is gaining important data which will be helpful in solving the engineering problems it faces and in successfully meeting the "public demand" factor mentioned above.

The Group has been particularly active in sponsoring technical sessions at most of the national meetings held throughout the country during the year: the Radio Fall Meeting, the Spring Television Conference in Cincinnati, and the IRE National Convention, to mention but a few.

The Group also publishes its own technical publication, called *Transactions*, which is distributed to some 1200 members as a part of their $2.00 assessment fee. The *Transactions*, which is published on a quarterly basis, has become a chief source of information on the latest technical developments in the field of broadcast and television receivers.

W. R. G. Baker

Chairman, Professional Groups Committee

trifles make PERFECTION... *but* PERFECTION *is no trifle*

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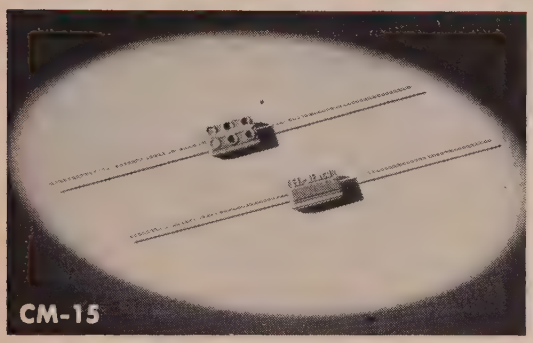
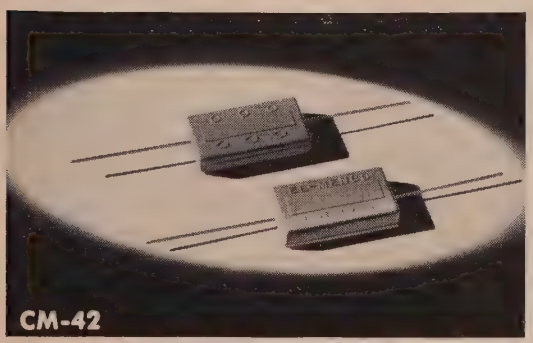
is VITAL *to* 61 TONS *of* MAGNIFICENT PERFORMANCE



Smallest Molded Mica Capacitors 9/32" x 1/2" x 3/16"



Made to Meet All MIL-C-5 Requirements. Largest Molded Mica Capacitors of Wire Terminal Type. 13/16" x 1-1/2" x 5/16"

Jobbers and Distributors are requested to write for information to Arco Electronics, Inc., 103 Lafayette St., New York, N. Y. — large stocks on hand — spot shipments for immediate delivery. Sole Agent for Jobbers and Distributors in U. S. and Canada.

When the mighty giants of the air lift their massive wings to fly, a thousand and more "tremendous trifles" instantly go to work in harmonious unison to give life and power. It is the perfection of these "trifles" that makes possible the magnificent performance of today's luxurious air liners.

The EL MENCO Capacitor—CM-15—is one of these "tremendous trifles" that plays such a vital part in the efficient operation of aircraft communication.

EL MENCO IS THE ONE OUT OF MANY CHOSEN FIRST

Superiority of manufacture and dependability of performance make EL MENCO first choice on the specification sheet . . . because EL MENCO Capacitors are factory-tested at *double their working voltage* — they are *guaranteed stable* under the most adverse conditions. Whether you use our *high* capacity CM-42 (10-25,000 mmf) or our midget *low* capacity CM-15 (2-525 mmf) you have guaranteed assurance of job-tested, job-rated capacitors — tremendous trifles of perfection so vital to the magnificent performance of YOUR product.

ELECTRO MOTIVE is now supplying special silvered mica films for the electronic and communication industries — just send us your specifications.

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El-Menco

MOLDED MICA CAPACITORS

MICA TRIMMER

Foreign Electronic Manufacturers Get Information Direct from our Export Dept. at Willimantic, Conn.

THE ELECTRO MOTIVE MFG. CO., INC. **WILLIMANTIC, CONNECTICUT**

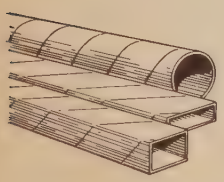
MADE TO YOUR EXACT SPECIFICATIONS
IN ANY SIZE • SHAPE • QUANTITY

Precision coil bobbins are fabricated from high dielectric materials and quality controlled to the most minute tolerances . . . Yet, because they are made on special high production equipment, they're available to you for prompt delivery at low unit cost.

Cores are spirally wound dielectric kraft, fish paper, acetate, phenol impregnated or combinations. Flanges are cut to any specification for all types of mountings.

Request illustrated bulletin. Send specifications for samples.

High Strength Low Cost Paper Tubes



Accurately fabricated in any size, shape, ID or OD. Spirally wound from select dielectric materials. Crush resistant, with excellent dimensional stability. Subject to rigid control and inspection for tolerance and uniformity.

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Philadelphia, Pa., Chestnut Hill 8-0282

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Plant No. 2: 79 Chapel St., Hartford, Conn.

News—New Products

These manufacturers have invited PROCEEDINGS readers to write for literature and further technical information. Please mention your I.R.E. affiliation.

(Continued from page 52A)

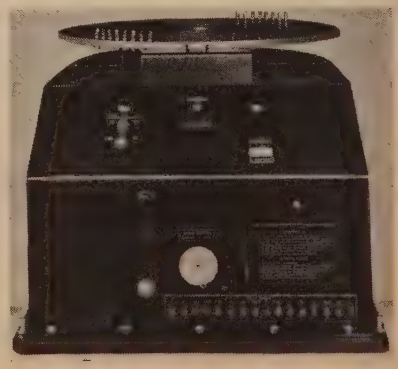
Sub-Miniature Resistor

A new resistor, Type A3606CG, has been developed by the **I-T-E Circuit Breaker Co., Resistor Div.,** 1924 Hamilton St., Philadelphia 30, Pa., to meet the needs of sub-miniaturizing programs being carried out with many electronic and electronic-actuated devices throughout industry. Hermetically sealed, this resistor measures 3/16 inch in diameter by $1\frac{1}{8}$ inches long.

Resistors are built to customers' specifications, rated at 0.10 w. Maxium resistance, using Evenohm wire or its equivalent, is 500,000 ohms. Tolerances down to 0.1 per cent are standard. Axial lead wires are of 22-gauge tinned copper, $1\frac{1}{2}$ inches long.

Synchronous Motor For Rate-of-Turn Table

A new 115-volt, 60-cps, single phase 180 rpm synchronous motor for the Genisco Model C, Rate-of-Turn Table, a precision component used for calibrating and evaluating rate gyros, and for calibrating accelerometers up to ten G's, has been developed by **Genisco, Inc.,** 2233 Federal Ave., Los Angeles 64, Calif.



The new motor, designed specifically for this low-speed application, eliminates the usual belt-and-pulley or gear reduction systems. The motor switch may be

(Continued on page 68A)

Fast Delivery AlSiMag®

DIE PRESSED Ceramics

FAST DELIVERY . . . Our own die shop and four modern plants speed deliveries.

LARGE OR SMALL QUANTITIES . . . We have the most complete press facilities in the industry.

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SEE OUR DISPLAY

BOOTH NO. 340

BASIC MATERIALS EXPOSITION

The Product Development Show

CHICAGO • MAY 17-20, 1954

Preformed Contact Finger Stock is an ideal electrical weather stripping around doors of equipment cabinets as well as being excellent for use with VHF and UHF circuitry. Silver plated, it comes in three widths—$\frac{17}{32}$, $\frac{31}{32}$ and $1\frac{7}{16}$ inches.

Variable vacuum capacitors come in three models, are lightweight, compact, eliminate the effects of dust and atmospheric conditions and have low inductance. Also available are eight types of fixed vacuum capacitors.

Air-system sockets, designed for Eimac tube types 4-400A, 4-1000A, 4X150A, and 4X150D, simplify cooling and assure adequate air-flow to various seals. The 4-400A socket can also be used with the 4-125A and 4-250A radial-beam power tetrodes if desired.

HR heat dissipating connectors provide efficient heat transfer from the tube element and glass seal to the air while making electrical connections to plate and grid terminals. Precision machined from dural rod, HR connectors come in ten sizes to fit most of Eimac's internal anode tubes.

High Vacuum Rectifiers come in eight models, are instant heating, have radiation-cooled pyrovac* plates and can be operated in a variety of rectifying and voltage multiplying circuits. Also available are four types of mercury-vapor rectifiers.

*An Eimac trade name.

• For further information write our Application Engineering department

EITEL-McCULLOUGH, INC.
SAN BRUNO · CALIFORNIA
Export Agents: Frazar & Hansen, 301 Clay St., San Francisco, California

News—New Products

These manufacturers have invited PROCEEDINGS readers to write for literature and further technical information. Please mention your I.R.E. affiliation.

(Continued from page 66A)

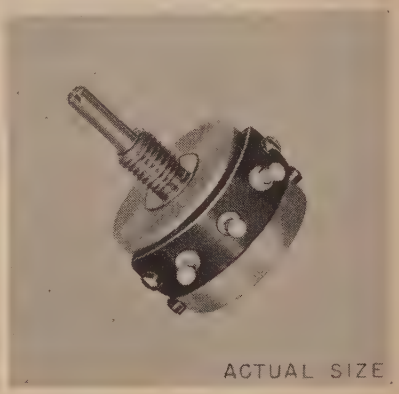
reversed instantly, however, the complete stop-start cycle is programmed. Twelve seconds are required to reach full synchronous speed; eight seconds to stop. The rotor is energized with direct current in the synchronous run condition and the resulting torque is in excess of fifty-inch pounds. As a result, the rate-of-turn machine operation is extremely smooth. Output from the constant speed motor to the 24-inch cast aluminum table is infinitely variable through a ball-disc integrator drive from 0.01 to 1,200 degrees per second. Thirty-six turns of the single control are required for full-scale traverse.

Table capacity is 100 lbs. 108 $\frac{1}{4}$ inch-20 tapped mounting holes are arranged in a regular 2 inch equilateral pattern for mounting.

Write the manufacturer for detailed specification sheets.

Miniature Precision Potentiometer

Jet Electronics, Inc., 93 Massachusetts Ave., Boston, Mass., announces production of their new $\frac{7}{8}$ inch precision potentiometer, available with up to 360° electrical rotation as well as mechanical rotation, in resistance ranges up to 50,000 ohms (± 1 per cent). The weight is $\frac{1}{2}$ ounce.



Cataloged as the D-100, it has an onodized aluminum case $\frac{7}{8}$ inch in diameter, with phosphor bronze bushing. A centerless, ground stainless steel shaft $\frac{1}{8}$ inch in diameter is supported by front and rear sleeve bearings (or ball bearings if specified). Paliney contact material is used between the rotor and resistance element. The rotor take-off brushes are made of gold alloy and rotate on a coin-silver ring. The manufacturer claims a life in excess of one million cycles.

Windings are available, with linearity guaranteed to within ± 1 per cent of total resistance. The units are rated at 1 watt at 25°C and the resistance alloy has temperature coefficient of 0.00002 parts/°C for values of 500 ohms or more. The ambient temperature rating is from −67°F to +250°F. Terminal board material withstands 250° F. Terminals are available with silver or gold plating.

(Continued on page 70A)

leading Hi-Fi producers specify **CRUCIBLE PERMANENT MAGNETS**

for maximum energy...minimum size

Crucible alnico permanent magnets are unsurpassed in their magnetic properties. They provide *consistently higher* energy product . . . which results in smaller, more powerful magnets.

That's why not only leading producers of high fidelity sound equipment, but hundreds of other discriminating manufacturers of instruments, controls, motors and other magnet-equipped devices, prefer Crucible alnico permanent magnets.

Crucible has been one of the largest producers of this type of magnet since the alnico alloys were first developed. And backing up this twenty-year magnet experience are 54 years of *quality* steelmaking. For highest quality alnico magnets and technical help in their application, *call Crucible.*

CRUCIBLE first name in special purpose steels

54 years of *Fine* steelmaking

ALNICO PERMANENT MAGNETS

CRUCIBLE STEEL COMPANY OF AMERICA, GENERAL SALES OFFICES, OLIVER BUILDING, PITTSBURGH, PA.

REX HIGH SPEED • TOOL • REZISTAL STAINLESS • ALLOY • MAX-EL • SPECIAL PURPOSE STEELS

BOGUE

HI-CYCLE GENERATORS

For precision Hi-Cycle power generating equipment to meet one of a myriad of special power requirements ranging from highest quality laboratory power to precision electronic testing, industry relies on Bogue Precision Power. Bogue as the recognized leader in the Hi-Cycle field offers these performance characteristics—

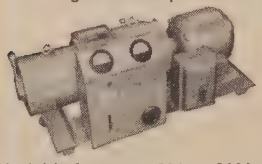
LOW HARMONICS

CLOSE VOLTAGE REGULATION

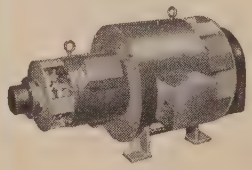
400 CYCLE REGARDLESS OF LOAD & INPUT VARIATIONS

For example, Bogue special 400 cycle single shaft, two-bearing synchronous motor driven units eliminate belts, gears and other special speed changers, yet, faithfully deliver 400 cycles—exactly—no load to full load regardless of voltage variations .... truly the standard of 400 cycle power .... the reason so many prominent companies have been depending on equipment built by Bogue Electric Manufacturing Company ...

5 KW portable regulated 400 cycle motor-generator set with integral control panel.

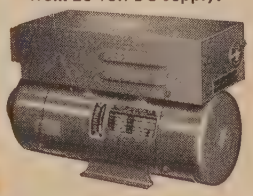


Variable frequency 320 to 1000 cycle M-G set. Bogue magnetic amplifier maintains voltage and frequency to within one-half of 1% of any preset value.



5 KW low harmonic set. 400 cycle regardless of input voltage, loading or heating.

400 cycle voltage & frequency regulated inverter. Operates from 28 volt DC supply.



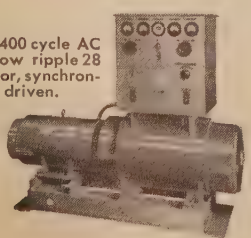
The Authority on High Cycle Power

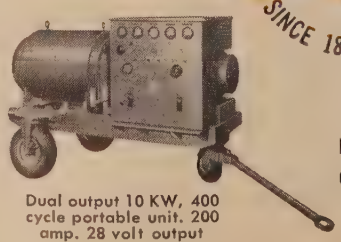
BOGUE

PRECISION ELECTRICAL EQUIPMENT — Bogue — SINCE 1892

50 IOWA AVENUE • PATERSON 3, NEW JERSEY

Low harmonic 400 cycle AC generator and low ripple 28 volt DC generator, synchronous motor driven.





Dual output 10 KW, 400 cycle portable unit. 200 amp. 28 volt output

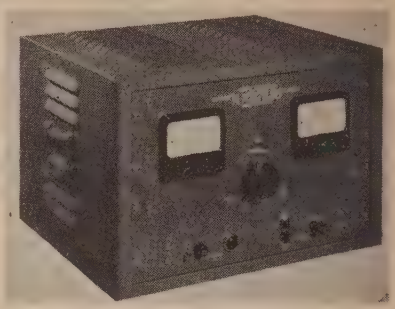
REQUEST ILLUSTRATED CATALOG 440

News—New Products

These manufacturers have invited PROCEEDINGS readers to write for literature and further technical information. Please mention your I.R.E. affiliation.

(Continued from page 68A)

Power Supply

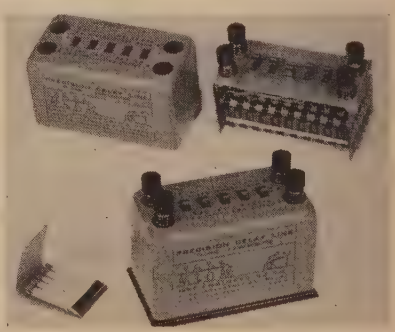


A new type of magnetic-amplifier-regulated power supply has been developed by the **Perkin Engineering Corp.**, 345 Kansas St., El Segundo, Calif. Identified as MR532-15, this unit is rated at 5–32 volts dc at 15 amperes and has a regulation accuracy of ±1 per cent from 5 to 32 volts dc, from 1/10 full load to full load, from 105 to 125 volts ac. Its ripple voltage is 1 per cent rms at 32 volts at full load and its recovery time is 1/10 second. The unit is provided with a 4½-inch ammeter and voltmeter and is mounted in a bench-type cabinet which can also be used for 19-inch rack panel mounting by removing the panel from the cabinet.

The special features of this unit are, the fact that it has no tubes, is regulated and stabilized, is continuously variable, uses a choke input filter, and it has a magnetic type circuit-breaker on the dc side with time-delay provision for starting motors. Further information and literature can be obtained by writing the Perkin Corp.

Precision Delay Line

A laboratory-type delay line, variable in additive increments of 0.2 microsecond each, and used for electronic circuit work such as color television, instrumentation, pulse forming net-works, computer circuits, is being manufactured by **May Engineering Co.**, 6055 Lankershim Blvd., N. Hollywood, Calif. This instrument is the lumped constant type with a total maximum delay of 1.0 microsecond measured at ½ amptitude, with a rise time of 0.05 microseconds (measured at 10 per cent and



90 per cent amptitude). The characteristics impedance is 50, 75, or 100 ohms, and maximum peak voltage is 500.

(Continued on page 72A)

ATR AUTO RADIO VIBRATORS

Have Ceramic Stack Spacers

A COMPLETE LINE OF VIBRATORS

Designed for Use in Standard Vibrator-Operated Auto Radio Receivers. Built with Precision Construction, featuring Ceramic Stack Spacers for Longer Lasting Life. Backed by more than 22 years of experience in Vibrator Design, Development, and Manufacturing.

Free

"A" Battery Eliminators, DC-AC Inverters, Auto Radio Vibrators

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✓ NEW MODELS ✓ NEW DESIGNS ✓ NEW LITERATURE

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Quality Products Since 1931

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when the going gets

ROUGH

VARIAN klystrons can take it

The true test of a production klystron is the ability to operate successfully when subjected to severe vibration and shock under field conditions. That's why manufacturers of mobile radar insist on VARIAN klystrons—klystrons that stand up when the going gets rough.

VARIAN KLYSTRONS ARE **RUGGED**

Varian makes **sure** that its klystrons meet field performance requirements by testing each one under severe high amplitude vibration. This production test, accurately duplicating field conditions, is rough — so rough that **ordinary** klystrons can't take it.

VARIAN MEANS **PROVED** PERFORMANCE

From design to finished product, Varian builds quality into every klystron. And quality means dependability — the reason why leading system manufacturers specify Varian when klystron performance is a critical factor in the operational reliability of their product.

For rugged, dependable, production klystrons, specify:

VA-6310/V-260
VA-6312/V-270
VA-6313/V-280
VA-6314/V-290
VA-6315/V-153
VA-6316/V-151

IN KLYSTRONS, THE MARK OF LEADERSHIP IS

VARIAN associates

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Representatives in all principal cities.

THE CORRECT COMBINATION OF SIZE AND POWER RATING offered by Jennings Vacuum Capacitors

The wide choice of size and power ratings for a given capacity range is illustrated by these four units all having maximum capacities of 1000 mmfd. JENNINGS functional designs thus permit you to select the smallest vacuum capacitor that will meet your voltage and current requirements.

Please let us suggest the capacitor that will best meet your specific circuit conditions.

UCSXF 10-1000 mmfd, 7.5, 10, 12 KV 60 amps rms

USCL 7-1000 mmfd, 3, 5 KV 42 amps rms

Jennings RADIO VACUUM ELECTRONIC COMPONENTS

VMMC 50-1000 mmfd 10, 15 KV, 125 amps rms

VMMHC 60-1000 mmfd 35, 40, 45 KV 125 amps rms

Literature mailed upon request

JENNINGS RADIO MANUFACTURING CORPORATION · 970 McLAUGHLIN AVE.
P.O. BOX 1278 · SAN JOSE 8, CALIFORNIA

News—New Products

These manufacturers have invited PROCEEDINGS readers to write for literature and further technical information. Please mention your I.R.E. affiliation.

(Continued from page 70A)

This delay line is constructed so that the individual switches each control a step of 20 coils and 20 matched capacitors. Low attenuation is a feature of this unit. It is also produced in the distributed parameter type for high characteristic impedance. The over-all dimensions are $2\frac{1}{2}\times3\times4\frac{1}{2}$ inches. These "Type A" precision delay lines are in production and priced at $75.00.

In addition to the above, the same chassis is used and the steps are 0.05, 0.01, 0.15, 0.3, and 0.5 microsecond, to total 1.0 microsecond, thereby making it possible to switch in any delay from 0 to 1.0 microseconds in increments of 0.05. These are "Type B" precision delay lines, and are priced at $90.00 each.

Single Side Band Adapter

Type 76 Single-Sideband Adapter manufactured by **Crosby Laboratories, Inc.,** Box 233, Hicksville, L. I., N. Y., incorporates the new Burnell toroidal coil filters in place of the crystal filters used in the Type 51 Adapter.

One of the advantages claimed of the toroidal coil filter (Adapter) is the substantial reduction in cost as compared to the present crystal filter units, a factor that will accelerate the trend to the much superior single-sideband type of communication systems.

In addition to the reduced cost feature of the toroidal coil, this type filter is considerably smaller in size than the crystal filter, resulting in a saving in weight and chassis space. The unit is more suitable for field use where equipment may be subject to operation under shock and vibration conditions.

Moreover, since the Type 76 Adapter has none of the multiple crystals and L-C network circuits, alignment procedures are substantially eliminated.

For more complete information write to the firm direct.

Plant Expansion

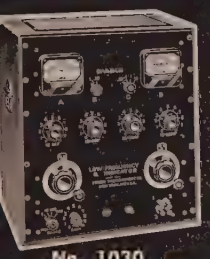
Geo. Stevens Mfg. Co., Pulaski Rd. at Peterson, Chicago 30, Ill., has just completed an addition which doubles the area of the present modern plant. The new space is being used for manufacturing high-speed coil winders and for an enlarged engineering and design department.

(Continued on page 74A)

FREED Instruments & Transformers

Famous For

QUALITY · DEPENDABILITY · ACCURACY



No. 1030
Low Frequency
"Q" Indicator

No. 1020B
Megohmmeter

Decade
Inductors

No. 1040
Vacuum Tube Voltmeter

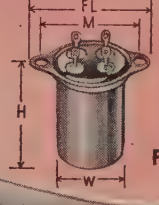
FREED MINIATURE AUDIO TRANSFORMERS

PROFESSIONAL MINIATURE AUDIO TRANSFORMERS

These high quality, miniature transformers feature hermetic sealing for maximum protection from moisture penetration with subsequent electrolysis and corrosion of fine wires. While primarily intended for non-military equipment, these units are constructed in accordance with MIL-T-27 Specifications.

CATALOG NO.	APPLICATION	IMPEDANCE LEVEL OHMS PRIMARY	IMPEDANCE LEVEL OHMS SECONDARY	MAXIMUM POWER LEVEL V.U.*	RATIO	MAX PRI DC PER SIDE Ma.	D.C. UNBAL. Ma.	FREQ. RESPONSE C.P.S.	CASE NUMBER
PMA 1	Line or microphone to single or push-pull grids	50/ 200/ 500	60,000 C.T.	+8	1:11	0	0	±2.0 DB 30-20000	DM-12
PMA 2	Dynamic microphone or speaker voice coil to single or P.P. grid	4/8	60,000 C.T.	+8	1:86.6	0	0	±2.0 DB 30-20000	DM-12
PMA 3	Line or microphone to single or push-pull grids. Magnetically shielded.	50/ 200/ 500	60,000 C.T.	+8	1:11	0	0	±2.0 DB 30-20000	DM-12
PMA 4	Single triode plate to single or push-pull grids	15,000	60,000 C.T.	+8	1:2	0	0	±2 DB 30-10000	DM-12
PMA 5	Single triode plate to push-pull grids	15,000	60,000 C.T.	+8	1:2	2	2	±2 DB 200-10000	DM-12
PMA 6	Single triode plate to multiple line	15,000	50/200/500	+8	5.48:1	0	0	±2 DB 30-20000	DM-12
PMA 7	Single triode plate to multiple line	15,000	50/200/500	+8	5.48:1	2	2	±1 DB 200-10000	DM-12
PMA 8	Push-pull triode plates to multiple line	30,000 C.T.	50/200/500	+8	7.75:1	2	0.25	±2 DB 30-20000	DM-12
PMA 9	Crystal mike or pickup to multiple line	60,000	50/200/500	+8	11:1	0	0	±2 DB 30-20000	DM-12
PMA 10	Mixing or matching	50/200	50/200/500	+8	1:1.50	0	0	±2 DB 30-20000	DM-12
PMA 11	Parallel Feed Reactor	40 hy.	3 ma dc,	3500 ohms D.C. resistance					DM-12

* 1 mw. reference level.



DM-12 CASE DIMENSIONS

FL - 1 1/2
FD - 1 1/32
W - 15/16
H - 1 15/32
M - 1 7/32

Screws - 4-40
Cut out - 7/8
Wgt. - 1.5oz.

Send for complete catalog on
Freed Instruments and Transformers

No. 1210
Null Detector &
Vacuum Tube Voltmeter

No. 1010
Comparison Bridge

No. 1110A
Incremental Inductance
Bridge

FREED TRANSFORMER CO., INC.

SHOW TEMPERATURE AT THE CONTROL PANEL

4½" Cat. #4511
3¾" Cat. #1654
2½" Cat. #2613

WARN OF TROUBLE WHEN IT STARTS — USE SIMPLYTROL INDICATING PYROMETERS.

WHITE FACE

Medium resistance pyrometers (4 ohms per millivolt). Automatic, bimetal cold junction correction. Compensated for copper error. With adjusting resistor to take thermocouples up to external resistance shown in table. When specified will be adjusted for mounting in steel panel. Accuracy 2%.

Contact pyrometers have same size and appearance but include contacts, adjustable to close at any place on the scale arc. Ask for Bulletin G-5.

Model B

Model C

4½" Cat. #4514
3¾" Cat. #1652
2½" Cat. #2612

BLACK FACE

PORTABLES

Any of these pyrometers can be supplied in the Model B carrying case or in the Model C bench stand. Ask for Bulletin G-3.

Temperature Range Degrees		Int. Res.	Ext. Res.	Thermo-couple	2½ inch Size 261 Cat. No.	Price	3¾ inch Size 351 Cat. No.	Price	4½ inch Size 451 Cat. No.	Price
0—3000° F	0—1650° C	70	10	Pt13%Rh	2618	$19.90	1658	$22.75	4518	$30.25
0—2500	0—1370	212	10	C/A	2611	18.40	1651	21.25	4511	28.75
0—2000	0—1100	172	10	C/A	2616	18.40	1656	21.25	4516	28.75
0—1500	0— 800	179	10	I/C	2612	18.40	1652	21.25	4512	28.75
0—1000	0— 500	112	10	I/C	2615	19.50	1655	21.80	4515	29.90
0— 750	0— 400	81	10	I/C	2614	19.50	1654	22.45	4514	31.05
0— 500	0— 260	52	8	I/C	2613	20.10	1653	23.00	4513	32.20
— 75 to 225	— 60 to 110	30	5	I/C	2617	23.00	1657	26.50	4517	35.65
—200 to 100	—130 to 40	30	5	C/C	2619	29.00	1659	32.50	4519	41.65

Several intermediate ranges not listed.

Thermocouples:—(not included). Chromel-alumel (C/A) 14 ga. insulated 5′ leads $4.00. Iron-constantan (I/C) or copper-constantan (C/C) 20 ga. insulated 5′ leads $2.90.

ASSEMBLY PRODUCTS, INC., P. O. BOX 191, CHAGRIN FALLS 2, OHIO

News—New Products

These manufacturers have invited PROCEEDINGS readers to write for literature and further technical information. Please mention your I.R.E. affiliation.

(Continued from page 72A)

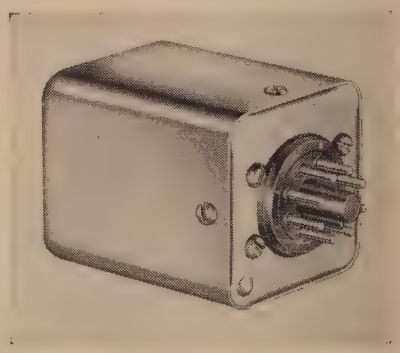
Frequency Multiplier

Power Frequency Multipliers which change the frequency of line power without the use of rectifiers are now being produced **D & R Ltd.,** 402 E. Gutierrez St., Santa Barbara, Calif., Units are hermetically sealed, and operate directly from a 400-cps, 3-phase power line, delivering 2,000-cps single phase output at a conversion efficiency of approximately 75 per cent.

The Model F-10 (illustrated) delivers 5 watts at 2,000 cps with good regulation, and with a harmonic content less than 7 per cent. The Model F-12 is only slightly larger than the F-10 yet delivers 25 watts at 2,000 cps, while the Model F-11 unit is capable of delivering 250 watts to the load with low harmonic content and an operating efficiency in excess of 80 per cent.

Relay Dust Cover

A new, low-cost, dust cover for their Model DOS Relay has just been announced by **American Relay & Controls, Inc.,** 3645 Howard St., Skokie, Illinois.



This component is a general-purpose relay that is light and compact but will handle relatively large power loads and is protected against dust. The plug-in mounting is also in line with the modern trend toward unit replacement.

The enclosure is of deep-drawn, one-piece aluminum, which fits snugly over the aluminum plate on which the relay and its plug-in base are mounted. No finishing is needed and, the enclosure will not rust.

(Continued on page 76A)

TAPE WOUND CORES

Are you ready for a revolutionary concept in the electrical and electronic industry—the Magnetics, Inc. "Performance-Guarantee" on Tape Wound Cores. Guaranteed to meet your specifications, *and sold at standard prices;* these Cores mean truly economical production of high permeability magnetic devices in your plant.

TABLE A

BASIC PHYSICAL CONSTANTS OF COMMON MAGNETIC MATERIALS

Trade Name	% Ni	% Fe	Other	Grain Structure	Satur. Flux Density Gausses	Resistivity Microhm-Cm	Curie Point °C	Dens. Grams per cc
Hy Mu 80	79	17	4 Mo	"random"	8,700	57	420	8.72
48 Alloy	48	52	....	"random"	16,000	45	500	8.3
Orthonol	50	50	....	oriented	15,500	45	500	8.25
Magnesil	..	97	3 Si	oriented	20,000	48	700	7.65

TABLE B

TRADE NAMES OF SIMILAR MATERIALS

Hy-Mu 80	48 Alloy	Orthonol	Magnesil
4-79 Permalloy	Carpenter 49	Orthonik	Armco Oriented T
Mo-Permalloy	Allegheny 4750	Permeron	Hypersil
Mu Metal*	Hypernik	Deltamax	Orthosil
		Hypernik V	Silectron

Typical of the unusual scope of the material contained in Catalog TWC-100 are Tables A and B, reproduced from Page 4 of "Performance-Guaranteed Tape Wound Cores."

GET THE COMPLETE STORY

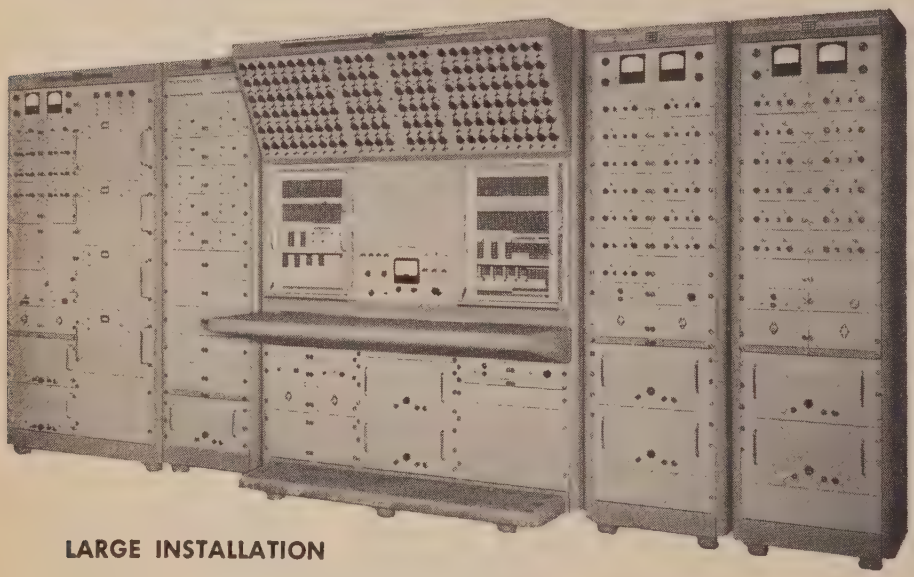
A wealth of new and unusual material on Tape Wound Cores is available to you in Catalog TWC-100, "Performance-Guaranteed Tape Wound Cores." Tables A and B of the catalog, reproduced on this page, present a striking illustration of material not to be found compiled together elsewhere.

Data and descriptive details on high permeability materials . . . factory core matching . . . free engineering design services . . . pages of characteristic graphs and tables . . . are yours for the asking. Simply write on your company letterhead.

MAGNETICS inc.

DEPT. I-7, BUTLER, PENNSYLVANIA

MODERN Problems Demand ... MODERN SOLUTIONS

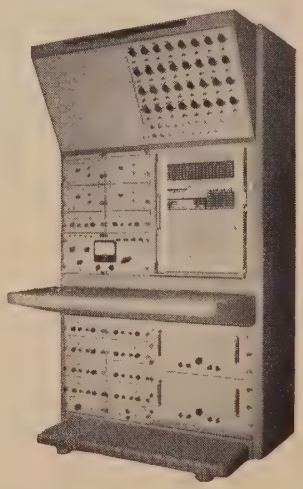


LARGE INSTALLATION

This large computer is used for the rapid solution of aero-dynamic problems. It consists of 50 operational amplifiers, 10 servo multiplying channels, 4 resolving channels, and a control console with two pre-patch bays, 156 attenuators, two voltmeters, and all necessary operational controls.

SINGLE PACKAGE COMPUTER

Our Type 16-31R Computer is a single package computer capable of solving differential equations with many simultaneous elements which are often encountered in the simulation of dynamic systems. It contains 20 operational amplifiers, 4 servo multipliers, thirty-two attenuators, all-metal removable problem board, and complete control panel.



PLOTTING EQUIPMENT

For presentation of problem solutions, the Variplotter Plotting Boards provide an accurate inked record. Typical uses include the automatic plotting of: Analog Computer output; guided missile data; engine performance characteristics; and control of manufacturing processes. With accessory equipment the range of applications can be greatly extended.

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News—New Products

These manufacturers have invited PROCEEDINGS readers to write for literature and further technical information. Please mention your I.R.E. affiliation.

(Continued from page 74A)

The enclosure can be removed for inspecting the relay.

The relay itself is Amrecon's Model DOS, which meets the rigorous standards for aircraft relays. The weight of relay and enclosure is seven ounces, and its dimensions are $2\times2\times2\frac{1}{2}$ inches high over-all, excluding the plug. Coil rating is 2.5 watts dc or 3.0 watts ac, with voltages up to 230 volts dc or 440 volts ac. Contact rating is 15 amperes at 115 volts ac or 32 volts dc, non-inductive.

Themistor-Type Temperature Indicator

Sensitivity and rapid response in temperature measurements are provided by the new Climate-Survey Temperature Indicators, developed by **Beckman and Whitley, Inc.,** 963 E. San Carlos Ave., San Carlos, Calif. The basic instrument is available in three models, as follows: Model 196, six overlapping ranges spanning from −70° to +60°F; Model 197, six overlapping ranges spanning from 10° to 140°F; and Model 198, having a single range from 60° to 90°F.

All are battery-operated and provide 500 hours of continuous operation.

Each instrument is provided with two input receptacles and a selector switch so that reference-temperature or wet- and dry-bulb readings can be made. Sensing elements can be located 800 feet from the instrument with standardized calibration, or with special calibration, cables as long as 8,000 can be used.

Direct-reading accuracy of all models is ±0.5°F through all ranges. Readings are made from a 3-inch square temperature-compensated meter at a response time of six seconds for 95 per cent of steady state condition.

Precision Phase Detector

Advance Electronics Co., P.O. Box 394, Passaic, N. J., has developed the Type 205 to meet the increasing need of a single instrument for detecting phase angle with error less than 0.1 degree in communication systems such as color television, and

(Continued on page 78A)

FACTS ON MAGNETIC RECORDING

the flexible "MEMORY" for science and industry

- In industry today magnetic recorders can "remember" and re-create the motions of skilled machinists, the forces encountered by a truck driving down a test road, the reflections from underground shock waves, the complex control of chemical processes.
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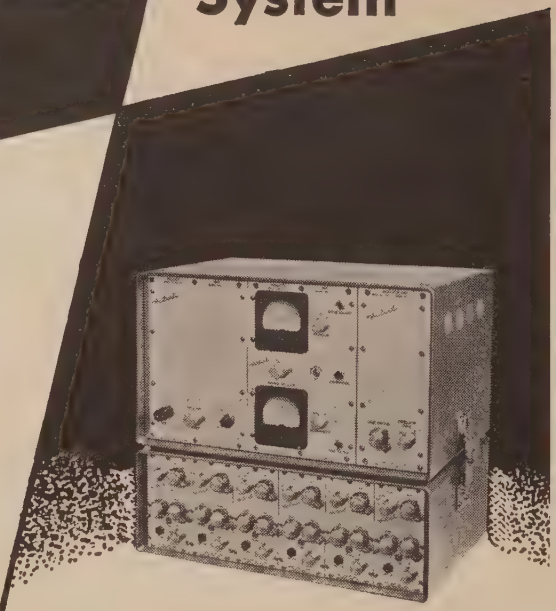
Heiland Amplifier System

The most complete, yet easiest to operate amplifier system ever developed for oscillographic recording



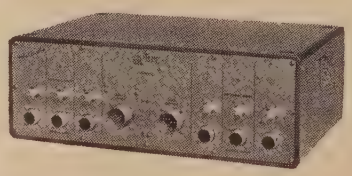
Model 119 Carrier and Linear or Integrating Amplifier System.

Heiland's model 119 Amplifier System, used in conjunction with Heiland Recording Oscillographs, has received wide acclaim from engineers for its extreme versatility, accuracy and simplicity of operation in the amplification of static and dynamic current phenomena.

This small, compact instrument, which can be provided for either rack, table, or shock mounting with available accessories, is housed in a rugged, yet lightweight cast aluminum case finished in attractive silver-gray gloss enamel. For complete specifications write or wire for our Bulletin 107.

Complete information on other Heiland products will be supplied on request.

Power Supply Assembly (Rear View)



Amplifier Assembly (Rear View)

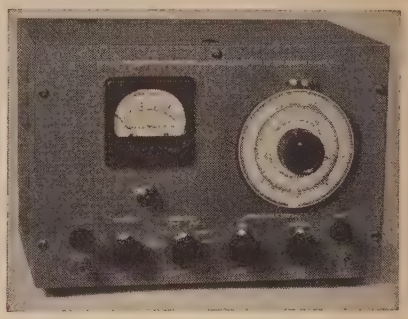
Heiland Research Corporation 130 East Fifth Ave., Denver, Colorado

News—New Products

These manufacturers have invited PROCEEDINGS readers to write for literature and further technical information. Please mention your I.R.E. affiliation.

(Continued from page 76A)

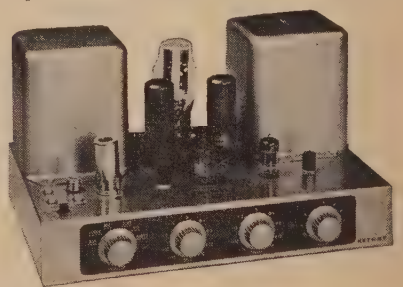
other communication networks at high frequencies. Essentially Type 205 consists of two cathode followers, a continuously variable delay line, a balanced phase detector, and a sensitive output indicator.



The accuracy is ± 0.1 degree (6 minutes) in phase reading, or ± 1 per cent of the time delay indicated on the dial of the continuously variable delay line. The frequency range is 10 kc to 10 mc. The lower limit may be extended to 1 kc with an additional delay line or phase shifting network. The upper limit may be extended with relaxing accuracy. The time delay of the two continuously variable delay lines which are supplied with the unit are 0 to 0.25 microsecond and 0 to 0.05 microsecond. Continuously variable delay lines with different time delay can be obtained upon request. The indicator sensitivity is 0.02, 0.04, shunted with 12 $\mu\mu$f on both input channels. The price is $445.00 each, F.O.B., Passaic.

24-Watt Amplifier

Allied Radio Corp., 100 N. Western Ave., Chicago 80, Ill., announces the release of a new 24-watt amplifier, the "Golden Knight." Frequency response at 24 watts output is ± 0.75 db from 20 to 40,000 cps. Harmonic distortion is less than 1 per cent at rated output; intermodulation distortion is less than 2 per cent at rated output.



The Knight amplifier has four inputs. One is for any high-impedance microphone; two inputs are for radio tuner, crystal cartridge, television, or tape recorder; another input takes General Electric, Pickering or Audak magnetic cartridges. Separate bass and treble tone controls, calibrated from -16 to $+16$ db are provided. Another feature is a three-

(Continued on page 116A)

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Julius A. Stratton

DIRECTOR, 1954

Julius A. Stratton was born on May 18, 1901 in Seattle, Washington. He received the B.S. degree in electrical engineering from the Massachusets Institute of Technology in 1923 and spent the following year in study in France. He returned to M.I.T. to become Research Assistant in the newly-established Communications Laboratory under the direction of Professors Arthur E. Kennelly and E. L. Bowles, and in 1925 he received the M.S. degree in electrical engineering. In 1927 he was awarded the D.Sc. degree in mathematical physics from the Technische Hochschule in Zurich, and he remained in Europe on an M.I.T. traveling fellowship, studying under Professor Summerfeld in Munich, until the Fall of 1928. He then became Assistant Professor in the Department of Electrical Engineering at M.I.T. In 1931 he transferred to the Department of Physics, becoming full Professor in 1941, and in 1945 he assumed Directorship of the Research Laboratory of Electronics at M.I.T. Dr. Stratton was appointed Provost in 1949, and is currently Vice President and Provost at M.I.T.

During World War II Dr. Stratton was a Staff Member of the Radiation Laboratory, Expert Consultant in the Office of the Secretary of War, and received the Medal for Merit in 1946. From 1946 to 1949 he was Chairman of the Committee on Electronics of the Joint Research and Development Board.

Dr. Stratton became a Member of the IRE in 1942, a Senior Member in 1943, and was elected a Fellow in 1945. He served as a Director from 1948–1951. He is currently on the Editorial Board.

Dr. Stratton is a member of the National Academy of Sciences, the American Physical Society, the American Academy of Arts and Sciences, Zeta Psi, Tau Beta Pi, and Sigma Xi.

The Engineering Report

FRANK G. KEAR

One important link between the engineer and his professional and commercial associates is the Engineering Report. Its significance and helpfulness are often underestimated, and its preparation is occasionally and regrettably neglected. Many readers of these PROCEEDINGS will accordingly benefit from the following helpful analysis of report preparation by a Fellow of the Institute, who is a consulting engineer and a member of Kear & Kennedy, of Washington, D. C.—*The Editor.*

The trade mark of the engineer is, by common acceptance, the slide rule, an instrument also generally assumed to be his most important tool. There is, however, a tool even more basic to an engineer, but one which is frequently neglected: this is the Engineering Report.

From the beginning of his professional education, the engineer receives training in the writing of reports. Many of these are in the nature of laboratory reports on specific tests, but a large number cover the investigation of more general problems. Familiarity with the framework of these reports is vital to clearly organized thinking, and only by means of such reports can the engineer adequately present the results of his work. Upon completion of his formal education, the engineer is often prone to overlook the continuing need for using the report framework as part and parcel of his thinking, with the result that much effort and time expended in performing work is lost to industry because of failure to prepare an adequate and complete presentation. Unless an engineer can convey to his management or his client a correct and adequate presentation of the work which he has performed, he has failed in his obligations.

Since the PROCEEDINGS OF THE I.R.E. constitutes the media whereby engineering information is disseminated within the profession, it seems appropriate that it include in its pages a reminder of the importance of the proper use of the engineering report.

The basic report outline is simple enough, having only five subdivisions—(1) the statement of the problem, (2) the facts bearing upon the problem, (3) an analysis and discussion of these facts, (4) the conclusions drawn therefrom, and (5) recommendations. Let us consider these.

(1) Formulation of a correct statement of the problem presented is obviously fundamental, yet frequently overlooked. The engineer must take ample time to determine exactly what is required of him in order that no time is wasted on issues which are not relevant. The statement should be reduced to writing before any work is begun.

(2) The facts bearing upon the problem should be carefully marshalled and great care taken to include every pertinent fact which the engineer will use. By so doing it becomes a simple matter at a later date to look back and determine the reasons for a recommended course of action. Frequently an engineer finds it advisable to change his recommendations materially because of the discovery of facts not originally considered. Reference to the original report will immediately disclose the absence of these new facts, the inclusion of which would justify the change in recommended action.

(3) The analysis and discussion of the facts will vary greatly in size and format, depending on the type of problem. It could run into several notebooks dull of pertinent information and data, or it may be a mere paragraph. If the analysis is lengthy, it should be reduced to summary form for inclusion in the report. However, brevity should not be favored at the expense of clarity of exposition.

(4) The conclusions reached after completion of the analysis should be stated concisely and clearly, pointing out the line of reasoning employed in reaching these conclusions and, of course, making certain that this reasoning is substantiated by the preceding portions of the report. This subdivision must be carefully limited to the conclusions of the engineer making the report. It should not contain any new material previously undisclosed, nor should it duplicate material which will be included in the fifth subdivision. Many engineering reports will close at this point. The average paper published in the PROCEEDINGS customarily contains these four subdivisions.

(5) The last and most important subdivision, although as in the case of a paper for the PROCEEDINGS it is not always required, is the one containing recommendations. The person or group to whom this report is made is customarily interested only in this last paragraph. It must contain clear and concise recommendations as to the proper course of action to be taken as a result of the engineering work which has been performed. Brevity is desirable. The context should not be confused by quotations from previous portions of the report. The ability to present this paragraph properly represents the dividing line between the technician and the professional engineer. When an engineer is assigned a task the presumption is that he is competent to perform this task. Accordingly, he should not waste time and space in establishing this competency. The recipient of the report merely wants to know what he should do and how he should do it. The recommendations should therefore be limited to simple, clearly-worded sentences, avoiding multiple or conflicting recommendations unless the writer clearly indicates the relative merits of each.

If the engineer follows carefully the foregoing outline in each problem presented to him, he will find his work is simplified, his reports present a permanent record of the project, and management or the client will have the necessary information on which to base their action.

Frequency Variations of Junction-Transistor Parameters*

R. L. PRITCHARD†, ASSOCIATE, IRE

Summary—A Theoretical solution for the frequency variation of the four small-signal parameters of a junction transistor has been obtained by extending Shockley's analysis and by taking account of space-charge-layer widening as suggested by J. M. Early. From the results of this solution, it has been possible to explain the experimentally observed frequency variation of open-circuit collector-base admittance of fused-junction *p-n-p* junction transistors that was reported earlier and is described here in more detail. The theoretical frequency variation of the current-amplification factor and of the other two small-signal parameters (open-circuit voltage-feedback factor h_{12} and short-circuit input impedance h_{11}, for the grounded-base connection) also is discussed in some detail. Numerical results are included for each of the four parameters in the form of curves of normalized parameters versus relative frequency. Derivation of the voltage-current relations for the theoretical model is given separately in section II. The effect of base-spreading resistance is taken into account in the usual manner by addition of a lumped resistance to the base contact of the theoretical model.

Introduction

FROM AN electric-circuit point of view, a transistor may be considered as an active four-terminal network, the behavior of which may be described by four parameters. One of the most commonly used sets of parameters for the junction transistor comprises the elements of an equivalent-Tee network[1] as shown in Fig. 1. The four parameters in this representation are the emitter resistance r_e, the base resistance r_b, the parallel combination of collector resistance r_c and collector capacitance C_c, and the effective current-amplification factor α_e. For small-signal linear operation of the transistor, the four parameters are considered as constants determined by the dc operating point.

The frequency variation of the current-amplification factor has been discussed in some detail.[2] On the other hand, it had been a popularly accepted fact that the other three parameters were independent of frequency.[3] However, on the basis of a theoretical study of the junction transistor, J. M. Early pointed out that both r_e and r_b also may vary with frequency. In fact, there is both a high-frequency and a low-frequency value for each of these parameters. Further frequency effects were not considered by Early.[4]

During the initial stages of an investigation of the electric-circuit properties of fused-junction *p-n-p* junction transistors developed by Saby[5] it was found that collector resistance r_c decreased with increasing frequency. Decreases of the order of a factor of 10 from the low-frequency value were observed. Subsequently, S. K. Ghandhi[6] found that the collector capacitance also decreased with increasing frequency.

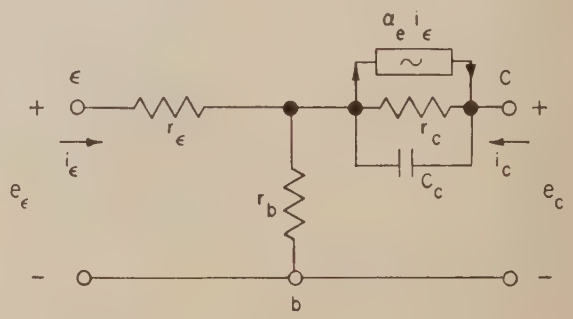

Fig. 1—Equivalent-tee circuit for junction transistor.

By extending Early's analysis to include frequency effects, it has been possible to explain the observed behavior of the collector impedance with frequency. Frequency variations of the other three parameters also have been calculated. The results of the theoretical analysis for the frequency variations of the four parameters are presented in the first section of the paper together with experimental results for the collector-base impedance. Calculation of the voltage-current relations for the theoretical model of the transistor from physics of the structure is given separately in section II.

I. Frequency Variation of Small-Signal Parameters

There are a large number of small-signal parameters that may be calculated for a junction transistor, e.g., either grounded-base, grounded-emitter, or grounded-collector operation may be considered, and six sets of parameters may be calculated for each connection. However, if any four *independent* parameters are known,

* Decimal classification: R282.12. Original manuscript received by the IRE, September 18, 1953. Presented at the Semiconductor Symposium of the Electrochemical Soc. Meeting, New York, N. Y., April 15, 1953, and at the Transistor Research Conference, State College, Pa., July 6, 1953. The research was supported by the Government Tri-Services Contract AF 33(600) 17793, Semiconductors.

† Communications Research Section, Genl. Elec. Research Lab., Schenectady, N. Y.

[1] See, for example, R. L. Wallace and W. J. Pietenpol, "Some circuit properties and applications of *n-p-n* transistors," PROC. I.R.E., vol. 39, Fig. 26, p. 765; July, 1951. Also *Bell Sys. Tech. Jour.*, vol. 30, p. 558; July, 1951. The voltage generator shown in their Fig. 26 can be replaced in a straightforward manner by the current generator shown in Fig. 1 above.

[2] See: a) W. Shockley, M. Sparks, and G. K. Teal, "*p-n* junction transistors," *Phys. Rev.*, vol. 83, section IX, p. 161; July 1, 1951, b) R. L. Pritchard. "Frequency variations of current-amplification factor for junction transistors," PROC. I.R.E., vol. 40, pp. 1476–1481; November, 1952, c) D. E. Thomas, "Transistor amplifier—cut-off frequency," *ibid.*, pp. 1481–1483.

[3] This is based on experience with point-contact transistors for which the three resistive parameters of the Tee network were found to be essentially independent of frequency up to several mc. See J. A. Becker and J. N. Shive, "The transistor—a new semiconductor amplifier," *Elec. Eng.*, vol. 68, p. 221; March, 1949, or J. Bardeen and W. H. Brattain, "Physical principles involved in transistor action," *Bell Sys. Tech. Jour.*, vol. 28, p. 258; April, 1949; also in *Phys. Rev.*, vol. 75, pp. 1215–1216; April 15, 1949.

[4] a) J. M. Early, "Effects of space-charge layer widening in junction transistors," PROC. I.R.E., vol. 40, pp. 1401–1406; November, 1952.
b) Since submission of this paper for publication Early has written, "Design theory of junction transistors," *Bell Sys. Tech. Jour.*, vol. 32, pp. 1271–1312; November, 1953.

[5] J. S. Saby, "Recent developments in transistors and related devices," *Tele-Tech*, vol. 10, p. 34; December, 1951. See also, J. S. Saby, "Fused impurity *p-n-p* junction transistors," PROC. I.R.E., vol. 40, pp. 1358–1360; November, 1952.

[6] S. K. Ghandhi, Electronics Laboratory, Syracuse, N. Y., personal communications, March, 1952.

any of the other parameters may be calculated. From the point of view of simplicity of presentation, this writer prefers to use the series-parallel, or $[h]$, parameters for the grounded-base connection of the transistor. These are defined by the equations[7]

$$\left.\begin{aligned} e_1 &= h_{11}i_1 + h_{12}e_2, \\ i_2 &= h_{21}i_1 + h_{22}e_2. \end{aligned}\right\} \quad (1)$$

There also are other advantages of using these parameters, e.g., in general they are the most easily measured quantitites for the junction transistor, which has a low input impedance and a high output impedance.

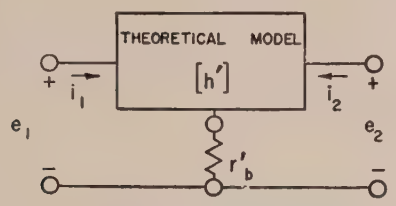

Fig. 2—Effect of adding base spreading resistance to theoretical model for junction transistor.

The current-voltage relations for the theoretical model of the transistor are obtained in the form $[i]=[y]\cdot[e]$, i.e., the admittance parameters are calculated. However, when the effect of the resistance of the base region is taken into account in the usual manner by adding an external resistance in series with the base contact to the theoretical model, the $[y]$ representation becomes somewhat more complicated. Accordingly, the four $[h]$ parameters for the theoretical transistor have been calculated from the $[y]$ parameters and are given in (59)–(63) of section II. In this section the parameters of the theoretical model will be characterized by primes, e.g., h_{11}'. When the base spreading resistance r_b' is added to the theoretical model as shown in Fig. 2, the $[h]$ parameters of the over-all transistor are given by the matrix equation (2).

$$\begin{bmatrix} e_1 \\ i_2 \end{bmatrix} = \begin{bmatrix} h_{11} & h_{12} \\ h_{21} & h_{22} \end{bmatrix} \times \begin{bmatrix} i_1 \\ e_2 \end{bmatrix}$$

$$[h] = \begin{bmatrix} h_{11}' + \dfrac{r_b'(1+h_{21}')(1-h_{12}')}{(1+r_b'h_{22}')} & \dfrac{h_{12}'+r_b'h_{22}'}{(1+r_b'h_{22}')} \\ \dfrac{h_{21}'-r_b'h_{22}'}{(1+r_b'h_{22}')} & \dfrac{h_{22}'}{(1+r_b'h_{22}')} \end{bmatrix} \quad (2)$$

It turns out that the output-admittance parameter h_{22} is least affected by the addition of r_b', whereas the feedback parameter h_{12} is most affected. Next most affected is the input impedance h_{11}, and third most affected is the current-amplification factor $h_{21}(=-\alpha)$. Each of the four parameters will be discussed in detail below.

A. Current-Amplification factor h_{21}

1. Theoretical model:

$$-h_{21}'=\alpha'=\gamma\cdot\beta. \quad (3)$$

For a junction transistor with a simple p-n junction collector, current-amplification factor α' may be written[2a] as the product of two factors γ and β. Emitter efficiency γ is the ratio of minority-carrier current (hole current for p-n-p transistor) into the base from the emitter junction to the total emitter current. Thus, the departure of γ from unity is a measure of undesired current from the base to the emitter (electron current for a p-n-p transistor), plus the reactive current through the capacity of the emitter-base junction. The factor β is that fraction of the minority-carrier current from the emitter that reaches the collector.

The frequency variation of β has been discussed[2a,b] previously and will be reviewed here only briefly. If the low-frequency value β_0 is unity, the frequency variation of β may be described by the expression

$$\beta = \operatorname{sech}(j\omega\tau_D)^{1/2}$$

where τ_D is a diffusion time for carriers in the base, i.e., a built-in transit time. The β-cut-off frequency $\omega_c/2\pi$ is defined as the frequency for which $|\beta|^2$ has decreased to $\frac{1}{2}$. This occurs for $\omega\tau_D=2.43$. In terms of the thickness w expressed in mils (10^{-3} inch) of the base region of the transistor, ω_c may be calculated from the expression[2b]

$$(\omega_c/2\pi)=\begin{cases} 5.6\times10^6/w^2 \text{ for an } n\text{-}p\text{-}n \text{ Germanium unit} \\ 2.6\times10^6/w^2 \text{ for a } p\text{-}n\text{-}p \text{ Germanium unit.} \end{cases} \quad (4)$$

If β_0 is different from unity, the term $j\omega\tau_D$ must be replaced by a complex variable. However, an excellent approximation for β in this case is simply

$$\beta = \beta_0 \operatorname{sech}(j\omega\tau_D)^{1/2}, \quad (5)$$

with

$$\tau_D \equiv 2.43/\omega_c, \quad (6)$$

where the β-cut-off frequency $\omega_c/2\pi$ *is defined as the frequency for which* $|\beta|^2/\beta_0^2$ *has decreased to* $\frac{1}{2}$. Exact values of β/β_0 as a function of ω/ω_c have been calculated and have been found to be in excellent agreement with respect to both phase and amplitude with values calculated from the approximate equation (5) and (6). However, in general, the relation between ω_c and w is a function of β_0. (This would be of consequence only if ω_c were to be calculated accurately from a measured value of w.) Numerical values of ω_c as a function of w and β_0 are given in Appendix A.

On the other hand, the frequency variation of the emitter efficiency γ cannot be described in such simple terms. In general, three parameters plus the parameter τ_D (or ω_c) introduced above for β, are necessary to define γ at all frequencies. A useful approximate expression for γ may be written as

$$\gamma = \{1 + [A(1+j\omega\tau_e)^{1/2} + (j\omega C_e'r_e'/\gamma_0)] \cdot [\tanh(j\omega\tau_D)^{1/2}/(j\omega\tau_D)^{1/2}]\}^{-1}. \quad (7)$$

Here A is a parameter involving relative concentrations of minority carriers in base and emitter regions, base thickness w, and the parameter τ_e, which is a second

[7] See, for example, E. A. Guillemin, "Communication Networks," vol. II, chap. IV, John Wiley & Sons, New York, N. Y., 1935. Also J. S. Brown and F. D. Bennett, "The application of matrices to vacuum-tube circuits," Proc. I.R.E., vol. 36, p. 845; July, 1948. Note however that Guillemin considers the special case of a bilateral network; hence most of his results are not directly applicable here.

characteristic time constant of the transistor (lifetime of minority carriers in the emitter). Alternatively, A may be defined in terms of the low-frequency value γ_0 of γ as

$$A \equiv (1 - \gamma_0)/\gamma_0. \tag{8}$$

In (7) C_e' is the capacity of the emitter-base junction, and r_e' is the low-frequency short-circuit input impedance to be discussed below ($r_e' \doteq kT/q_e I_e$ as given by Shockley et al.[2a]). Finally, the hyperbolic-tangent function is an approximation for a more complicated expression just as β was approximated above.

For the moment, it may be assumed that the effect of $\omega C_e' r_e'$ is negligible, but that $A \neq 0$, i.e., $\gamma_0 \neq 1$. As the frequency is increased, the hyperbolic function remains substantially constant, but the diffusion-admittance term $(1+j\omega\tau_e)^{1/2}$ discussed by Shockley[8] increases, causing a decrease in $|\gamma|$. If ω_c were infinitely high (τ_D infinitely small), ultimately $|\gamma|$ would decrease as $\omega^{-1/2}$, and the phase angle of γ would approach -45 degrees. This is shown by the curves of Fig. 3 for several values of A. However, with finite ω_c, the hyperbolic-function term within the brackets in (7) decreases in magnitude and becomes complex with increasing frequency. This tends to counteract the increase in magnitude of the term $(1+j\omega\tau_e)^{1/2}$, and $|\gamma|$ decreases less rapidly than $\omega^{-1/2}$. Ultimately, the hyperbolic-function term decreases with frequency at the same rate as $(1+j\omega\tau_e)^{1/2}$ increases, so that the product becomes constant. Hence, $|\gamma|$ attains a limiting high-frequency value that is determined by the ratio τ_e/τ_D and by the parameter A. (Actually, the limiting value is determined only by the ratio of minority-carrier concentrations in emitter and base respectively and by the ratio of their mobilities, which is a constant.) This is shown by the curves of Fig. 4 for several ratios of τ_D/τ_e for a particular value of A corresponding to $\gamma_0 = 0.93$. Note that the high-frequency value of $|\gamma|$ may be higher or lower than γ_0 in general.

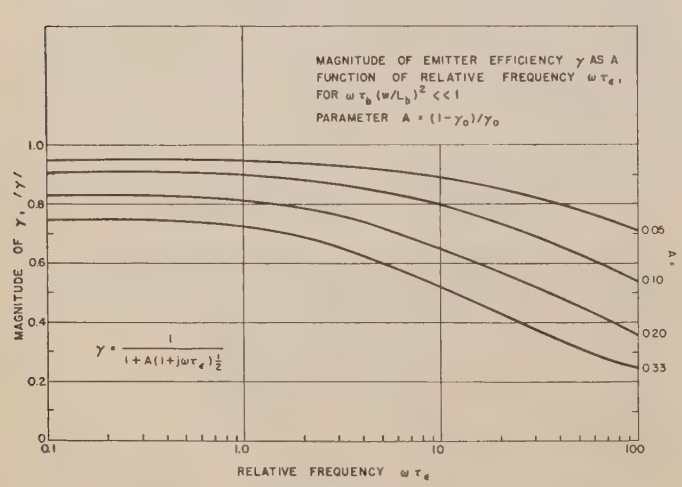

Fig. 3—Theoretical variation of the magnitude of the emitter efficiency γ with relative frequency, for case of $\omega\tau_D \ll 1$.

[8] W. Shockley, "The theory of p-n junctions in semiconductors and p-n junction transistors," *Bell Sys. Tech. Jour.*, vol. 28, pp. 459–465; July, 1949. Also "Electrons and Holes in Semiconductors," pp. 313–318; D. Van Nostrand, New York, N. Y., 1950.

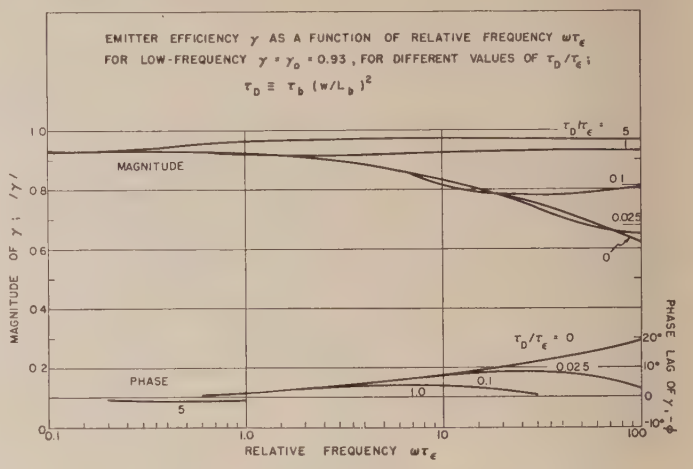

Fig. 4—Theoretical variation of emitter efficiency γ with relative frequency, for $\gamma_0 = 0.93$.

The effect of both γ and β in determining the overall frequency behavior of α is illustrated by the curves of Fig. 5 for a particular case. Numerical values of the parameters were selected to provide a fit with experimental results of α versus frequency obtained for an experimental grown-junction transistor.[9] Note that γ alone determines the low-frequency behavior of $|\alpha'|$, resulting in a gradual fall-off of $|\alpha'|$, whereas the effect of β becomes significant at higher frequencies. This "slow-drool" type of behavior of $|\alpha'|$ at low frequencies appears to be typical of the effect of γ upon α'.

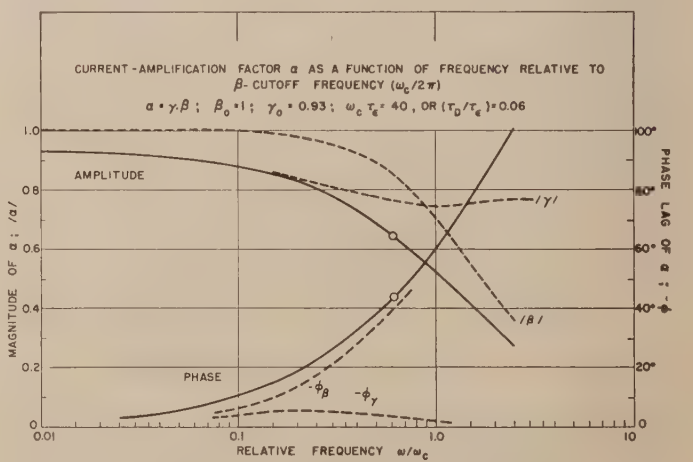

Fig. 5—Theoretical variation of current-amplification factor α with relative frequency for $\gamma_0 = 0.93$, for ratio of diffusion time τ_D to emitter lifetime τ_e of 0.06.

Note also that the cut-off frequency of α', indicated by the circles, has been reduced relative to the β-cut-off frequency $\omega_c/2\pi$. Furthermore, the magnitude of the phase angle at α'-cut-off frequency is substantially less than the 57 degrees which occurs for β-type cut-off only.

[9] This transistor was fabricated by R. N. Hall in the course of his experimental work with crystal growing. The β-cut-off frequency $\omega_c/2\pi$ was of the order of 600 kc.

Suppose now that $A=0$, i.e., $\gamma_0=1$, in order to consider the relatively simple effect of C_e'. The capacity of the emitter-base junction may be fairly large because of the small-voltage drop across this junction (of the order of a few tenths of a volt); hence, C_e' could be of the order of 50 $\mu\mu$f for a junction for which reverse-bias capacitance was of the order of 10 $\mu\mu$f. Therefore, for low values of emitter current for which r_e' may be fairly high (of the order of 10^3 ohms for $I_e=25$ μa) the product $\omega C_e' r_e'$ may be appreciable at operating frequencies. In this case, the reactive ac current through the emitter-base junction causes an increase in emitter current at increasing frequency without producing a corresponding increase in collector current; hence γ will decrease in magnitude, ultimately as $1/\omega$. However, at frequencies approaching the β-cut-off frequency $\omega_c/2\pi$, the forward-bias impedance of the emitter-base junction decreases as the square root of increasing frequency. This tends to reduce the bypassing effect of C_e' in a manner decreased mathematically by a hyperbolic-tangent term in (7). Hence, at the higher frequencies $|\gamma|$ will decrease less rapidly. A typical calculated variation of γ with frequency due to the effect of C_e' is shown in Fig. 6 for several values of $\omega_c C_e' r_e'$. For the (relatively large) value of $\omega_c C_e' r_e'=1$, the decreased rate of fall-off of $|\gamma|$ with frequency that occurs at the higher frequencies is quite evident.

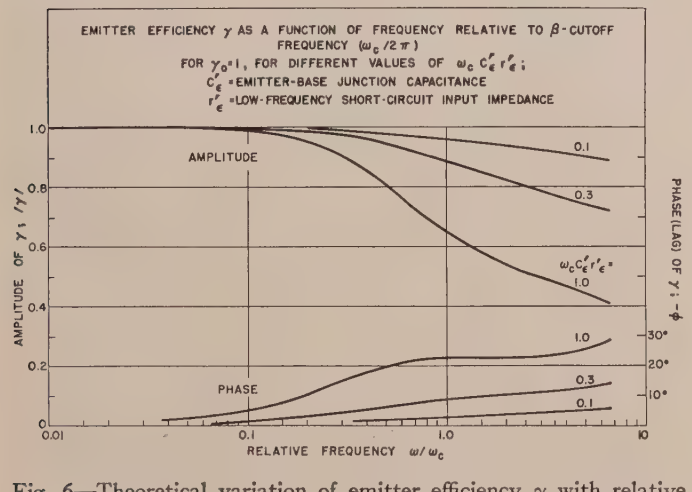

Fig. 6—Theoretical variation of emitter efficiency γ with relative frequency illustrating effect of emitter junction capacitance, with $\gamma_0=1.0$.

2. Modification due to r_b':

$$-h_{21}=\alpha=(\alpha'+r_b'h_{22}')/(1+r_b'h_{22}'). \qquad (9)$$

At high frequencies, the open-circuit output admittance h_{22}' of the theoretical transistor is essentially capacitive reactive and hence increases in magnitude with increasing frequency. If the product $r_b'h_{22}'$ is appreciable (of the order of 0.1 or more) at the β-cut-off frequency $\omega_c/2\pi$, the frequency variation of α for the over-all transistor may be significantly different from that of the theoretical model. In particular, $r_b'h_{22}'$ has a positive reactive component which subtracts from the negative reactive component of α' to reduce the magnitude of α relative to that of α'. Furthermore, the rate of change of $|\alpha|$ with frequency generally is increased.[10]

A typical variation of α with frequency of this type is shown in Fig. 7 for a particular case, in which it is assumed that $\gamma=1$, and that $h_{22}'=j\omega C_{22}'$ only, i.e., output conductance is neglected. The base resistance r_b' is assumed to be such that at the β-cut-off frequency

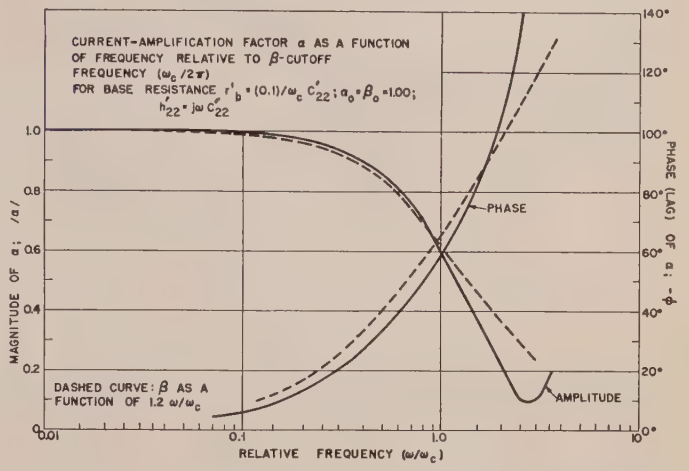

Fig. 7—Theoretical variation of current-amplification factor α with relative frequency illustrating effect of base spreading resistance, for $\alpha_0=\beta_0=1.0$, for negligible collector-base conductance.

$\omega_c r_b' C_{22}'=0.1$ (e.g., $\omega_c/2\pi=5\times10^6$, $r_b'=400$, $C_{22}'=8$ $\mu\mu$f). In order to show the rate of decrease of $|\alpha|$ relative to that of $|\beta|$, the dashed curve of β in Fig. 7 has been plotted to a slightly different abscissa scale. Note that the cut-off frequency of α is reduced approximately 20 per cent relative to that of $\alpha'=\beta$ because of r_b', and that the rate of cut-off of $|\alpha|$ is more rapid than the cut-off of β. Such more-rapid-than-theoretical decreases in α with frequency have been observed in practice with transistors having relatively large values of base spreading resistance.[11]

B. Open-Circuit Collector-Base Admittance h_{22}

1. *Theoretical model:* The exact expression for h_{22} which is given in section II is quite complicated for

[10] From a circuit point of view, this behavior may be explained as follows. Total impressed voltage in the collector-base mesh is the sum of an $i_e r_b'$ drop in the base resistance plus the voltage appearing across the collector capacitance due to injected emitter current (cf. the current generator of the equivalent circuit in Fig. 1). For a constant ac emitter current, the first term is constant with frequency, whereas, the voltage across the capacitance decreases with increasing frequency because of the decrease in the magnitude of both α and of the capacitive reactance. Of considerably more importance, however, is the fact that the voltage across the collector capacitance lags the $i_e r_b'$ drop by more then 90 degrees because of the phase lag of α. Hence, these two impressed voltages tend to cancel each other, and the voltage available in the collector-base mesh decreases, causing α to fall off with frequency faster than if r_b' were zero. At very high frequencies $i_e r_b'$ voltage predominates, and α actually increases toward unity (but there is essentially no power gain at these frequencies).

[11] Note, for example, the sharper α-cut-off of the junction transistor tetrode when operated as a triode with high base resistance relative to the tetrode operation in which base resistance is reduced. See R. L. Wallace, Jr., L. G. Schimpf, and E. Dickten, "A junction transistor tetrode for high frequency use," PROC. I.R.E., vol. 40, Fig. 7, p. 1397; November, 1952. However, this effect of r_b' on α described above should not be confused with the effect of r_b' on power gain described by Wallace, et al.

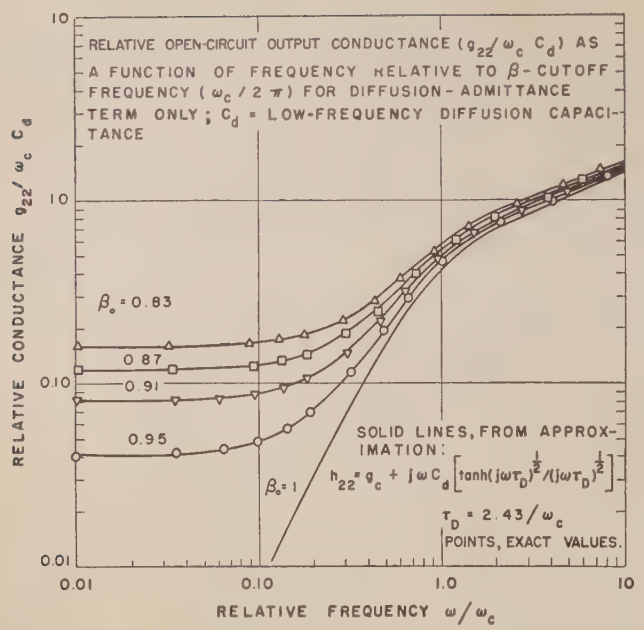

Fig. 8—Theoretical variation of relative open-circuit output conductance with relative frequency, for diffusion-admittance term only.

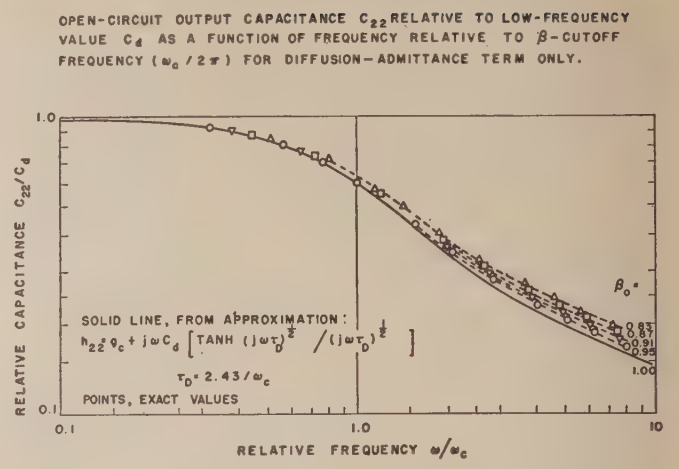

Fig. 9—Theoretical variation of relative open-circuit output capacitance with relative frequency, for diffusion admittance term only.

purposes of calculation. The result may be simplified somewhat by expanding the hyperbolic-tangent in a Taylor Series and retaining only the first few terms, in order to make use of the function of the semi-imaginary argument $(j\omega\tau_D)^{1/2}$ as was done in the case of the current-amplification factor β. The result may be written in the form

$$h_{22}' \doteq j\omega C_c' + g_{c\beta} + j\omega C_d[\tanh (j\omega\tau_D)^{1/2}/(j\omega\tau_D)^{1/2}] + g_{c\gamma}(\gamma/\gamma_0)\cdot\beta^2\cdot(1 + j\omega\tau_\epsilon)^{1/2}, \quad (10)$$

where C_c' is the capacity of the collector-base junction, $g_{c\beta}$ and $g_{c\gamma}$ are two parts of the expression for collector conductance calculated by Early due to the effect of space-charge-layer widening, and C_d is a diffusion capacity due to charge stored in the base layer by the dc emitter current. (See Appendix B for more details of C_d.) In particular[12]

$$g_{c\beta} \doteq 2\alpha_0(I_\epsilon/w_0)(\partial w/\partial E_c)(1 - \beta_0)$$
$$g_{c\gamma} \doteq \alpha_0(I_\epsilon/w_0)(\partial w/\partial E_c)(1 - \gamma_0). \quad (11)$$

where I_ϵ is the dc emitter current, w_0 is the thickness of the base region of the transistor, and $(\partial w/\partial E_c)$ is rate of change of base thickness with dc collector voltage.

In (10), the first and fourth terms are substantially exact, whereas the second and third terms result from the approximation mentioned above.[13] However, numerical values of these terms have been calculated and have been compared with values calculated from the corresponding exact expression. Results of the comparison of real and imaginary parts of $(h_{22}' - j\omega C_c')$ are shown in normalized form in Figs. 8 and 9 respectively for values of $\beta_0 = 0.95, 0.91, 0.87, 0.83$. (Fig. 9 actually shows $1/j\omega$ times the imaginary part, i.e., the effective capacity.) For the smaller values of β_0 the approximate values may be in error by ten to fifteen per cent, whereas the error is less for the larger values of β_0.

When γ is essentially unity, as in the case of fused-junction transistors,[14] the fourth term of (10) vanishes, and h_{22}' reduces to the form described earlier.[15] In this case, the diffusion capacitance C_d may be determined from the difference between the value of the low-frequency collector-base capacitance corresponding to a given emitter current and that for zero-emitter current. For an abrupt, or step, junction, as in fused-junction transistors, theory given in Appendix B indicates that C_d is directly proportional to dc emitter current and inversely proportional to square root of collector voltage. Experimental values of collector-base capacitance as a function of dc operating bias, which are shown in Fig. 10, are in reasonably good agreement with these theoretical results. Note, however, the departure of linearity with respect to I_ϵ at the larger currents. In this case, the applied current is sufficiently great that the number of injected (minority) carriers is comparable with the number of (majority) carriers normally present in the base. Thus, the condition of low-level carrier injection implicitly assumed in the theoretical analysis is violated, and the theoretical results can no longer be expected to agree perfectly with experiment.

As the frequency is increased, the term involving C_d becomes complex, and a conductance term results from this diffusion capacitance. This conductance term increases as the square of the frequency at first; hence, the conductance g_c normally predominates at low frequencies. At higher frequencies, the conductance due

[12] See Early (footnote 4), p. 1403; the expressions given here for g_c are approximately correct for $\alpha > 0.95$. For $\alpha < 0.95$, the more exact results given by Early (p. 1405) should be used.

[13] More accurately, the term $g_{c\beta}$ should read $(g_{c\beta}/2)\{[\tanh (j\omega\tau_D)^{1/2}/(j\omega\tau_D)^{1/2}] + \beta/\beta_0$. However, the frequency variation of this term generally is relatively unimportant compared to that of the C_d term above, and (10) as given is sufficiently accurate for most practical purposes.

[14] E. L. Steele, "Theory of alpha for p-n-p diffused junction transistor," PROC. I.R.E., vol. 40, p. 1427; November, 1952.

[15] R. L. Pritchard, "Collector-Base Impedance of a Junction Transistor," PROC. I.R.E., vol. 41, p. 1060; August, 1953.

to C_d predominates, and increases ultimately as the square root of the frequency. This is shown by the curves of Fig. 8. The curve for $\beta_0=1$ can be used with any low-frequency value of g_c to calculate g_{22} versus frequency according to the approximate equation (10).

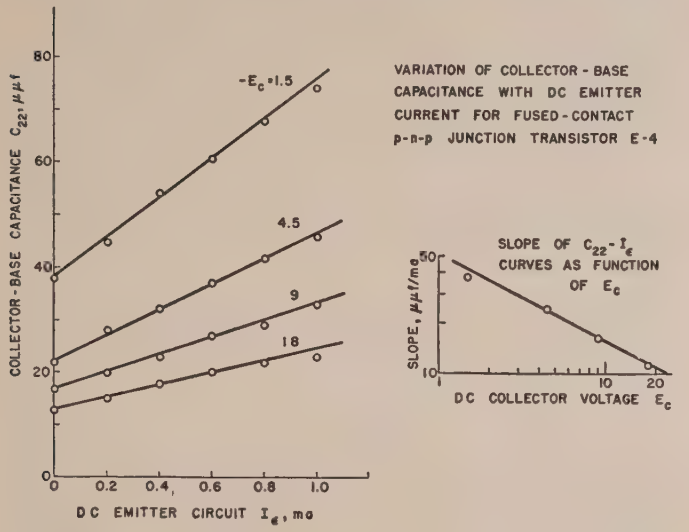

Fig. 10—Variation of collector-base capacitance with dc bias for *p-n-p* fused-junction transistor.

On the other hand, the capacitive part of the term involving C_d decreases with increasing frequency, ultimately as the square root of the frequency, as shown in Fig. 9. Hence, the collector-base capacitance decreases from its low-frequency value of $(C_c'+C_d)$ to an ultimate high-frequency value of C_c'.

It is especially important to note the following: If C_d is small *relative* to C_c', the frequency variation of C_{22} may be relatively unimportant; however, if C_d is appreciable in *magnitude* (e.g., several $\mu\mu$f), the high-frequency collector-base conductance of the transistor may be quite large, causing an over-all decrease in the effective Q of the device.

Experimental results for the frequency variation of h_{22} are shown in Fig. 11 for an early model fused-junction *p-n-p* transistor developed by Saby.[5] The solid curves have been calculated from the approximate equation (10) using as data only the low-frequency values of C_c, $C_d \equiv C_{22}-C_c$, g_{22}, and the α-cut-off frequency f_c. Agreement between experimental and theoretical results is surprisingly good.

When $\gamma_0 \neq 1$, in addition to the admittance terms discussed above, h_{22}' contains an additional diffusion admittance term of the type $(1+j\omega\tau)^{1/2}$ discussed by Shockley,[8] multiplied by the frequency-sensitive factors γ and β^2. Because of the rapid decrease of $|\beta|^2$ with increasing frequency, the effect of this extra term becomes less important at the higher frequencies.

2. Modification due to r_b':

$$h_{22} = h_{22}'/(1 + r_b'h_{22}'). \tag{12}$$

When $r_b'h'_{22}$ is appreciable, h_{22} becomes slightly different from h_{22}' as decribed above. The principal effect of r_b' on h_{22} is to increase the high-frequency conductance, and to a lesser extent to reduce high-frequency capacitance. Thus, to a first approximation, if $h_{22}' \approx j\omega C_{22}'$, $h_{22} \approx (\omega^2 r_b' C_{22}'^2 + j\omega C_{22}')/[1+(\omega r_b' C_{22}')^2]$. In some transistors, with very small values of C_d, the conductance due to r_b' may be the dominant factor in the conductive component of h_{22} at high frequencies.

C. Short-Circuit Input Impedance h_{11}'

1. Theoretical model:

$$h_{11}' \doteq r_e'(\gamma/\gamma_0)\left[\tanh(j\omega\tau_D)^{1/2}/(j\omega\tau_D)^{1/2}\right] \tag{13}$$

Alternatively h_{11}' may be represented as the parallel combination of 3 admittances:

$$1/h_{11}' \doteq (\gamma_0/r_e')\left[(j\omega\tau_D)^{1/2}\coth(j\omega\tau_D)^{1/2}\right] + \left[\frac{1-\gamma_0}{r_e'}\right](1+j\omega\tau_e)^{1/2} + j\omega C_e'. \tag{14}$$

As before, the hyperbolic-tangent function is an approximation that is accurate to within 5–10 per cent for $\beta_0>0.9$. The emitter resistance r_e' is the quantity calculated by Shockley, et al.,[2] as $r_e' = (kT/q_eI_e)$; actually I_e should be replaced by the sum of the input emitter current plus a small emitter+saturation current (the current that flows between emitter and base when both collector and emitter are biased in the reverse direction.)

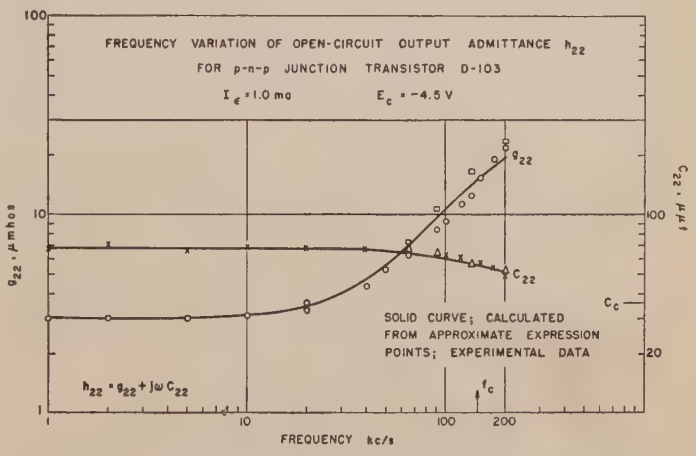

Fig. 11—Variation of open-circuit output admittance with frequency for *p-n-p* fused-junction transistor.

Hence, r_e' has a limiting value for very small emitter currents. It should be emphasized that r_e' is *not* the usual emitter resistance of the equivalent Tee circuit. Early has shown that the Tee emitter resistance may be smaller than r_e' by a factor of 2 or more.

Consider first that $\gamma=1$. As the frequency is increased, the term within the brackets of (13) decreases in magnitude and has an associated phase shift. The real and imaginary parts of the *reciprocal* of the expression within the brackets of (13) (which is the same as the first term in brackets in (14)) are shown in Fig. 12; thus, Fig. 12 shows the relative input *admittance* of the

the theoretical model. Note that for frequencies up to approximately the β-cut-off frequency the (relative) conductance is essentially constant, while the susceptance behaves as that of a capacity having an approximate value of

$$C_e'' = 0.81/\omega_c r_e'. \tag{15}$$

In general, this capacity will be quite large (e.g., $r_e' = 25$ ohms, $\omega_c/2\pi = 5\times10^6$, $C_e'' \doteq 0.001$ μf). At frequencies greater than the β-cut-off frequency, the real and imaginary parts of the admittance tend to become equal, and to vary directly as the square root of the frequency. Thus, the over-all frequency behavior of the first term of (14) is very similar to that of the conventional diffusion admittance[8] (except for the absolute frequency scale, which is proportional in this case to diffusion time through the base rather than to lifetime).

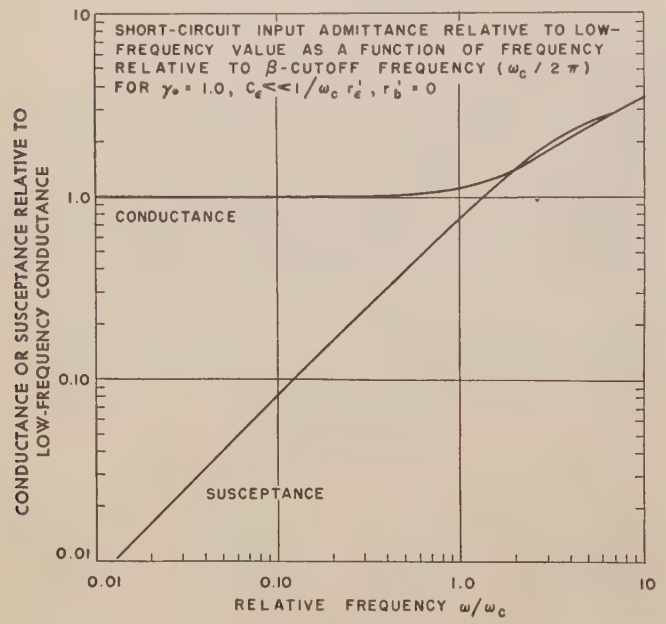

Fig. 12—Theoretical variation of relative short-circuit input admittance with relative frequency for theoretical model, for $\gamma_0 = 1$; effect of emitter-junction capacitance not included.

When $\gamma_0 \neq 1$, a conventional diffusion admittance, the magnitude of which is proportional to $(1-\gamma_0)/r_e'$, must be added in parallel with the admittance term described above. In addition, the short-circuit input admittance includes the effect of emitter-base junction capacitance C_e'. This will be important only if C_e'' is quite small, as in the case of low emitter currents.

2. Modification due to r_b':

$$h_{11} = h_{11}' + [r_b'(1-\alpha')(1-h_{12}')/(1+r_b'h_{22}')]. \tag{16}$$

In general the term $(1-h_{12}')$ can be approximated by unity at all frequencies.

The second term in this equation will increase with frequency as the quadrature-component of α' becomes significant. This will occur at frequencies of the order of $(1-\alpha_0)$ times the α-cut-off frequency.[16] If r_b' can be considered as a real constant independent of frequency (which may not be valid for some transistors), this second term should continue to increase with frequency until the limiting value of r_b' is reached, or until $r_b'h_{22}'$ becomes appreciable. Furthermore, the reactive part of the bracketed term will be inductive for frequencies up to approximately $5\omega_c$, provided $r_b'h_{22}'$ is negligible. For values of emitter current $I_e = 1$ma., or greater, the h_{11}' term generally will be negligible at the higher frequencies. Hence, the input impedance of the over-all transistor is determined essentially by the base spreading resistance multiplied by $(1-\alpha)$ and should be inductive.[17]

D. Open-Circuit Voltage Feedback Parameter h_{12}

1. Theoretical model:

$$h_{12}' = \mu_{ec}(\alpha'/\alpha_0). \tag{17}$$

At low frequencies, h_{12} becomes identical with the μ_{ec} parameter introduced by Early.[4] As the frequency is increased, h_{12}' varies with frequency in exactly the same manner as does α'. However, generally this is of no significance because of the modification introduced by r_b'.

2. Modification due to r_b':

$$h_{12} = (h_{12}' + r_b'h_{22}')/(1 + r_b'h_{22}'). \tag{18}$$

As frequency is increased from very low values, h_{22}' becomes essentially capacitive reactive (this may occur between 500 and 50,000 cps, for example). Hence, the term $r_b'h_{22}'$ will increase linearly with frequency. The magnitude of this term generally is such that an increase in h_{12} relative to its low-frequency value of $(\mu_{ec}+g_{22}r_b')$ may be noted at frequencies of the order 5–50 kc. (e.g. with $r_b' = 400$ ohms, $C_{22} \approx 8$ $\mu\mu$f, $(\mu_{ec}+g_{22}r_b') \approx 3\times10^{-4}$, $|h_{12}|$ would increase to $\sqrt{2}$ times its low-frequency value at a frequency of 15 kc.). Beyond this frequency range, $|h_{12}|$ increases linearly with frequency, up to frequencies for which $r_b'h_{22}'$ becomes appreciable relative to unity.

E. Equivalent Circuit

The frequency variation of the junction-transistor parameters also may be described by means of an equivalent circuit, if desired. A substantially exact circuit is shown in section II (Fig. 16), and an approximate circuit for the case of $\gamma \equiv 1$, applicable for the fused junction transistor, has been given previously.[15]

For the more general case, an approximate circuit is shown in Fig. 13. This circuit is based on the approximations used above for the parameters h_{ij}'. Thus, dis-

[16] The frequency behavior of $(1-\alpha)$ is very similar to that of the grounded-emitter current-amplification factor $\alpha/(1-\alpha)$ discussed earlier by the writer; see footnote 2b.

[17] The inductive input reactance of the junction transistor at high frequencies was pointed out by Wallace et al., (footnote 11), p. 1400.

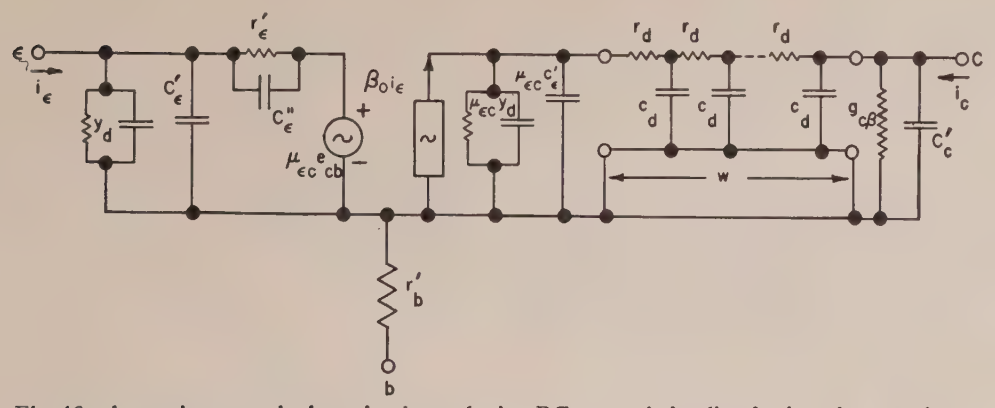

Fig. 13—Approximate equivalent circuit employing RC transmission line for junction transistor.

tributed shunt conductance in the collector-circuit transmission line has been replaced by a lumped conductance $g_{c\beta}$. Also, the transmission line in the emitter circuit, which is indicated by the form of h_{11}', has been replaced by a lumped resistance and capacitance in parallel, in accordance with the frequency variation of h_{11}' noted above (cf. Fig. 12). Note that the effect of nonunity γ is indicated by the presence of the diffusion admittance y_d (of the conventional form $(1+j\omega\tau)^{1/2}$) and the emitter junction capacitance C_ϵ' in the emitter circuit *and* at the input to the collector-circuit transmission line. When y_d and C_ϵ' can be neglected, the circuit of Fig. 13 reduces to that given previously,[15] with $z_b \equiv r_b'$ and $z_e \equiv r_\epsilon'/(1+j\omega r_\epsilon' C_\epsilon'')$.

F. Summary

In general, the frequency variation of the junction-transistor parameters cannot be described in simple terms. However, for the special and important case for which the emitter efficiency γ can be assumed to equal unity at all frequencies, some simplification is possible. This is normally the case in fused-junction transistors.

It has been shown that the current-amplification factor varies essentially as the hyperbolic-secant function described earlier. However, a more rapid fall-off with frequency can be obtained at higher frequencies if base-spreading resistance is appreciable. The open-circuit collector-base admittance can be described as the parallel combination of the collector-base junction capacity, the conductance due to space-charge-layer widening, and a frequency-sensitive (complex) diffusion capacity also due to space-charge-layer widening. At high frequencies, the principal effect of the diffusion capacitance is to increase the conductance of the collector-base admittance. Base-spreading resistance also may increase high-frequency collector-base conductance. Short-circuit input impedance consists of two impedances in series. The first of these is the parallel combination of the original Shockley, et al., emitter resistance $(kT/q_e I_\epsilon)$ and a capacitance that is inversely proportional to this emitter resistance, while the second impedance is the base-spreading resistance multiplied by the factor $(1-\alpha)$. As frequency is increased, the latter term increases in magnitude and becomes inductive reactive, due to the increase of the quadrature component of α. Finally, the voltage-feedback parameter is the sum of two terms. The first of these is essentially frequency independent and describes the voltage feedback calculated by Early plus current feedback through the effective collector conductance and base-spreading resistance. On the other hand, the second term, which describes the current feedback through the effective collector capacitance and base-spreading resistance, is proportional to frequency and dominates at frequencies greater than a few tens of kilocycles.

II. Derivation of Voltage-Current Relations for a Theoretical Junction Transistor

A theoretical analysis for the voltage-current relations of a one-dimensional transistor was given originally by W. Shockley.[8] Subsequently, J. M. Early[4] extended this analysis to include the effect of space-charge-layer widening. The analysis to be described below is an extension of this work to include the effect of frequency variations. Accordingly, the basic features and assumptions of the original analysis will not be reviewed here.

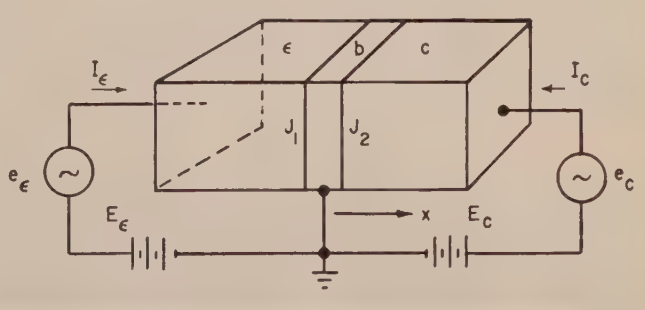

Fig. 14—Theoretical model of one-dimensional junction transistor.

A. Formulation

Consider the theoretical model of the junction transistor shown in Fig. 14. A *p-n-p* structure will be assumed for convenience; however, the final results also can be used for the *n-p-n* transistor by making only slight changes in notation, as will be indicated below.

The minority-carrier (hole) density p in the base region must satisfy the equation

$$\frac{\partial p}{\partial t} = \frac{p_B - p}{\tau_B} + D_p \frac{\partial^2 p}{\partial x^2}, \tag{19}$$

where p_B is the equilibrium minority-carrier concentration in the base region, τ_B is the lifetime of the minority carriers in the base, and D_p is the diffusion constant for holes.

A solution for p may be assumed in the form of a Fourier series in time; thus:

$$p(x, t) = p_B + \sum_{m=-\infty}^{\infty} p_m(x) \cdot \epsilon^{jm\omega t}. \tag{20}$$

Substitution of (20) in (19) yields an infinite set of equations of the form

$$\frac{\partial^2 p_m}{\partial x^2} - \Gamma_m{}^2 p_m = 0, \tag{21}$$

where

$$\Gamma_m \equiv (1/L_B)(1 + jm\omega\tau_B)^{1/2} \tag{22}$$

is a "propagation constant" associated with the frequency $m\omega$. Here L_B is the diffusion length for minority carriers in the base and is defined as

$$L_B \equiv (D_p\tau_B)^{1/2}. \tag{23}$$

Note that at low frequencies $\Gamma_m \approx \Gamma_0 = 1/L_B$ whereas at higher frequencies, Γ_m increases as the square root of the frequency. A solution for well-known equation (21) is

$$p_m(x) = A_m \cosh \Gamma_m x + B_m \sinh \Gamma_m x, \tag{24}$$

where A_m and B_m are constants to be determined by the boundary conditions.

1. *Boundary Conditions:* The boundary conditions to be satisfied by $p(x, t)$ are as follows:

$$\left.\begin{aligned} &\text{at } J_1,\ p(x, t) = p_B\epsilon^{a(E_\epsilon + e_\epsilon)} \\ &\text{at } J_2,\ p(x, t) = p_B\epsilon^{a(E_c + e_c)} \end{aligned}\right\} \tag{25}$$

where

$$a \equiv (q_e/kT), \tag{26}$$

and E_ϵ, e_ϵ, E_c, e_c, are the dc and ac emitter and collector voltages, respectively. It is assumed that each of the ac voltages may be expressed as a Fourier series in time, with a fundamental frequency ω. By choosing a suitable origin for the x-co-ordinate system and by imposing (25) on (20) and (24), solutions for A_m and B_m could be obtained in a straightforward manner. However, in general, the exact locations of junctions J_1 and J_2 will depend upon the thickness of the space-charge layer at each junction. Since the thickness is a function of the voltage across the junction, the exact locations of J_1 and J_2 will be instantaneous functions of time. Accordingly, in the (much-distorted) co-ordinate system shown in Fig. 15, the instantaneous location of junction J_1 is defined as $\delta \equiv -\delta_1$, where

$$\delta_1 = \left(\frac{\partial \delta}{\partial E_\epsilon}\right) e_\epsilon + \frac{1}{2!}\left(\frac{\partial^2 \delta}{\partial E_\epsilon{}^2}\right) e_\epsilon{}^2 + \cdots, \tag{27}$$

and the relation between δ and E_ϵ depends upon the nature of the emitter-base junction.[4] Similarly, the location of J_2 is defined as

$$w \equiv w_0 + w_1 \tag{28}$$

where

$$w_1 = \left(\frac{\partial w}{\partial E_c}\right) e_c + \frac{1}{2!}\left(\frac{\partial^2 w}{\partial E_c{}^2}\right) e_c{}^2 + \cdots, \tag{29}$$

and the relation between w and E_c depends upon the nature of the collector-base junction.

2. *Linearity.* In order to obtain the desired linear, small-signal ac operation of the transistor *for the general case* it will be assumed that the magnitudes of e_ϵ and e_c are small relative to the thermal voltage $1/a = kT/q_e$. Hence, (25) may be written as

$$\text{at } J_1\text{: } p(\delta, t) \doteq p_B\epsilon^{aE_\epsilon}[1 + ae_\epsilon], \tag{30a}$$

$$\text{at } J_2\text{: } p(w, t) \doteq p_B\epsilon^{aE_c}[1 + ae_c], \tag{30b}$$

$$|ae| \ll 2.$$

Furthermore, it will be assumed that $|e_\epsilon|$ and $|e_c|$ are small relative to $|E_\epsilon|$ and $|E_c|$ respectively. In particular, if $|e| \ll 2|E|$, the relations between w_1 and e_c and between δ_1 and e_ϵ may be taken to be linear, i.e.,

$$w_1 \doteq (\partial w/\partial E_c)e_c;\ \delta_1 \doteq (\partial\delta/\partial E_\epsilon)e_\epsilon,\ |e| \ll 2|E|. \tag{31}$$

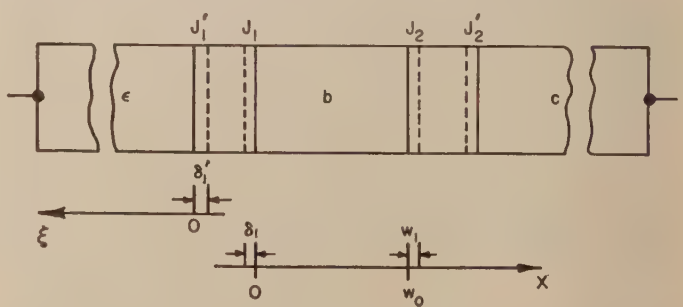

Fig. 15—Co-ordinate system for theoretical model of a junction transistor.

For *normal* transistor operation with the emitter-base junction biased in the forward direction, the approximation (30) is more restrictive than (31) for e_ϵ, since the terms involving δ_1 become insignificant compared to those involving ae_ϵ. On the other hand, with the collector-base junction biased in the reverse direction, the approximation (30) is not necessary for e_c, since the terms involving $(a\epsilon^{aE_c})e_c$ are insignificant compared to those involving w_1. In other words, to obtain results for *normal* transistor operation, the effect of space-charge-layer widening at the emiter-base junction may be neglected, whereas the classical reverse-bias admittance of the collector-base junction may be neglected, as was done by Early.

Returning to the more general case, additional assumptions that will be made to insure linearity are as follows:

a) Magnitude of ac junction displacements δ_1 and w_1 small relative to diffusion length L_B, i.e., $|w_1|, |\delta_1| \ll L_B$.

b) Also, $|\Gamma_1 w_1| \ll 1$, $|\Gamma_1 \delta_1| \ll 1$, is assumed at all frequencies of interest. This requires that the diffusion time across the distance $|\delta_1|$ or $|w_1|$ be small relative to periods of all frequencies of interest. Both a) and b) above are easily satisfied in general.

B. Solution for Minority-Carrier Current

By imposing the requirements for small-signal linear operation upon the more general results given above in (20) through (29), (approximate) first-order solutions may be obtained for the dc and fundamental ac ($m=0$, 1, respectively) minority-carrier concentrations as follows. Equation (24) may be substituted in (20) with only $m=0$, 1 terms retained. Imposing (30a, b) upon the result evaluated once at $x=-\delta_1$ and once at $x=w_0+w_1$ respectively, with δ_1 and w_1 defined by (31), yields two equations involving the four constants A_0, B_0, A_1, B_1. (In applying both (30) and (31), e_e and e_c should be replaced by $e_e\epsilon^{j\omega t}$ and $e_c\epsilon^{j\omega t}$ respectively.) Since dc and ac solutions must hold simultaneously, these two equations actually yield four simultaneous equations from which A_0, B_0, A_1 and B_1 may be determined. Note, however, that *before* splitting the equations into dc and ac parts, it is necessary to expand *each* of the hyperbolic functions into dc and ac terms, which arise from the "moving" boundary conditions, e.g.,

$$\begin{aligned}\sinh \Gamma w &= \sinh \Gamma(w_0 + w_1)\\ &= \sinh \Gamma w_0 \cosh \Gamma w_1 + \cosh \Gamma w_0 \sinh \Gamma w_1.\end{aligned}$$

By imposing assumptions (a) and (b) above, $\cosh \Gamma w_1$ may be replaced by unity and $\sinh \Gamma w_1$ may be replaced by Γw_1. Finally, second-order terms of the form $A_1\delta_1$, $B_1 w_1$ (which are proportional to e_e^2 or $e_e e_c$, for example) may be neglected.

Under these conditions, the above-mentioned simultaneous equations may be solved as follows:

$$\left.\begin{aligned}A_0 &= p_B(\epsilon^{aE_e}-1)\\ B_0 &= p_B[(\epsilon^{aE_c}-1)\ \operatorname{csch}\ \Gamma_0 w_0 - (\epsilon^{aE_e}-1)\ \coth\ \Gamma_0 w_0]\end{aligned}\right\}. \quad (32)$$

$$\left.\begin{aligned}A_1 &= [(ap_B\epsilon^{aE_e}) + B_0\Gamma_0(\partial\delta/\partial E_e)]e_e\\ B_1 &= \{[(ap_B\epsilon^{aE_c}) - \Gamma_0(\partial w/\partial E_c)\cdot(A_0 \sinh \Gamma_0 w_0\\ &\quad + B_0 \cosh \Gamma_0 w_0)](\operatorname{csch} \Gamma_1 w_0)e_c - A_1 \coth \Gamma_1 w_0\}\end{aligned}\right\}. \quad (33)$$

By substituting these values in (24) for $m=0$, 1, the dc and first-harmonic ac solutions respectively for $p(x, t)$ may be obtained if desired.

The minority-carrier (hole) current $I_p(x) \equiv I_{p0}(x) + i_{p1}(x)$ in the base may be obtained from $p(x, t)$ from the relation

$$I_p(x) = -Sq_e D_p(\partial p/\partial x) \quad (34)$$

where S is the cross-section area of the transistor. The hole current into the base from the emitter is obtained from this relation evaluated at $x=\delta$, while $I_p(w)$ represents the hole current into the collector from the base. After some algebraic manipulation, it follows that:

$$\left.\begin{aligned}I_{p0}(\delta) &= -Sq_e D_p\Gamma_0 B_0\\ i_{p1}(\delta) &= -Sq_e D_p[\Gamma_1 B_1 - \Gamma_0^2 A_0(\partial\delta/\partial E_e)e_e]\end{aligned}\right\} \quad (35)$$

$$\left.\begin{aligned}I_{p0}(w) &= -Sq_e D_p\Gamma_0[A_0 \sinh \Gamma_0 w_0 + B_0 \cosh \Gamma_0 w_0]\\ i_{p1}(w) &= -Sq_e D_p\{\Gamma_1[A_1 \sinh \Gamma_1 w_0 + B_1 \cosh \Gamma_1 w_0]\\ &\quad + \Gamma_0^2(\partial w/\partial E_c)(e_c)[A_0 \cosh \Gamma_0 w_0\\ &\quad + B_0 \sinh \Gamma_0 w_0]\}\end{aligned}\right\}. \quad (36)$$

Explicitly, $i_{p1}(\delta)$ may be written as

$$\begin{aligned}i_{p1}(\delta) = +Sq_e D_p p_B &\left\{(a\epsilon^{aE_e})(\Gamma_1 \coth \Gamma_1 w_0) - \left(\frac{\partial\delta}{\partial E_e}\right)\right.\\ &\times\Gamma_0(\Gamma_1 \coth \Gamma_1 w_0)[(\epsilon^{aE_e}-1)\coth \Gamma_0 w_0 - (\epsilon^{aE_c}-1)\\ &\left.\times \operatorname{csch} \Gamma_0 w_0] + \left(\frac{\partial\delta}{\partial E_e}\right)\Gamma_0^2(\epsilon^{aE_e}-1)\right\} e_e - Sq_e D_p p_B\\ &\times\left\{(\Gamma_1 \operatorname{csch} \Gamma_1 w_0)(a\epsilon^{aE_c})\right.\\ &+\left(\frac{\partial w}{\partial E_e}\right)\Gamma_0(\Gamma_1 \operatorname{csch} \Gamma_1 w_0)[(\epsilon^{aE_e}-1)\\ &\left.\times \operatorname{csch} \Gamma_0 w_0 - (\epsilon^{aE_c}-1)\coth \Gamma_0 w_0]\right\} e_c. \quad (37)\end{aligned}$$

The expression for $i_{p1}(w)$ is similar to that for $i_{p1}(\delta)$ and in fact may be obtained from the latter by interchanging E_e and E_c, e_e and e_c, δ and $-w$ and by replacing w_0 by $-w_0$.

1. *Minority-Carrier Density in Emitter and Collector Regions:* In addition to the hole current in the base region, there are minority-carrier (electron) currents in the emitter and collector regions which must be included in calculating the total emitter and collector currents, respectively. Referring to the ξ-co-ordinate system of Fig. 15, the electron current $I_n(\xi)$ may be determined from the electron density $n(\xi)$ in the emitter from a relation similar to (34):

$$I_n(\xi) = +Sq_e D_n(\partial n/\partial\xi). \quad (38)$$

The electron density $n(\xi)$ must satisfy a diffusion equation of the same form as (19) for $p(x)$, and a general solution of the type assumed for $p(x)$ in (20) and (24) may be obtained for $n(\xi)$.

One of the boundary conditions which is applied to $n(\xi)$ is quite similar to that applied to $p(x)$ at J_1, viz., that at the emitter-base junction J_1', $(\xi=\delta'\equiv-\delta_1')$,

$$\text{at } J_1',\ n(\xi, t) = n_e\epsilon^{a(E_e+e_e)},$$

where n_e is the equilibrium concentration of electrons in the emitter region. On the other hand, the second boundary condition applied to $n(\xi)$ is that at the contact

end of the emitter region $\xi \equiv d_\epsilon$, $n(\xi, t) = n_\epsilon$, i.e., the excess concentration shall be zero. By applying these boundary equations, by employing the same assumptions as in the case of $p(x)$, and by substituting the resulting expression for $n(\xi)$ in (38), the electron current $I_n(\delta') = I_{n0}(\delta') + i_{n1}(\delta')$ across J_1' (in the direction of positive ξ) may be calculated as:

$$\left.\begin{aligned} I_{n0}(\delta') &= -Sq_eD_n\Gamma_{\epsilon 0}n_\epsilon(\epsilon^{aE_\epsilon} - 1)\coth \Gamma_{\epsilon 0}d_\epsilon \\ i_{n1}(\delta') &= -Sq_eD_nn_\epsilon\{(a\epsilon^{aE_\epsilon})\Gamma_{\epsilon 1}\coth \Gamma_{\epsilon 1}d_\epsilon \\ &\quad - \Gamma_{\epsilon 0}(\epsilon^{aE_\epsilon} - 1)(\partial\delta'/\partial E_\epsilon)[\Gamma_{\epsilon 1}(\coth \Gamma_{\epsilon 1}d_\epsilon) \\ &\quad \times (\coth \Gamma_{\epsilon 0}d_\epsilon) - \Gamma_{\epsilon 0}]\}e_\epsilon \end{aligned}\right\} \quad (39)$$

where

$$\Gamma_{\epsilon m} \equiv (1/L_\epsilon)(1 + jm\omega\tau_\epsilon)^{1/2}, \quad m = 0, 1, \quad (40)$$

while L_ϵ and τ_ϵ are the diffusion distance and lifetime, respectively, for electrons in the emitter.

An exactly similar set of results may be obtained for the electron current $I_n(w')$ into the collector region by replacing Γ_ϵ, n_ϵ, E_ϵ, e_ϵ, d_ϵ and δ' with the corresponding quantitites for the collector region.

It should be noted in either case that if d is greater than approximately 3 or 4 diffusion lengths, i.e., $\Gamma_0 d \geqq 3$, coth Γd may be replaced by unity with an error of less than one per cent. This is equivalent to describing $n(\xi)$ by a single exponential $\epsilon^{-\Gamma\xi}$, as in Shockley's[8] analysis, rather than using hyperbolic functions for $n(\xi)$.

Total dc and ac emitter currents I_ϵ and i_ϵ respectively (positive direction into the emitter as in Fig. 14) may be obtained by adding the hole current $I_p(\delta)$ across J_1 and the electron current $-I_n(\delta')$, given by (35) and (39) respectively. Furthermore, the reactive current due to the barrier capacitance C_ϵ of the emitter-base junction must be included. Thus:

$$\left.\begin{aligned} I_\epsilon &= I_{p0}(\delta) - I_{n0}(\delta') \\ i_\epsilon &= i_{p1}(\delta) - i_{n1}(\delta') + j\omega C_\epsilon e_\epsilon \end{aligned}\right\}. \quad (41)$$

Similar equations may be written for the dc and ac collector currents. With positive direction of collector current taken into the collector, these become:

$$\left.\begin{aligned} I_c &= -I_{p0}(w) - I_{n0}(w') \\ i_c &= -i_{p1}(w) - i_{n1}(w') + j\omega C_c e_c \end{aligned}\right\}, \quad (42)$$

where C_c is the barrier capacitance of the collector-base junction.

C. Normal Transistor Operation

The general results obtained above now will be restricted to the case of normal transistor operation. In this case the emitter-base junction is biased in the forward direction, and the collector-base junction is biased in the reverse direction such that $|E_c| \ll 1/a$, i.e., with a magnitude of dc collector voltage greater than, say, 0.5 volt. Under these conditions, ϵ^{aE_c} may be neglected relative to unity; furthermore, terms involving $a\epsilon^{aE_c}$ may be neglected relative to those involving $(\partial w/\partial E_c)$. It also can be shown[18] that the terms involving $(\partial\delta/\partial E_\epsilon)$ may be neglected relative to those involving $a\epsilon^{aE_\epsilon}$.

Under these conditions, and with the additional assumption that emitter and collector regions are sufficiently wide that coth $\Gamma d \approx 1$, (35), (36) and (32) may be combined, and the results may be substituted in (41) and (42), together with (39) plus similar equations for $I_n(w')$ to yield for the dc equations:

$$\left.\begin{aligned} I_\epsilon &= I_0(E_\epsilon)[\coth \Gamma_0 w_0 + K_{\epsilon 0}] - I_0(E_c)[\operatorname{csch} \Gamma_0 w_0] \\ I_c &= -I_0(E_\epsilon)[\operatorname{csch} \Gamma_0 w_0] + I_0(E_c)[\coth \Gamma_0 w_0 + K_{c0}] \end{aligned}\right\} \quad (43)$$

where

$$\left.\begin{aligned} I_0(E) &\equiv I_0[\epsilon^{aE} - 1] \\ I_0 &\equiv (Sq_eD_p p_B\Gamma_0) \end{aligned}\right\} \quad (44)$$

$$K_{\epsilon 0} \equiv \left(\frac{D_n n_\epsilon \Gamma_{\epsilon 0}}{D_p p_B \Gamma_0}\right) \quad (45)$$

and

$$K_{c0} \equiv \left(\frac{D_n n_c \Gamma_{c0}}{D_p p_B \Gamma_0}\right). \quad (46)$$

Except for the difference in notation, these are identical with the equations given by Shockley.[19] By a similar combination, the ac equations may be written as

$$\left.\begin{aligned} i_\epsilon &= \{F_1(E_\epsilon)\cdot(\Gamma_1/\Gamma_0)[\coth \Gamma_1 w_0 + K_{\epsilon 1}] + j\omega C_\epsilon\}e_\epsilon \\ &\quad - F_2(E_\epsilon)\cdot(\Gamma_1/\Gamma_0)[\operatorname{csch} \Gamma_1 w_0]e_c, \\ i_c &= -F_1(E_\epsilon)\cdot(\Gamma_1/\Gamma_0)[\operatorname{csch} \Gamma_1 w_0]e_\epsilon + \{F_2(E_\epsilon) \\ &\quad \times (\Gamma_1/\Gamma_0)[\coth \Gamma_1 w_0] - F_3 + F_4[(\Gamma_{c1}/\Gamma_{c0}) - 1] \\ &\quad + j\omega C_c\}e_c, \end{aligned}\right\} \quad (47)$$

where

$$F_1(E_\epsilon) \equiv (Sq_eD_p\Gamma_0 p_B a\epsilon^{aE_\epsilon}), \quad (48)$$

$$F_2(E_\epsilon) \equiv Sq_eD_p\Gamma_0{}^2p_B[(\epsilon^{aE_\epsilon} - 1)\operatorname{csch} \Gamma_0 w + \coth \Gamma_0 w_0]\left(\frac{\partial w}{\partial E_c}\right), \quad (49)$$

$$F_3(E_\epsilon) \equiv Sq_eD_p\Gamma_0{}^2p_B\left(\frac{\partial w}{\partial E_c}\right), \quad (50)$$

$$F_4 \equiv Sq_eD_n\Gamma_{c0}{}^2n_c\left(\frac{\partial w'}{\partial E_c}\right) \quad (51)$$

and

$$K_{\epsilon 1} \equiv (D_n\Gamma_{\epsilon 1}n_\epsilon/D_p\Gamma_1 p_B). \quad (52)$$

[18] For example, detailed consideration of (37) and the corresponding equation for $i_{p1}(w)$ shows that the ratio of the terms involving ϵ^{aE_ϵ} to those involving $(\partial\delta/\partial E_\epsilon)$ may be expressed approximately as $[a \div (1/w_0)(\partial\delta/\partial E_\epsilon)]$. If a value of $(\partial\delta/\partial E_\epsilon)$ of the order of 2×10^{-6} meters/volt is assumed (for a fused-junction transistor, having an internal emitter-base potential of say 0.5v with a forward bias voltage of say 0.2v), and if $w_0 \geqq 10^{-5}$ meters, then this term becomes $\sim[40 \div 0.2]$. Hence, the magnitude of the term involving space-charge-layer widening at the emitter is only of the order of 1 per cent of that of the term due to forward-bias admittance.

[19] Shockley (footnote 8), (6-1), (6-2), p. 472.

It should be emphasized that these results also may be used for an n-p-n transistor merely by replacing E_ϵ by $-E_\epsilon$ (or by $|E_\epsilon|$), $(\partial w/\partial E_c)$ by $|\partial w/\partial E_c|$, and by interchanging n and p, e.g., $F_1 \propto D_n n_B$, $K_{\epsilon 1} \propto (D_p p_\epsilon / D_n n_B)$. (Note that the sign of $(\partial w/\partial E_c)$ is + for the *p-n-p* transistor and − for the case of *n-p-n*.)

1. *Simplification of the ac Equations:* In order to help simplify further calculations, the parameters F_1 and F_2 may be expressed in terms of the dc emitter and collector currents as follows: Equation (43) for I_c (with E_c negative and large) may be manipulated to yield

$$I_0[(\epsilon^{aE_\epsilon} - 1)(\text{csch } \Gamma_0 w_0) + \coth \Gamma_0 w_0] = |I_c + I_{cn}| \equiv I_c', \quad (53)$$

where

$$I_{cn} \equiv I_0 K_{c0} = q_e D_n n_c \Gamma_{c0} \quad (54)$$

is that part of I_c due to *electron* current from the collector into the base. In general, I_{cn} will be small (e.g., a few μa). Furthermore, it can be shown that $I_{cn} < I_{c0}$, where I_{c0} is the usual collector current corresponding to zero dc emitter current.

Similarly, (43) for I_ϵ may be manipulated to yield

$$[Sq_e D_p p_B \Gamma_0 \epsilon^{aE_\epsilon}(\coth \Gamma_0 w_0 + K_{\epsilon 0})] = |I_\epsilon + I_{\epsilon s}| \equiv I_\epsilon' \quad (55)$$

where

$$I_{\epsilon s} \equiv I_0[\coth \Gamma_0 w_0 - \text{csch } \Gamma_0 w_0 + K_{\epsilon 0}] \quad (56)$$

is a saturation emitter current that flows when both emitter and collector are biased in the inverse direction. Here also it can be shown that $I_{\epsilon s} < I_{c0}$. For most practical purposes with values of dc emitter current much greater than I_{c0}, I_ϵ' I_c' may be replaced by $|I_\epsilon|$ and $\alpha_0 |I_\epsilon|$ respectively.

By combining (53) and (55) with (48) and (49), F_1 and F_2 may be written as

$$F_1 = (\gamma_0 I_\epsilon' a \Gamma_0 w_0)[\tanh (\Gamma_0 w_0)/\Gamma_0 w_0], \quad (57)$$

$$F_2 = I_c' \Gamma_0 (\partial w/\partial E_c) \approx \alpha_0 |I_\epsilon| \Gamma_0 (\partial w/\partial E_c). \quad (58)$$

For moderately good transistors (e.g., $\alpha > 0.93$) the bracketed term of (57) may be approximated by unity (with an error <5 per cent). In this case, and when $I_\epsilon \gg I_{c0}$, F_1 may be simplified further to read

$$F_1 \approx (\gamma_0 |I_\epsilon| a \Gamma_0 w_0) = (\gamma_0 |I_\epsilon| \Gamma_0 w_0 q_e / kT). \quad (57a)$$

Further simplification of the results will be made by neglecting the terms F_3 and F_4. F_3 is a pure conductance and may be neglected with respect to the conductance term involving F_2, except at extremely small emitter bias voltages. On the other hand, the term involving F_4 is essentially zero at low frequencies, but will have a real and an imaginary part at higher frequencies. However, it should be kept in mind that this term might be of some importance at very low emitter-bias voltages in some types of transistors (e.g., if the minority-carrier concentration in the collector is much greater than that in the emitter $[n_c \gg p_B]$; here, too, the space-charge layer extends more deeply into the collector region than into the base region so that $(\partial w'/\partial E_c) \gg (\partial w/\partial E_c)$).

2. *Effect of Base Spreading Resistance; Equivalent Circuit for Theoretical Model:* In an actual transistor the effect of spreading resistance of the base region must be taken into account. To the first approximation, this may be done in the usual manner by adding a lumped resistance r_b' in series with the base terminal of the theoretical model. This is shown in Fig. 16, together with an exact equivalent circuit for the theoretical model. The circuit contains two "leaky" RC transmission lines and is based on the voltage-current equations (47) with the terms involving F_3 and F_4 omitted. Each line has a propagation constant Γ_1, but the characteristic impedance of the line in the emitter circuit is lower than that in the collector circuit by the ratio F_2/F_1. Actually the circuit of Fig. 16 is of more academic than practical interest. A considerably more useful and fairly accurate circuit, described earlier,[15] can be obtained by replacing (low impedance) line in the emitter circuit by a lumped resistance and capacitance in parallel, by replacing the voltage generator in the collector-circuit line by a current generator, and by neglecting shunt leakage in the collector-circuit line. (See Fig. 13 also.)

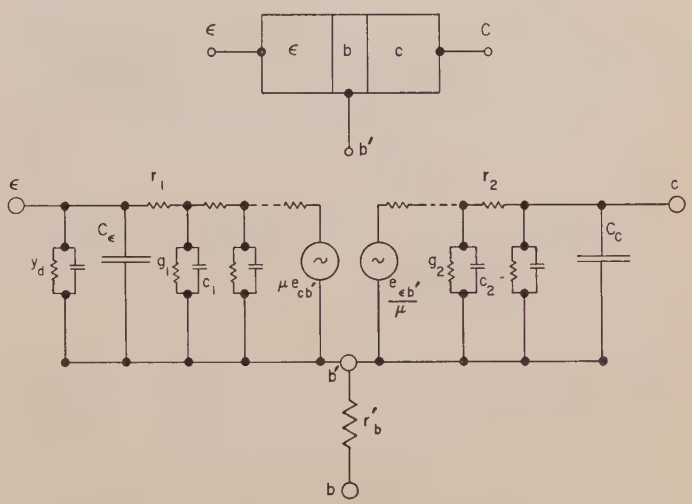

Fig. 16—Exact equivalent circuit for junction transistor.

In order to calculate the performance of the over-all transistor including the effect of r_b', this writer prefers to use the $[h]$, or series-parallel, parameters. These parameters have been defined by (1) of section I. The admittance, or $[y]$, representation which is so useful for the theoretical model is not so simple when r_b' is included. However, other parameters may be calculated for the over-all transistor as desired.[20]

Accordingly, the admittance parameters of the theoretical model given by (47) have been converted to $[h]$ parameters, and the results are given below. Modifica-

[20] H. Johnson of the RCA Laboratories has described a calculation of the grounded-emitter admittance parameters (from an independent derivation of the voltage-current relations of the theoretical transistor) at the Transistor Research Conference, State College, Pa., July 6, 1953.

tions in the $[h]$ parameters due to the addition of r_b' are given in section I, where each parameter is discussed in detail.

By applying the transformation,[21]

$$\begin{bmatrix} h_{11} & h_{12} \\ h_{21} & h_{22} \end{bmatrix} = \begin{bmatrix} (1/y_{11}) & -(y_{12}/y_{11}) \\ (y_{21}/y_{11}) & y_{22} - (y_{12}y_{21}/y_{11}) \end{bmatrix}$$

to (47) with F_3 and F_4 terms neglected, it follows that

$$h_{11} = (\Gamma_0 w_0/F_1)\cdot\gamma\cdot[\tanh(\Gamma_1 w_0)/\Gamma_1 w_0], \tag{59}$$

where

$$\gamma \equiv \{1 + [K_{\epsilon 1}\Gamma_1 w_0 + j\omega C_e(\Gamma_0 w_0/F_1)] \cdot [\tanh(\Gamma_1 w_0)/\Gamma_1 w_0]\}^{-1}, \tag{60}$$

$$h_{21} = -\gamma\cdot\operatorname{sech}\Gamma_1 w_0, \tag{61}$$

$$h_{12} = -(F_2/F_1)\cdot h_{21}, \tag{62}$$

$$h_{22} = j\omega C_c + (F_2/\Gamma_0 w_0)\{(\Gamma_1 w_0)^2[\tanh(\Gamma_1 w_0)/\Gamma_1 w_0] + (\Gamma_1 w_0)\cdot K_{\epsilon 1}\cdot\gamma\cdot(\operatorname{sech}\Gamma_1 w_0)^2\} + j\omega C_\epsilon(F_2/F_1)\gamma\operatorname{sech}^2\Gamma_1 w_0. \tag{63}$$

In general, the latter term in (63) is extremely small relative to $j\omega C_c$ and may be neglected.

D. Acknowledgments

The writer is very grateful to R. N. Hall for helpful discussions concerning this subject, for providing the physical picture of the diffusion capacity, and for providing grown-junction transistors used in some of the experimental work. The fused-junction transistors were supplied by J. S. Saby. Assistance of Miss Joyce Gillespie in carrying out most of the numerical calculations for the exact forms of the hyperbolic functions also is gratefully acknowledged.

E. Appendix A

Relationship between β-cut-off frequency and base width w for theoretical model of junction transistor.

In general,[2a,b] $\beta = \operatorname{sech}[(1+j\omega\tau_b)^{1/2}w/L_b]$.

The β-cut-off frequency $\omega_c/2\pi$ is defined as the frequency for which $|\beta|$ has decreased to $\frac{1}{2}\beta_0^2$, where β_0 is the low-frequency value of $\beta(\omega=0)$. Numerical calculations of β versus $(\omega w^2/D)$ (where $D\equiv L_b^2/\tau_b$) for several values of β_0 yielded the following values of ω_c:

β_0	$(\omega_c w^2/D)$
1.00	2.43
0.952	2.53
0.908	2.63
0.867	2.73
0.830	2.81

F. Appendix B

1. *Discussion of Diffusion Capacity C_d:* The diffusion capacity C_d due to charge stored in the base region by the dc emitter current was introduced in the approximate relationship (10) for collector-base admittance h_{22}', in order to provide a convenient analytical description for the high-frequency behavior of h_{22}'. More accurately, C_d may be defined as $1/j\omega$ times the limiting low-frequency value of the collector-base admittance due to space-charge-layer widening for the case of $\gamma_0=1$. Thus, from (63) with $K_{\epsilon 1}=0$, the admittance $(h_{22}-j\omega C_c)$ due to space-charge-layer widening may be written as

$$(h_{22} - j\omega C_c) = (F_2/\Gamma_0 w_0)(\Gamma_1 w_0)^2[\tanh(\Gamma_1 w_0)/\Gamma_1 w_0].$$

The limiting low-frequency value of the right-hand side of this expression is

$$(F_2/\Gamma_0 w_0)[1 + j\omega\tau_b](w_0/L_b)^2[\tanh(w_0/L_b)/(w_0/L_b)].$$

From this equation, and by using the relation $D_b\equiv L_b^2/\tau_b$, the diffusion capacity C_d as defined above becomes

$$C_d \equiv (w_0^2/D_b)(F_2/\Gamma_0 w_0)[\tanh(w_0/L_b)/(w_0/L_b)]. \tag{64}$$

Substituting for $F_2/\Gamma_0 w_0$ from (58) the diffusion capacity C_d may be written more explicitly as

$$C_d \approx \alpha_0 |I_\epsilon| (w_0/D_b)[\tanh(w_0/L_b)/(w_0/L_b)](\partial w/\partial E_c). \tag{65}$$

Hence, C_d is directly proportional to the thickness of the base region, to the dc emitter current, and to the rate of change of the base thickness with respect to collector voltage.

A useful approximate expression for C_d also may be written in terms of the concentration p_B of minority carriers in the base by using the approximation

$$I_\epsilon \approx (Sq_e D_p p_B/w_0)\epsilon^{aE_\epsilon},$$

which is valid for moderately high α (e.g., $\alpha>0.93$) and for moderately high currents (such that $\epsilon^{aE_\epsilon}\gg 1$), cf. (55); thus:

$$C_d \approx q_e S p_B(\partial w/\partial E_c)\epsilon^{aE_\epsilon}. \tag{66}$$

An alternative calculation of C_d may be obtained by direct consideration of the physics of the transistor in a manner suggested by R. N. Hall. Consider the model of the transistor shown in Fig. 17, for which a value of $\alpha_0=1$ is assumed for convenience. At dc equilibrium, $\Delta E_c=0$, the concentration of holes $p(x)$ through the base layer is essentially linear as shown by the solid line. At the emitter-base junction, $x=0$, $p(x)=p_\epsilon$, and at the base-collector junction, $x=w_0$, $p(x)=0$. The slope of the line $p(x)$ is proportional to the dc emitter current.

If now the magnitude of the collector-base voltage is increased by applying $-\Delta E_c$ as shown, the collector-base junction moves toward the emitter by an amount Δw. The emitter current is assumed to be held constant. Accordingly, the slope of the $p(x)$ line must remain the same. Furthermore, since the hole concentration must vanish at the new position of the base-collector junction, the hole concentration $p(x)$ corresponding to the new collector voltage must be as shown by the dotted line. This has two effects.[22] First the value of $p(x)$ at the emitter-base junction must decrease from p_ϵ by an amount Δp. The corresponding decrease in the emitter

[21] For example, Brown and Bennett (footnote 7), p. 851.

[22] If a value of α different from unity had been assumed, a third result would be collector-base conductance due to space-charge-layer widening as calculated by Early.

voltage produces the feedback voltage described by Early. Second, in order to permit the new boundary conditions to be attained, a quantity of charge ΔQ shown by the shaded area in Fig. 17 must be removed from the base region into the collector circuit. The ratio of this charge ΔQ to the change in voltage ΔE_c causing the charge to appear in the circuit is the effective diffusion capacity C_d. Thus

$$C_d = I_e(w_0/D)(\Delta w/\Delta E_c)$$

which is equivalent to the result obtained above in (65) for the case $(w/L_b)^2 \ll 1$, $\alpha \doteq 1$.

The term diffusion capacity is used to indicate that the charge contributing to C_d must diffuse from the base region to the base-collector junction. As the frequency of the ac collector voltage is increased, the time required for the charge at the emitter end of the base to diffuse to the base-collector junction becomes comparable with the period of the ac voltage, and C_d effectively becomes complex. The quantitative description of this process is best found by solution of diffusion equation above.

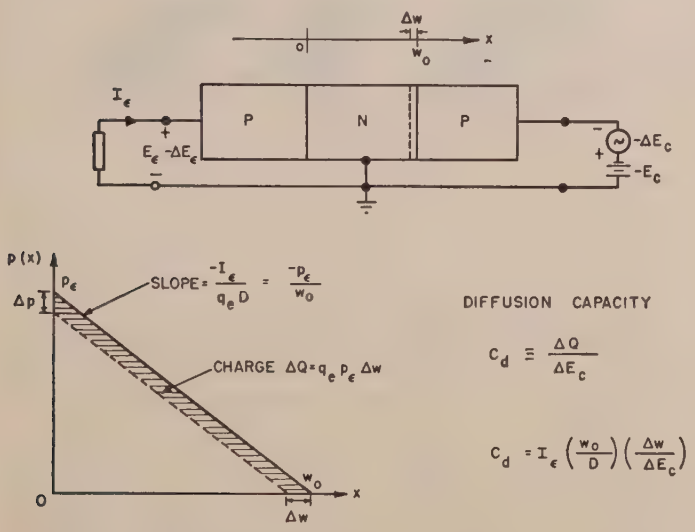

Fig. 17—Model of junction transistor to illustrate physical concept of collector-base diffusion capacitance.

2. *Variation of C_d with dc Bias:* The diffusion capacity is seen to be directly proportional to dc emitter current. Obversely, the variation of C_d with dc collector voltage E_c depends upon the variation of $(\partial w/\partial E_c)$, which is determined by the nature of the collector-base junction.

For the uniformly graded junction, Early has given an expression for $\partial w/\partial E_c$. In the case of a nonuniformly graded junction, such as usually exists in a grown junction transistor, the value of $(\partial w/\partial E_c)$ depends upon the variation of carrier concentration with distance through the base-collector region and is not calculated easily in general.

For a step junction having a high conductivity on one side and uniform but much lower conductivity on the other side, the space-charge layer is essentially confined to the low-conductivity material. This is the case that normally exists in fused-junction transistors, when the base material is of relatively low conductivity (e.g., $\rho_b \geqq 1$ ohm cm). Here, Early[23] has given the results

$$(\partial w/\partial E_c) = -x_m/2E_c, \tag{67}$$

$$x_m^2 = 2\epsilon_r\epsilon_0 E_c/q_e N_B, \tag{68}$$

where x_m is the thickness of the collector barrier, ϵ_0 is the free-space dielectric constant, ϵ_r is the *relative* dielectric constant of the semiconductor material, and N_B is the concentration of *majority* carriers in the base region (N_B = donor concentration N_d, for n-type semiconductor). Alternatively, x_m may be related to the collector barrier capacity C_c by the equation[24]

$$C_c = \epsilon_0\epsilon_r S/x_m \tag{69}$$

Also,[24]

$$C_c = S(q_e N_B \epsilon_r \epsilon_0/2\,|\,E_c\,|\,)^{1/2}. \tag{70}$$

If (67) and (68) are combined (with the minus sign suppressed) and are substituted back into (65) for C_d, it will be seen that for a step junction

$$C_d = (\alpha_0\,|\,I_e\,|\,/\,|\,E_c\,|^{1/2})(w_0/D_b)[\tanh(w_0/L_B)/(w_0/L_B)] \times (\epsilon_r\epsilon_0/2q_eN_B)^{1/2}, \tag{71}$$

i.e., C_d is inversely proportional to the square root of collector voltage, Fig. 10, to square root of majority carrier concentration in base, and is directly proportional to dc emitter current and to base thickness.

Another very interesting result may be obtained by calculating the ratio C_d/C_c, using the approximate (66) combined with (67), (68) and (70):

$$C_d/C_c \doteq p_B\epsilon^{aE_e}/N_B. \tag{72}$$

The numerator of this expression represents the minority-carrier concentration injected into the base at the emitter-base junction, whereas the denominator is the equilibrium majority-carrier concentration. Now, one of the fundamental assumptions of the theoretical analysis given by Shockley, and hence of the extension described above, is that minority-carrier concentration be small relative to majority-carrier concentration. However, if C_d is comparable with C_c, this low-level injection requirement obviously is not well satisfied, and the theory from which the expression for C_d was derived is not strictly valid![25] However, since the experimental results for C_d are in fairly good agreement with those predicted by (71), up to values of C_d/C_c somewhat less than unity (cf. Fig. 10), it appears that the violation of the low-level injection assumption is not particularly serious so long as the injected minority-carrier current concentration does not exceed the *equilibrium* majority-carrier concentration in the base. Alternatively, these results indicate that direct measurement of the ratio C_d/C_c may be used to estimate the injection level if it is not known from other considerations.

[23] Early, (footnote 4), (17) and (15), p. 1406. Note that $(\partial w/\partial E_c) < 0$ for the n-p-n structure considered by Early, whereas $(\partial w/\partial E_c) > 0$ for the p-n-p case.

[24] For example, W. Shockley, 1949 (footnote 8), (2.55)–(2.56) p. 451, modified for the MKS system used above.

[25] An extension of Shockley's analysis for the case of high-level injection has been made recently by W. M. Webster of the RCA Laboratories, who has calculated the variation of α with dc emitter current. This paper was presented at the Transistor Research Conference, State College, Pa., July 7, 1953.

Electron Beam Focusing with Periodic Permanent Magnet Fields*

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Summary—**This paper presents experimental and theoretical results on focusing a long electron beam with a permanent magnet structure which produces a periodic magnetic field along the axis of the electron beam. In principle, the periodic permanent magnet made up of N cells will be reduced in weight by a factor of N^2 over a permanent magnet producing a uniform field over the same length. The periodic focusing results are comparable to results obtainable with a uniform field. In addition the periodic fields for certain combinations of beam voltage and magnetic field exhibit regions of "stop bands" wherein the beam becomes unstable and no current reaches the collector.**

Introduction

TRAVELING-WAVE tubes and most present-day microwave tubes require long electron beams of high current density. The radial forces of space charge in such dense beams must be neutralized if the beam is to be held together over a distance, and generally this is accomplished by immersing the beam in a uniform magnetic field. In principle such a field is most efficiently utilized in "Brillouin Flow"[1] wherein the cathode is shielded from the magnetic field. In practice, however, beams with shielded guns are focused with fields about 50 per cent greater than that for theoretical Brillouin Flow. The fields necessary for such focusing can be obtained from either an electromagnet such as a solenoid, or a permanent magnet. Both structures involve an inconvenient amount of weight, and in addition solenoids consume power.

These disadvantages of weight and power consumption can be overcome by using a periodic magnetic field which is symmetric about the axis of the electron beam. Such a field can be produced, as shown in Fig. 1, by a series of short magnetic lenses. The reduction in weight made possible by the use of alternating magnetic field is vividly illustrated in Fig. 2. The large magnet weighing 38 pounds is necessary to produce the uniform field[2] for focusing a beam in a medium power traveling-wave tube. The periodic magnetic structure, weighing 1 pound 5 ounces, will focus the same beam equally well. Thus the weight has been reduced by a factor of 30.

The possibility of focusing with alternating transverse magnetic fields has been pointed out by Livingston, Courant, and Snyder.[3] In their article on "Strong Focusing" they discuss the use of quadrupole fields for focusing high energy electron beams in which space charge forces are unimportant. Periodic fields of axial symmetry produced by electrostatic lenses have been discussed by Hahn.[4] Clogston and Heffner have shown

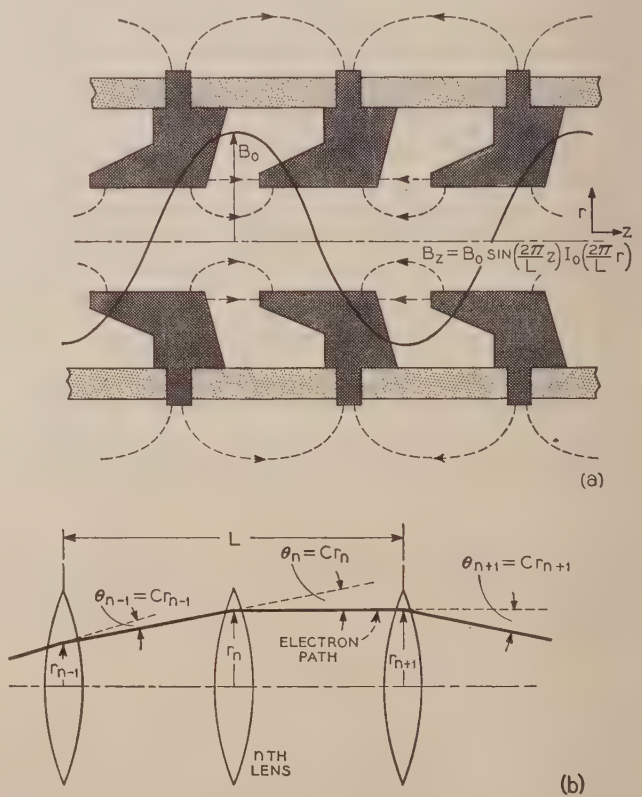

Fig. 1—A typical permanent magnet structure for producing an alternating magnetic field on the axis as shown in (a). A thin lens equivalent of this system is depicted in (b) to illustrate some of the properties of periodic focusing.

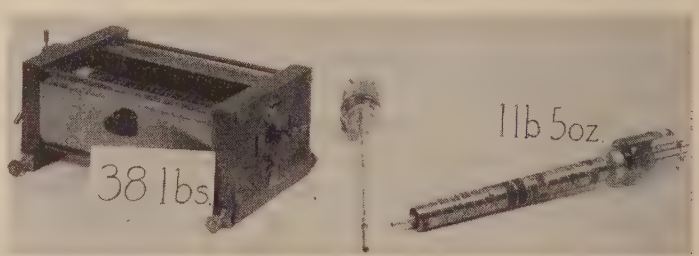

Fig. 2—This photograph illustrates the saving in weight which is possible in periodic focusing. The large magnet weighing 38 pounds is necessary to produce a field of 450 gauss over 9 inches of length. The periodic magnet weighing 1 pound 5 ounces will focus the same beam as the uniform field magnet.

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[1] J. R. Pierce, "Theory and Design of Electron Beams," D. Van Nostrand Co., New York, N. Y., chap. IX; 1954.

[2] The permanent magnet structure for producing uniform field shown in Fig. 2 was developed by C. C. Cutler and G. M. Eberhardt. The structure in the center is a series of mu metal discs necessary to eliminate the stray transverse fields which are present between the two Alnico bar magnets. Structures such as this can be used for fields up to 500 gauss. For higher fields the bars have to be replaced by horseshoe magnets.

[3] E. D. Courant, M. S. Livingston, and H. S. Snyder, "The strong focusing synchroton," *Phys. Rev.*, vol. 88, pp. 1190–1196; December, 1952.

[4] W. C. Hahn and G. F. Metcalf, "Velocity modulation tubes," PROC. I.R.E., vol. 27, pp. 106–117; February, 1939.

in a recent paper[5] that it is possible to overcome the forces of space charge and focus with periodic fields.

In order to focus a given beam with periodic magnetic fields it is necessary to use an rms value of the alternating field which is equal to the uniform "Brillouin" field required to focus the same beam. Therefore, since the effective field strength is not changed, only a minor improvement can be expected by the use of a series of electromagnets producing a periodic field in place of the use of a solenoid producing a uniform field. However, Pierce has pointed out that the weight of the permanent magnet required to produce an alternating field of a given rms value is much less than the weight of the magnet required to produce a uniform field. This saving in weight can be illustrated by the following argument.

In Fig. 3(a) there is illustrated a permanent magnet assumed to be of the most efficient design to produce a uniform axial field over the length of the magnet. Now suppose the same field is desired over a length N times as great. For the uniform field case it is necessary to scale all dimensions of the magnet by a factor N as shown in Fig. 3(b). For the periodic field, however, it is merely necessary to increase the length of the structure by a factor N as shown in Fig. 3(c). The volume of the uniform structure must, therefore, be increased by N^3 and the volume of the periodic structure by a factor N. Thus in the periodic structure the weight is smaller by a factor of N^2 over the uniform field where N is the number of lenses making up the periodic structure.[6] In practice such factors as the shape of the pole pieces and requirements for mechanical rigidity enter in such a way as to reduce the N^2 weight saving factor. In the experimental structure which will be discussed here the factor is about midway between N and N^2.

Before discussing the experimental results the important features of periodic focusing can be illustrated with the use of a few equations from Pierce's book on electron beams. Consider the periodic system of Fig. 1(b) which consists of a series of thin converging lenses. An electron moving through such a system of lenses will move in straight lines between successive lenses provided that the forces of space charge are not considered and secondly, that the fields are assumed to be closely confined to the region of the lens. The radial position of electron path at the nth lens is defined by r_n as in Fig. 1(b), the slope of the path by r_n' and the convergence of the lens C by the reciprocal of the focal length.

With these approximations the position of the path at the $n+1$ lens will be

$$r_{n+1} = r_n + \frac{L}{2} r_n'. \tag{1}$$

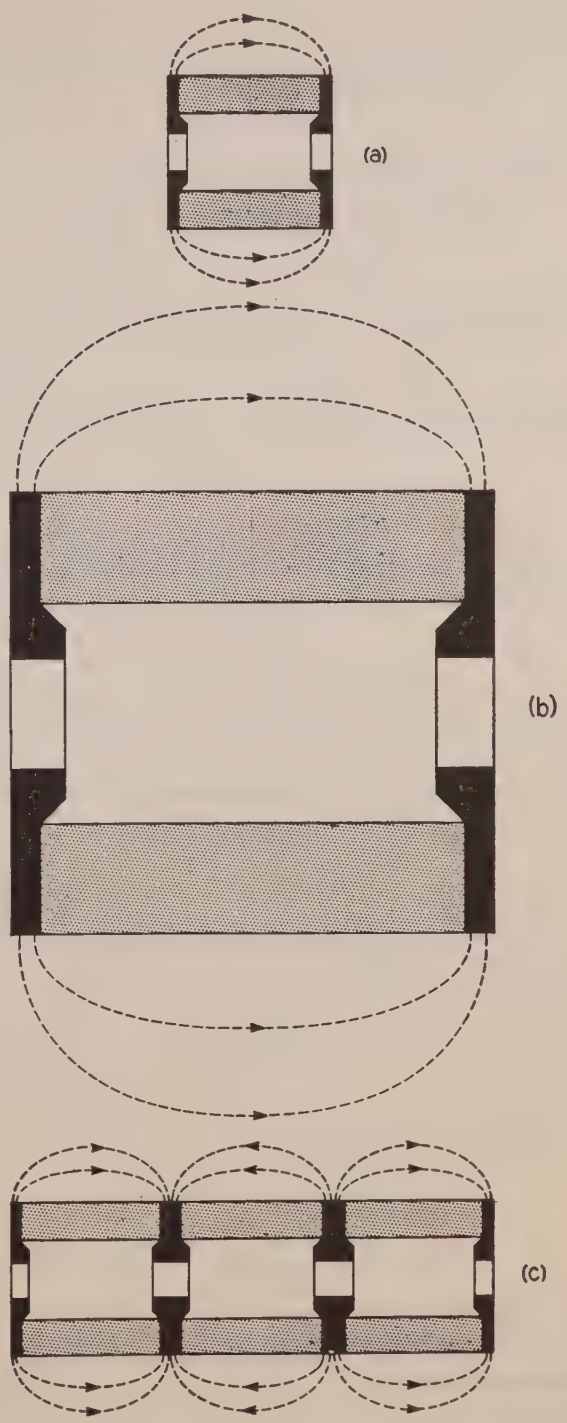

Fig. 3—The weight-saving feature of periodic focusing can be pointed out by increasing the length of the magnetic field by a factor of N. The uniform structure must be scaled in all dimensions by a factor of N as in (b), which means an increase in volume of N^3, whereas the periodic structure must be increased only in length by a factor of N as in (c). Thus the periodic structure has a weight advantage of N^2.

The change in slope of the path in passing through the nth lens will be Cr_n and thus

$$r_{n+1}' = r_n' - Cr_{n+1} \tag{2}$$

which gives from (1)

$$r_n' = \frac{2}{L}(r_{n+1} - r_n)$$

[5] A. M. Clogston and H. Heffner, "Focusing of an electron beam by periodic fields," to be published in the *Jour. of Appl. Phys.*

[6] The decrease in magnetic material required for the periodic structure can also be seen from a consideration of the external fields as discussed in Appendix A. Here it is shown that for the periodic case the external stored energy is reduced by a factor of $1/N^2$ over that of the uniform field structure. Also, it is important to note that periodic structures can be located in close proximity to each other without having appreciable interaction of the magnetic fields.

and similarly

$$r_{n+1}' = \frac{2}{L}(r_{n+2} - r_{n+1})$$

and hence (2) can be written

$$r_{n+2} - \left(2 - \frac{CL}{2}\right) r_{n+1} + r_n = 0. \tag{3}$$

The difference equation (3) has the solution of the form

$$r_n = A \cos n\theta + B \sin n\theta. \tag{4}$$

If the entering conditions are properly chosen, only the cos $n\theta$ solution is necessary, then (3) and (4) give

$$\cos\theta = \left(1 - \frac{CL}{4}\right). \tag{5}$$

Therefore the electron will move through a series of converging lenses with a stable orbit providing

$$CL < 8.$$

At $CL = 8$ the electron orbits are not periodic and the radius increases without limit and this point is defined as the beginning of a "stop band." For a magnetic lens system the convergence can be written[7]

$$C = \frac{\eta}{8V_0}\int_0^{L/2} B^2 dz. \tag{6}$$

With the sinusoidal variation of field along the axis as shown in Fig. 1(a) where

$$B_z = B_0 \cos\frac{2\pi z}{L} \tag{6a}$$

(6) becomes

$$CL = \frac{\eta L}{8V_0}\int_0^{L/2} B_0^2 \cos^2\frac{2\pi z}{L}\, dz$$

or

$$CL = \frac{\eta}{8V_0} B_0^2 \frac{L^2}{4}. \tag{7}$$

Therefore at the "stop band" given by $CL = 8$ (7) becomes

$$\frac{\eta B_0^2 L^2}{V_0} = 256 \tag{8}$$

and stable flow must occur for

$$\frac{\eta B_0^2 L^2}{V_0} < 256.$$

In a later section the thick lens case is treated including the effects of space charge and it is there shown that for stable flow with a minimum of ripple in the beam the rms value of the periodic field is equal to the dc magnetic field required for the uniform Brillouin case. Also when the thick lens is considered it is shown that (8) should be more accurately written

$$\frac{\eta B_0^2 L^2}{V_0} = 418. \tag{9}$$

[7] Pierce, *loc. cit.*, chap. VII.

Experimental Results

In the initial studies of periodic focusing it was desirable to vary the magnetic field over a wide range and for this purpose a system of electromagnetic lenses was constructed which produced a sinusoidal flux pattern at the axis, i.e.

$$B_z = B_0 \cos\frac{2\pi z}{L}.$$

The magnet spacing L was chosen to be 1 inch since this seemed a suitable value for a typical medium power traveling-wave tube (9 inch helix 0.080 inch in diameter) which was available for the experiment.

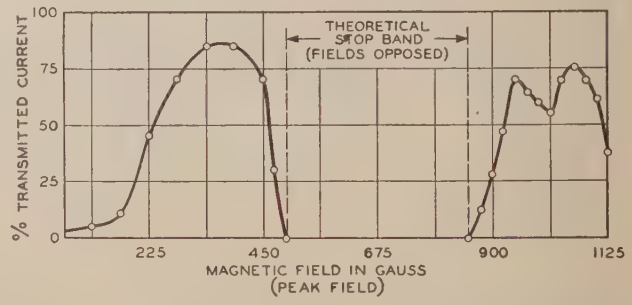

Fig. 4—Per cent collector current through a hole 9 inches in length and 0.080 inch in diameter as a function of the peak magnetic field of the periodic system. The first theoretical stop band is also shown for comparison. (Alternating opposed fields.)

Fig. 4 shows the results of changing the magnetic field with the beam current and voltage remaining constant. For small values of magnetic field the behavior is exactly like that of a uniform field solenoid, i.e., improved transmission with increasing magnetic field. However, at about 400 gauss peak field the transmitted current tends rapidly to zero and there is encountered a region of no transmission between 475 gauss and 850 gauss. At still higher magnetic fields alternate regions of electron transmission and no transmission are observed, although these are not shown in Fig. 4. This phenomenon has been found to be characteristic of periodic focusing. The regions of no transmission, termed "stop bands," occur for certain definite magnetic field strengths depending upon the magnet spacings and beam voltage. The first stop band begins when the magnetic field (peak value) is sufficiently high so that

$$\frac{\eta B_0^2 L^2}{V_0} = 418 \; MKS \text{ units} \tag{9}$$

as explained earlier. At the beginning of the next pass band the curve of percentage transmission is quite similar in shape to the first pass band. The width of the stop

and pass bands is calculated in a later section, and as can be seen from Fig. 4, the experimental curve agrees quite closely with theory.

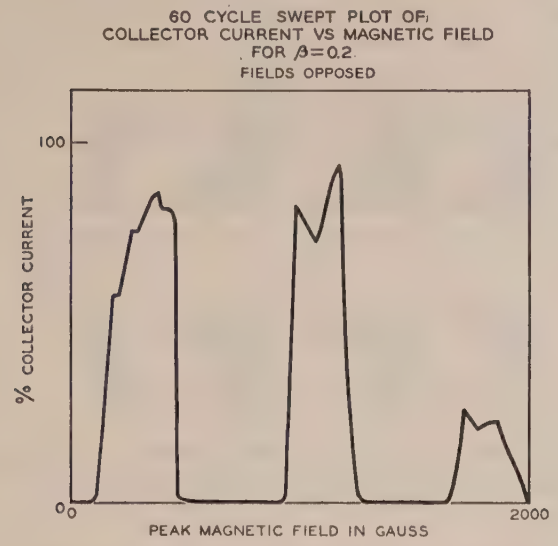

Fig. 5—Oscilloscope picture of the collector current when the magnetic field was swept at a 60 cycle rate. (Alternating opposed fields.) The third stop band and pass band appear in this picture.

A more graphic picture was obtained when the magnetic field was swept at 60 cps and the percentage collector current observed on an oscilloscope, as shown in Fig. 5. Had more magnetic field been available presumably more peaks and stop bands would have appeared to the right and the general pattern would have been repeated. In this experiment the higher pass bands were reached by increasing the magnetic field and for this condition the electrons cross the axis once in the second pass band and twice in the third pass band. As can be seen from Fig. 5, satisfactory focusing was achieved for the second pass band but not for the third. It should be pointed out that it may well be possible to operate in

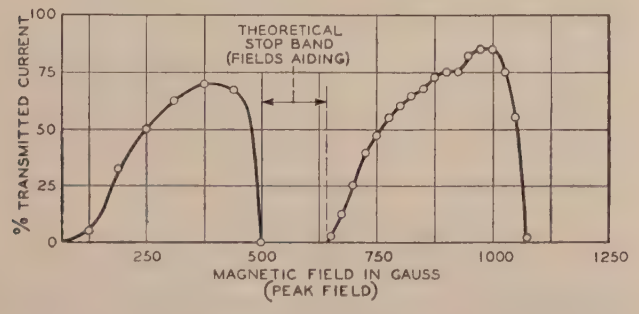

Fig. 6—Per cent collector current versus peak magnetic field when the fields are poled in the same direction (aiding fields) with same structure and beam as Fig. 4. This curve was taken for a structure with a 1-inch period, a beam voltage of 700 volts, and a beam current of 6 ma.

the second pass band by increasing the period rather than the magnetic field. The location of these higher order pass bands with respect to magnetic field, beam velocity, and magnet spacing is presented in a later section.

For the sake of completeness the magnetic field of each individual lens was directed in the same direction to give the "aiding field" case shown in Figs. 6 and 7.

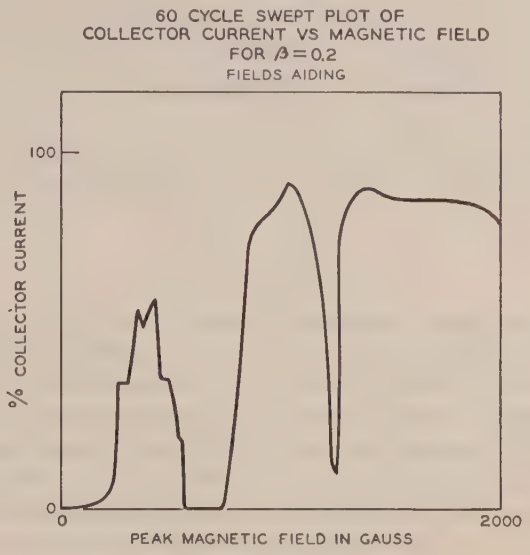

Fig. 7—Oscilloscope picture of the aiding field case with a swept magnetic field.

Here again the situation is of slight practical interest since such a configuration requires heavy permanent magnets.

This arrangement also gives alternate pass and stop bands which are slightly altered in position. The theory also predicts their location with good accuracy.

The optimum focusing conditions at the peak of the first pass band for this tube with alternately poled magnets were:

Helix potential	—700 volts
Collector current	—5.2 ma (88 per cent transmission)
Helix current	—0.7 ma
B_0	—380 gauss (peak field) 270 gauss (rms)
Brillouin field	—177 gauss (theoretical).

Although theory predicts perfect transmission with an rms field equal to the Brillouin field it has been determined experimentally that, for both uniform solenoidal fields and periodic fields, near perfect transmission is not realized until the magnetic field is about twice the Brillouin field. This is probably a result of thermal velocities and aberrations in the gun structure which preclude the attainment of ideal parallel flow at the start of the focusing structure.

The data presented in the previous figures were obtained from a traveling-wave tube whose gun structure was rather imperfectly designed. It was felt desirable to obtain information concerning the focusing properties of a tube with a higher perveance and also one whose gun produced a well-defined beam.

The tube selected for further testing was a medium power (10 watts) traveling-wave tube designed for application in radio-relays. The operating beam voltage is

about 1,760 volts with 38 ma, or a perveance of about 0.6×10^{-6}. A uniform field of about 350 gauss is the theoretical field for Brillouin focusing, so in the light of our experience the actual field required for good transmission (>99 per cent) should be approximately 700 gauss. To obtain an rms field of this value with a sinusoidal periodic field would necessitate a peak field of about 1,000 gauss.

With the magnet period of 1 inch used with the previous tube and a beam voltage of 1,700 volts,

$$\frac{\eta B_0^2 L^2}{V_0} > 418;$$

thus the operation would be beyond the first pass band and therefore no current transmission. The only parameter which could be conveniently changed was L, and so a system of closer lens spacing was constructed.

With a magnet period of $\frac{3}{4}$ inch and an rms magnetic field twice the Brillouin value, 99 per cent of the beam current was focused to the collector. Although this does not represent the maximum attainable percentage transmission it was considered sufficiently good for experimental verification of this type of focusing. It seemed desirable to obtain a careful check of the stability requirement $(\eta B_0^2 L^2/V_0)<418$ with this tube since the focusing was known to be good.

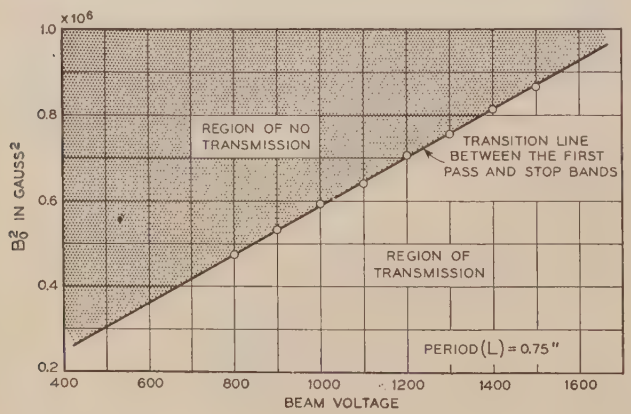

Fig. 8—Measurement of relationship of magnetic field and beam voltage for the beginning of the first stop band. For this test a beam with a perveance of 0.6×10^{-6} was used in conjunction with a 0.080-inch diameter helix 9 inches long.

The procedure followed was to adjust the magnetic field at each value of beam voltage so that the collector current just disappeared. These points represent the beginning of the first stop band and are plotted in Fig. 8. It will be noted that the relationship between B_0^2 and V_0 is a linear one and therefore establishes the criterion for stability. From this plot the stop band began at

$$\frac{\eta B_0^2 L^2}{V_0} = 360$$

rather than the analytically determined value of 418, but since this represents a difference of only 8 per cent in B (measurements of the field were only accurate to within about 5 per cent) it lies within the range of expected error. The important feature of this graph is the very close approximation to a straight line relationship between B_0^2 and V_0. This gives experimental confirmation to the theory that the mechanism of the stop band phenomenon is not dependent upon space-charge conditions.

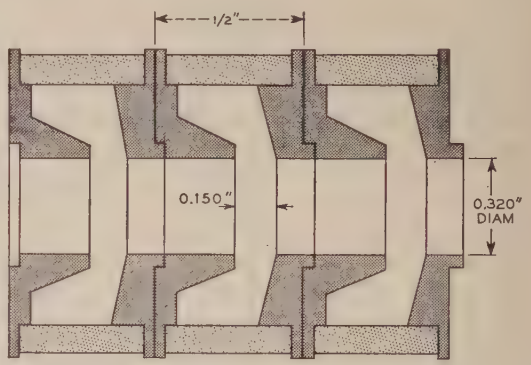

Fig. 9—Detail of a permanent magnet element used for periodic focusing. The Alnico V ring magnets and Permalloy pole pieces form a unit which can produce 700–800 gauss at the gap.

Experimental Results for a Permanent Magnet Circuit

The permanent magnet structure used to focus the medium power tube mentioned in the previous section shown in Fig. 9 has a 1-inch period and is made of Alnico V ring magnets with Permalloy pole pieces. Each cell was magnetized individually, and cells were stacked together to produce the alternating magnetic field. In these tests no great care was taken to equalize the strength of each cell and as a result there was about a 10 per cent variation in peak field from cell to cell as shown in Fig. 10.

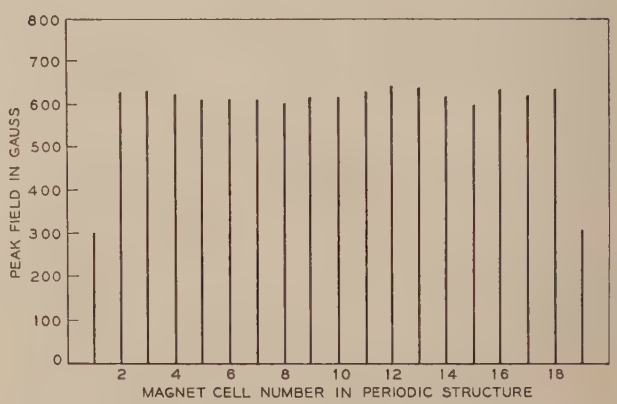

Fig. 10—Measurement of uniformity of peak fields occurring at each gap in a periodic structure made up of the elements shown in Fig. 9.

With the fixed magnetic field the collector current versus voltage was obtained as shown in Fig. 11 for a constant beam perveance. It should be noted that as the voltage is increased from zero all the stop bands are traversed, as predicted from (9), until the beginning of the primary pass band is reached near 1,000 volts, for an rms magnetic field of 423 gauss. As the voltage is increased above the stop band the collector current rises

rapidly. An increase of 100 volts above the edge of the stop band gives a current transmission of 70 per cent.

To reach higher values of current than shown in Fig. 11, it is necessary to go to higher values of magnetic field. Also, if the voltage at the stop band is to remain constant the period of the magnet structure must be decreased. If it were possible to scale the magnet shown in Fig. 9 the field strength would not vary. However, the tube in general does not permit the internal diameter of the structure to be decreased. Thus, as the length is decreased without a corresponding decrease in diameter the available peak field in the gap must decrease.

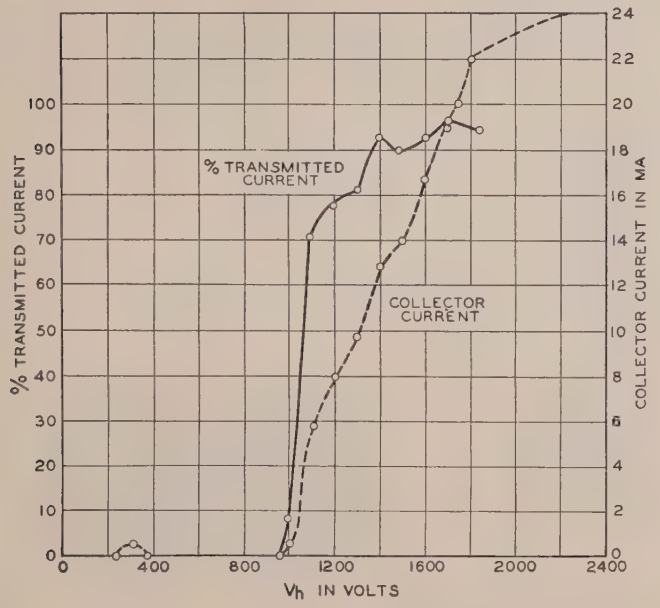

Fig. 11—Collector current versus beam voltage for an rms magnetic field of 423 gauss in the structure of Fig. 9. The beam transmission was measured through a hole 9 inches in length and 0.080 inch in diameter.

For example, in the magnet of Fig. 9 it is possible to obtain a peak field of 800 gauss near the axis which means a stable field of perhaps 700 gauss after partial demagnetization. For higher peak fields, it has been determined that Ferroxdure[8] magnets are suitable. With Ferroxdure cells as shown in Fig. 12 it was possible to achieve peak fields of 1,000 gauss with a period of 0.7 inch. It appears from Fig. 11 that it is desirable to operate at voltages at least 20–30 per cent above the stop band. The voltages at the stop band for the 1-inch period of Fig. 11 is computed from (9) to be 1,000 volts compared to 950 volts measured. Hence the agreement is within 5 per cent, which is within the experimental error of measuring the magnetic fields.

For the permanent magnet structures discussed here, it has been necessary for optimum focusing to use magnetic fields 50–100 per cent in excess of the value predicted in (9). However, in focusing with uniform fields it is also necessary to use similarly large fields. Thus, the initial experiments have resulted in focusing comparable to that obtained in the uniform magnetic field. It should be noted that in the uniform field case tubes have been built which focus satisfactorily with fields closer to the Brillouin values than mentioned above. However, it is expected that further refinement in the mechanics of the periodic structures will result in further improvement. For example in the structure of Fig. 9 it was possible to improve the focusing somewhat by rotating the individual magnet cells. This indicates the presence of transverse components of magnetic field on the axis which probably results from nonsymmetry in the pole pieces.

[8] J. J. Went, G. W. Rathenau, E. W. Gorter, and G. W. van Oosterhout, "Ferroxdure, a class of new permanent magnate materials," *Phillips Tech. Rev.*, vol. 13, p. 194; 1952.

RF Measurements

Since periodic magnetic-focusing structures introduce a ripple in the beam radius due to the periodicity of the magnets which is not characteristic of uniform field circuits, it was desirable to investigate the rf amplification properties of a typical traveling-wave tube under these circumstances. The tube tested was a medium power traveling-wave tube designed for operation at 3 cm. The input and output matches were constructed with standard $\frac{1}{2}\times 1$ inch waveguides since $\frac{1}{2}$ inch was the length of the magnetic lenses used in the periodic structure. Pole pieces were fitted into the waveguide and excited by external magnets, thus making an unbroken series of alternating lenses. With 5 ma of current focused to the collector with a beam voltage of 1,800 volts the rf gain of the tube at 3 cm was measured to be 22 db, with a power output of 200 milliwatts. This corresponds closely to operation in a uniform solenoid.

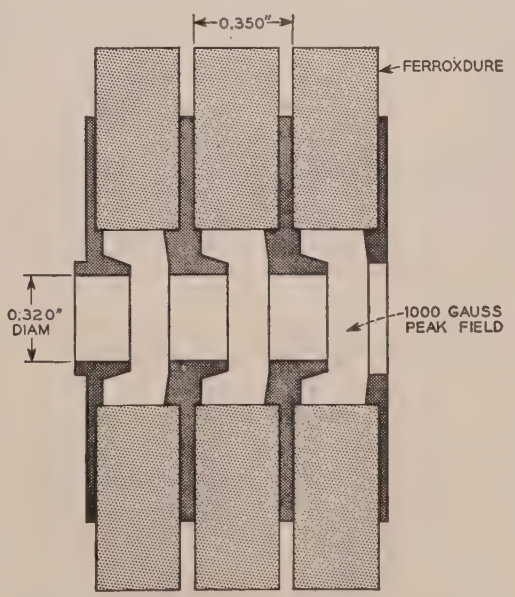

Fig. 12—Element of a periodic structure using Ferroxdure magnets to produce a field of 1,000 gauss with the period of 0.7 inch.

Dynamics

The equations of motion for the electron beam in the periodic magnetic field of Fig. 12 are derived with the following assumptions.

(a) The magnetic field is axially symmetric and uniform over the beam cross section. The longitudinal component can thus be written as,

$$B_z = B_0 \cos \frac{2\pi z}{L} \cdot \tag{6a}$$

(b) The electric field due to space charge acts only in a radial direction, i.e.,

$$E = E_r.$$

The Lagrangian for an electron in an electric and magnetic field is given by[9]

$$L = \frac{m}{2}(\dot{r}^2 + r^2\dot{\theta}^2 + \dot{z}^2) + e(v \cdot A) - eV \tag{10}$$

where A is the magnetic vector potential and V the electric potential. For the assumed conditions,

$$A = A_\theta = \frac{B_z r}{2} \cdot$$

For no magnetic field at the cathode one obtains from Busch's theorem[10]

$$\dot{\theta} = \frac{1}{2}\eta B_z \quad \text{where} \quad \eta = \left|\frac{e}{m}\right|. \tag{11}$$

Using (11), the expression for r derived from the Lagrangian is,

$$\ddot{r} + \frac{B_z^2\eta^2}{4} r - \eta \frac{\partial V}{\partial r} = 0. \tag{12}$$

For an electron at the edge of the beam,

$$\frac{\partial V}{\partial r} = \frac{r_0^2 \rho}{2\epsilon r}$$

where r_0 is the beam radius at the entrance of the focusing structure and ρ is the volume charge density. Equation (12) now becomes,

$$\ddot{r} + \frac{B_z^2\eta^2}{4} r - \frac{\eta r_0^2 \rho}{2\epsilon r} = 0. \tag{13}$$

It is convenient to make a change of variables. Let

B_0 = peak value of the longitudinal magnetic field at the axis

L = the magnet period (twice the distance between magnets)

$$\omega_\rho^2 = \frac{\rho\eta}{\epsilon} \qquad z = u_0 t$$

$$\omega_L = \frac{1}{2}\eta B_0 \qquad \sigma = \frac{r}{r_0}$$

$$\omega = \frac{2\pi u_0}{L}$$

[9] J. C. Slater and N. H. Frank, "Introduction to Theoretical Physics," McGraw-Hill Book Co., New York, N. Y., p. 77; 1933.
[10] Pierce, *loc. cit.*, p. 34.

where u_0 is the beam velocity which is assumed to be constant. Rewriting (13),

$$\frac{d^2\sigma}{dt^2} + \left(\frac{B\eta}{2}\right)^2 \sigma - \frac{\omega_\rho^2}{2}\frac{1}{\sigma} = 0. \tag{14}$$

Using the value of B given in (6a)

$$\ddot{\sigma} + \frac{1}{2}\left(\frac{\omega_L}{\omega}\right)^2 (1 + \cos 2T)\sigma - \frac{1}{2}\left(\frac{\omega_\rho}{\omega}\right)^2 \frac{1}{\sigma} = 0$$

where $T = \omega t$ and $\ddot{\sigma} = d^2\sigma/dT^2$.
For convenience let

$$\alpha = \frac{1}{2}\left(\frac{\omega_L}{\omega}\right)^2 \tag{15}$$

and

$$\beta = \frac{1}{2}\left(\frac{\omega_\rho}{\omega}\right)^2 \tag{16}$$

then

$$\ddot{\sigma} + \alpha(1 + \cos 2T)\sigma - \frac{\beta}{\sigma} = 0. \tag{17}$$

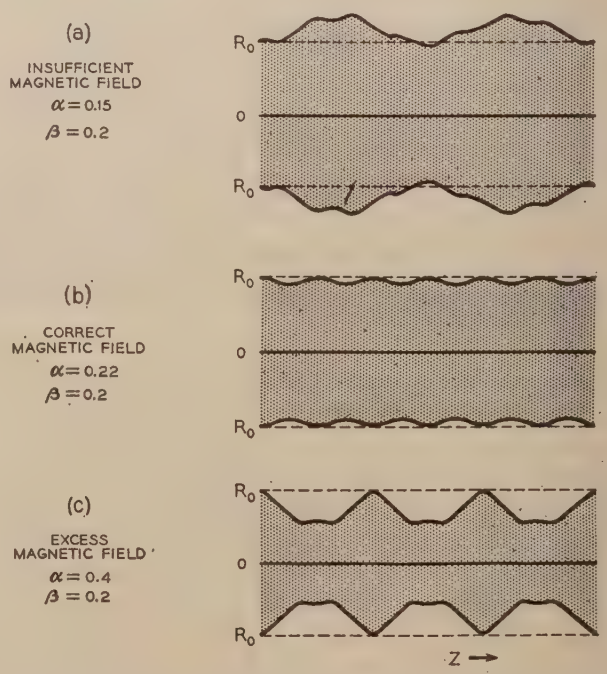

Fig. 13—Three typical analog-computer plots of beam radius as a function of distance, illustrating optimum flow conditions and deviations from this with an incorrect magnetic field. The small ripples in the three curves are associated with magnet spacing.

This nonlinear differential equation was solved with the aid of an analog computer for values of α and β of interest in the study of practical traveling-wave tubes. For practical reasons it was assumed that the electrons were injected with no radial velocity at the point $T = 0$, i.e., at the point where the magnetic field was a maximum. In principle, this condition should be attainable by proper positioning of the electron gun.

As might be expected, for a given set of beam parameters "optimum" focusing was obtained when the rms magnetic field strength was equal to the corresponding Brillouin field. By "optimum" it is inferred that the perturbations in the beam radius were a minimum. A complete set of beam contour curves were obtained from the analog computer with various values of α (the magnetic field parameter) and β (the space-charge parameter). Fig. 13 is typical of the type of beam contours encountered with a change in the magnetic field parameter α and a fixed value of β. Curve (b) is the so-called optimum point since this represents a minimum of beam ripple. Such curves were obtained for a wide range of the space-charge parameter β and from these curves it was possible to construct a universal curve for optimum focusing conditions as shown in Fig. 14.

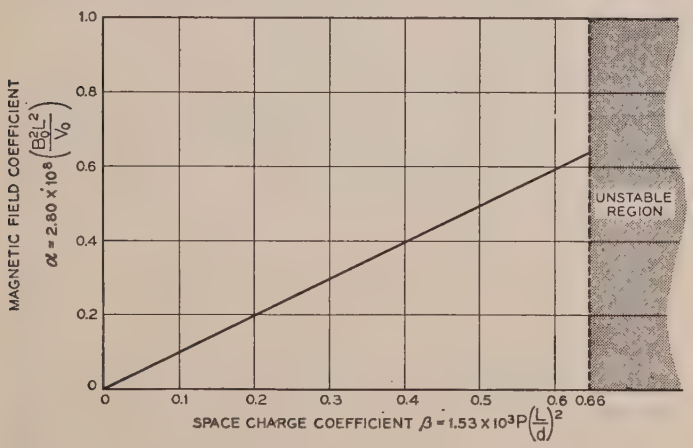

Fig. 14—Curve showing relation between beam parameters and magnetic field for the condition where the ripple on the beam boundary is a minimum.

In Fig. 14 it is evident optimum focusing occurs when

$$\alpha = \beta \tag{18}$$

or with (15) and (16) this can be written

$$I_0 = 1.45 \times 10^6 \, B_{\text{rms}}{}^2 a^2 V_0{}^{1/2} \tag{19}$$

where a = beam radius.

This will be recognized as the same relationship that exists with uniform field focusing in "Brillouin Flow" provided we associate the rms value of magnetic field in the periodic case with the dc field in the uniform field case and the average beam diameter in the periodic case with the beam diameter of the uniform case. Strictly speaking the required magnetic field applies only to the idealized beam that enters the structure with perfect parallel flow (no transverse velocities). In any physical system this is not the case and a higher magnetic field is required to hold the beam together within a prescribed radius.

In Fig. 15 the resultant minimum ripple is given for the optimum focusing conditions. This value will be exceeded in practice since a higher than optimum magnetic field is generally used. For high-percentage ripples it is difficult to obtain good transmission due to the rather violent fluctuations of the beam radius which results in the intensification of the slight aberrations inherent in any physical system. Although Figs. 14 and 15 represent an ideal system which can only be approached in practice, they do give valuable information concerning the ultimate that can be achieved in periodic focusing, as well as the correct relationship between the various beam parameters. These curves can be interpreted in practical cases by multiplying the indicated magnetic field by some empirical factor to account for transverse velocities.

In addition to the conditions of optimum focusing (17) must be used to yield information on the nature of the stop bands. Thus it is necessary to determine the conditions under which (17) yields diverging solutions

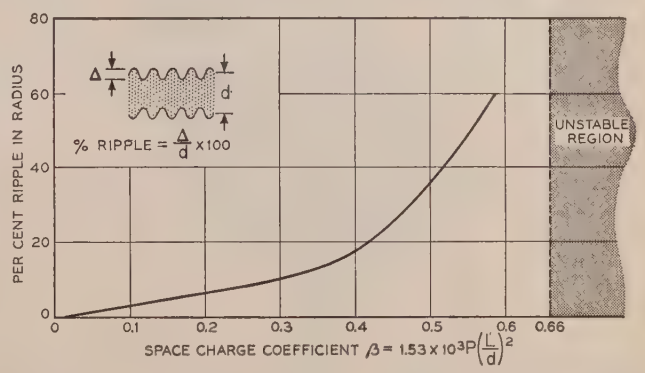

Fig. 15—Curve showing the minimum ripple which can be obtained on a beam which is focused with the optimum conditions shown in Fig. 15.

for σ. The equation can be simplified somewhat if it is noted that the last term on the left represents a diverging force. So, if the first two terms also give a diverging solution the entire solution of (17) must also diverge. By neglecting the space charge term (17) can be written

$$\ddot{\sigma} + \alpha(1 + \cos 2T)\sigma = 0. \tag{20}$$

This is a form of Mathieu's differential equation[11] with arbitrary constants specified, and is quite accurate for locating the position of the so-called stop bands. In (20) the nature of the solution for σ as a function of T (or z) depends upon the value of the constant α (which is always positive). For certain ranges of α, σ is periodic or quasi periodic and bounded while for other values of α, σ increases without limit thus indicating an unstable condition of electronic flow. These unstable regions can easily be found from the stability plot of (20)[11] given in Fig. 16, which applies to the general form of Mathieu's equation whereas the straight line on this plot applies to the particular form given in (17). As α is increased from zero along this straight line the solution for σ passes through successive regions of stability (pass band) and instability (stop band). It should be noted that this plot merely indicates the regions of α for which

[11] N. W. MacLachlan, "Theory and Application of Mathieu Functions," Oxford University Press, New York, N. Y.; 1947.

instability occurs since space charge has been ignored in obtaining this diagram. Therefore, the analog computer solutions are necessary to show that the stable regions indicated on Fig. 16 truly represent stable electron flow in the presence of space charge. Consequently, the Mathieu diagram only gives information concerning the location of the stop bands and is not intended as a guide for the conditions of good focusing.

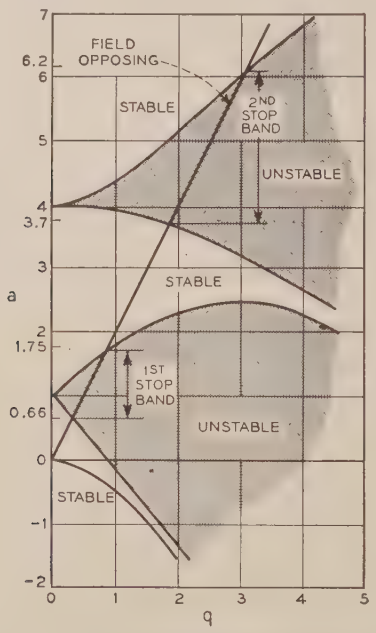

Fig. 16—Mathieu Stability Plot which gives the theoretical values of α (the magnetic field parameter) for the stop bands and pass bands. The general form of Mathieu's equation is,

$$\frac{d^2\sigma}{dT^2}+(a+2q\cos 2T)\sigma=0$$

and for the periodic structure $a=\alpha$, $q=\alpha/2$ yielding the straight line shown on the diagram.

As indicated in Fig. 16 instability and hence the beginning of the stop band occurs for

$$\alpha = 0.66$$

or with (15)

$$\frac{\eta B_0^2 L^2}{V_0} = 418. \tag{9}$$

The extent of the stop band together with the location of the second stop band is shown in Fig. 16. It is from this plot that the information for the theoretical stop bands is obtained for Figs. 4 and 6.

Design Considerations

In considering the design of a periodic-structure it is instructive first to note the restrictions imposed by the stop-band phenomenon as expressed in (9):

$$\frac{\eta B_0^2 L^2}{V_0} < 418. \tag{9}$$

To achieve a balance between space-charge forces and the inward focusing force of the magnetic field, it is necessary that (see Fig. 14)

$$\alpha = \beta. \tag{18}$$

Using the definition (15) and (16) for α and β, (18) becomes

$$\frac{\eta B_0^2 L^2}{V_0} = \frac{16}{\pi}\frac{1}{\epsilon\eta^{1/2}} P\left(\frac{L}{d}\right)^2.$$

Substituting (9) for the left-hand side of the above equation one obtains

$$P\left(\frac{L}{d}\right)^2 < 4.35 \times 10^{-4}. \tag{21}$$

This simple relationship gives the maximum value of $P(L/d)^2$ attainable with periodic system if it were possible to operate at the edge of the stop band. In practice, however, this value is less than that given in (21) because of two factors. First, for good focusing the operating voltage must be about 20 per cent in excess of the voltage at the stop band and secondly, it is necessary to use magnetic fields in excess of the Brillouin value by 50–100 per cent. These two factors combine to reduce the right-hand side of (21) to give

$$P\left(\frac{L}{d}\right)^2 \leqq 10^{-4}. \tag{22}$$

This condition must be satisfied if good focusing is to be realized. Since perveance and beam diameter are probably the most fundamental constants of any beam such a relationship is very useful for determining the limitations of a given periodic structure.

If one considers the limiting value of L/d which can be realized in a practical structure then it is possible to describe a rough limit to the perveance that can be employed.

With axial symmetry the longitudinal magnetic field varies radially as

$$I_0\left(\frac{2\pi r}{L}\right)$$

and, therefore, the ratio of the field at the inner diameter of the magnets to the field at the axis is given by Ratio $= I_0(\pi D/L)$ where D is the inner diameter of the magnets. In general D can be no less than three times the beam diameter. Let

$$D = 3d, \qquad d = \text{beam diameter.}$$

Hence,

$$\text{Ratio} = I_0\left(\frac{3\pi d}{L}\right).$$

Since there is a limit to the coercive force that can be achieved with permanent magnets, it is reasonable to ascribe an upper limit to the above ratio. Arbitrarily let

$$I_0\left(\frac{3\pi d}{L}\right) \leqq 3.$$

Consequently

$$\frac{3\pi d}{L} \leqq 2.5$$

or

$$\left(\frac{L}{d}\right)^2 \geqq 16.$$

Referring to (22) one can now write an approximate limit to the perveance that can be satisfactorily focused in a periodic structure.

$$P \max \cong 6 \times 10^{-6}.$$

Fortunately, present-day beams typically employ perveances of the order of unity, which is well within the above practical limit.

Conclusion

This paper describes a system of permanent magnet focusing which realizes such a saving in weight and size that most beam-type vacuum tubes can be focused with a structure weighing but a few pounds. This is to be compared with previous methods which required magnets which weighed more than 40 pounds. Also, the new system makes practical the design of much longer beam-type tubes than could previously be considered because of the cumbersome focusing structure required.

The limits of practicability of this system include beam perveances in excess of what is commonly used today with magnetic field strengths which are realizable with commercially available permanent magnet material. Although there is a stop-band phenomenon to consider, design procedures have been established which determine a suitable structure once the perveance and beam diameter are specified.

The experimental evidence presented shows very good agreement with the analytical results, and in addition demonstrates the feasibility of obtaining satisfactory focusing (>99 per cent transmission) with practical structures. It is felt that one of the prime objections to long beam-type amplifiers, namely the bulky magnets or power consuming solenoids, has been eliminated with the successful demonstration of permanent magnet periodic focusing.

Acknowledgment

The authors gratefully acknowledge the many cogent suggestions given by C. C. Cutler and R. Kompfner, and their encouragement throughout the course of this work. Also, the fine constructive efforts of L. E. Cheesman, N. E. Cram, and P. Hannes are greatly appreciated.

Appendix A

As approximations to the fields produced by a uniform magnet and by a periodic magnetic structure, consider Fig. 17. Fig. 17(a) shows a segment of length L (period $=2L$) of an infinitely long sinusoidal distribution of magnetic field. This, it is assumed, is approximately the distribution of flux from a bar magnet of effective length L at a radial distance fom the structure of $r=Na$. Fig. 17(b) shows a segment of length L from a sinusoidal distribution of field where the periodicity has been increased by N, i.e., $Nx=L$. This is approximately the magnetic field distribution from a periodic magnetic structure at a distance $r=a$. It is assumed that the two fields are derived from magnets or lenses having similar geometries, simply differing by a scale factor N. Hence the field in Fig. 17(a) at a distance $r=Na$ will equal that of Fig. 17(b) at a distance $r=a$.

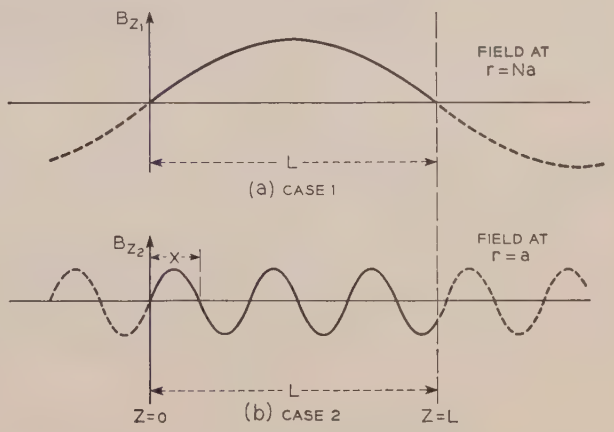

Fig. 17—Approximate form of magnetic fields for (a) a long bar magnet of length L, and (b) for a periodic magnetic structure of length L and periodicity N.

Since in a region of no current B_z in cylindrical coordinates is a solution of Laplace's equation we know that the fields for the two cases may be written

$$B_{z_1}(r, z) = C_1 \frac{\pi}{L} K_0\left(\frac{\pi r}{L}\right) \sin \frac{\pi z}{L} \text{ for case 1, and}$$

$$B_{z_2}(r, z) = C_2 \frac{\pi}{x} K_0\left(\frac{\pi r}{x}\right) \sin \frac{\pi z}{x} \text{ for case 2.}$$

From the above assumptions then

$$B_{z_1}(Na, L/2) = C_1 \frac{\pi}{L} K_0\left(\frac{\pi Na}{L}\right) = B_{z_2}\left(a, \frac{x}{2}\right)$$

$$= C_2 \frac{\pi}{x} K_0\left(\frac{\pi a}{x}\right)$$

or

$$\frac{C_1}{L} = \frac{C_2}{x}.$$

We are interested in the stored magnetic energy which in the above two cases leads to the following expressions. In case 1,

$$W_1 = \int_{Na}^{\infty} B_{z_1}{}^2 dA$$

$$= \int_{Na}^{\infty} B_{z_1}^2 r dr \quad \text{where } W_1 \text{ is proportional to energy}$$

$$= C_1^2 \int_{\pi(Na/L)}^{\infty} \nu K_0^2(\nu) d\nu$$

and similarly for case 2,

$$W_2 = C_2^2 \int_{\pi a/x}^{\infty} \nu K_0^2(\nu) d\nu.$$

Since $L = Nx$ the two integrals are equal and thus

$$\frac{W_2}{W_1} = \frac{C_2^2}{C_1^2} = \frac{1}{N^2}.$$

Thus the external stored energy in the periodic structure is $1/N^2$ times that of the uniform field magnet for this idealized case.

Appendix B

Derivation of the Magnetic Field at the Center of a Periodic Focusing Structure

In deriving the dynamic equations describing the motion of electrons in a beam focused with a periodic magnetic structure, we assume the magnetic field is constant across the beam diameter and is of the form

$$B_z = B_0 \cos \frac{2\pi z}{L}.$$

In actual structures the magnetic field does have a variation with r and also contains higher harmonics as shown in the following derivation. Fig. 18 shows a plot of the field and a sketch of the structure from which the field is derived.

In a current-free region B_z satisfies Laplace's equation. If cylindrical co-ordinates are used, this leads to solutions of the form:

$$B(z, r) = \sum_{1}^{\infty} C_n I_0\left(\frac{2n\pi r}{L}\right) \cos \frac{2n\pi z}{L}.$$

Evaluating C_n, we get,

$$C_n = \frac{4}{\delta I_0\left(\frac{2n\pi R}{L}\right)} \int_0^{\frac{\delta}{2}} B_{\text{gap}} \cos \frac{2n\pi z}{L} dz$$

where R is the radius of the whole through the pole piece.

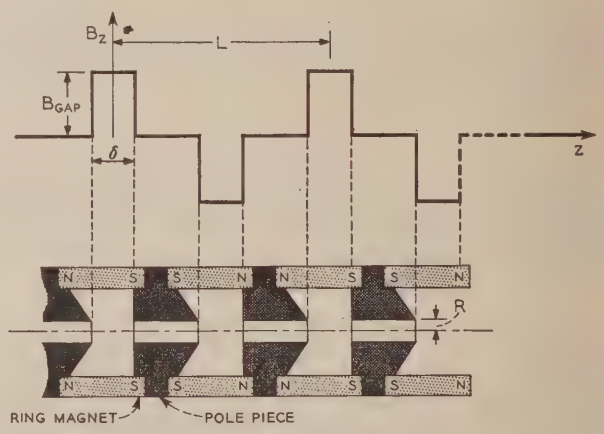

Fig. 18—Magnetic field strength at the edge of the pole pieces for a periodic magnetic structure.

$$C_n = \frac{4B_{\text{gap}}}{I_0\left(\frac{2n\pi R}{L}\right)} \frac{\sin \frac{n\pi\delta}{L}}{n\pi} \qquad n = 1, 3, 5, \cdots.$$

In the present case $R = 0.160$ inch, $L = 1.000$ inch, and $\delta = 0.150$ inch and the expression for the magnetic field on the axis of the structure is

$$B(z, 0) = B_{\text{gap}}\left[.456 \cos \frac{2\pi z}{L} + .0711 \cos \frac{6\pi z}{L} + .0066 \cos \frac{10\pi z}{L} + \cdots\right].$$

Antenna and Receiver Measurements by Solar and Cosmic Noise*

JULES AARONS†, ASSOCIATE, IRE

Summary—**Utilizing the most intense of the several hundred celestial sources of radio frequency energy, and known data about the effective diameter of the sun at various ranges of radio frequencies, a technique for plotting the directional characteristics of large antennas is outlined. Over-all system sensitivities (receiver, antenna, and transmission line) are checked by using values already obtained for sky temperatures. The general receiver characteristics necessary for such measurements are outlined. Patterns of antennas are illustrated and their analysis evaluated.**

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Background

Basic research in astronomy, within the past five years, has used two new tools extensively, the low-noise stable radio receiver and the highly directive antenna. Both the sun and the cosmic noise sources have been explored extensively with these tools.

The greatest amount of energy, as recorded by observations, from both of these sources is in the 18-mc to 300-mc region. The cosmic sources have been explored almost wholly in this frequency region. The sun maintains its irregular or disturbed component basically

in the low frequencies and has a rather quiet radiation at the microwave end of the radio spectrum. Receiver equipment designed to detect these sources is exceedingly stable.

Several workers in the field now use previously established facts about the sun and the cosmic sources to calibrate their antennas. For other antennas which have large effective areas, the cosmic and solar sources might prove to be useful. These sources are never, of course, in the near field; they can be located precisely at all times using astronomical tables; they present a means for checking the over-all system sensitivity over long periods of time in case of deterioration of equipment. This technique of sensitivity calibration would aid in maintaining similar sets of identical equipment at constant levels by providing a "primary" noise-source standard.

Receiver Principles Used in the Design of Radio Telescopes

For these measurements three basic factors are important.

1. The most vital consideration in over-all design of the receiver is its stability. Australian scientists[1] have designed 100-mc equipment which achieves stability of one part in 10,000 for the B plus voltage, and one part in 6,000 for the filament voltage. They, therefore, were enabled to obtain changes in sky temperatures of 0.3 degrees K. After detection bucking voltage is set up to rid the receiver of the steady-state noise. A milliammeter or microammeter records the signal.

2. In addition to constant supply voltages and dc-filament voltage (in some cases), these low-noise receiving systems use long time constants in their recorders.

The minimum signal which can be detected is proportional to the square root of the bandwidth-time constant product.

$$\frac{P_{\min}}{P_n} = \frac{1}{\sqrt{BC}} \qquad (1)$$

$P_{\min}$ = minimum detectable signal power
P_n = power generated in the receiver and delivered to the receiver
B = bandwidth in cycles per second
C = time constant.

From the analysis suggested by Dicke,[2] the output is the sum of a large number of successive random pulses, each of average duration $1/B$. In a time interval ΔT (the time constant of the output system) $B\Delta T$ pulses are received. Solar noise bursts have been of the order of one-half second or longer, and time constants have been maintained in that range. Bandwidths have varied, although in many cases bandwidths have been below radar or television bandwidth requirements.

3. Atmospheric noise, ignition noise, etc., in the vicinity of the antennas can be an important factor.

[1] C. W. Allen and C. S. Gum, "Survey of galactic radio-noise at 200 mc," *Aus. Jour. Sci. Res. A.*, vol. 3, p. 224; 1950.
[2] R. H. Dicke, "The measurement of thermal radiation at microwave frequencies," *Rev. Sci. Inst.*, vol. 17, p. 268; 1946.

The Sources

The Sun

The sun is a variable source of radiation emitting on top of its normal "quiet" level signal an enhanced base or long-term level plus shortlived bursts and outbursts.

Angular width and general activity vary on different frequencies (Fig. 1). The width of the photosphere, the optical width of 31 minutes, 59.3 seconds, approaches the radio-frequency width at 1 cm. However, at 90 per cent of the emergent intensity of its central ray at 60 mc the sun appears as a source of 47 minutes. The exact details of the width are still in doubt. New data taken by P. A. O'Brien at the Cavendish Laboratory in Cambridge, England would tend to increase the width considerably (Mon. Not. R.A.S. in press).

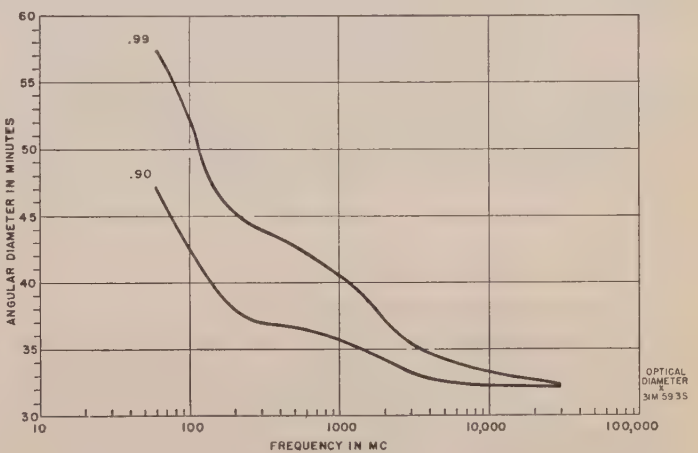

Fig. 1—Angular diameter of the sun at radio frequencies (Smerd).[3]

The general activity at the low end of the radio frequency spectrum is much greater than in the microwave region. For the former, during a disturbed day, several hundred bursts may occur within the space of one hour at 200 mc.

At centimeter wavelengths the picture is different. "Sudden increases of radiation at 3.18 cm caused by disturbed conditions in the sun, were found to be rare."[4]

These observations make the measurements outlined in this paper more feasible for the sun in the microwave range. At the low activity portion of the sunspot cycle measurements in both ranges are quite possible.

Cosmic Noise

Much exploratory work has been done in the radio frequency band of 18 mc to 400 mc in cosmic noise sources. At 1,200 mc, Cygnus A, the most "intense" point source is rather diffuse.[5] At 3,000 mc no point sources have been located. Therefore, cosmic noise is feasible only for radio-frequency measurements at the low end of the spectrum.

[3] S. F. Smerd, "Radio frequency radiation from the quiet sun," *Aus. Jour. Sci. Res. A.*, vol. 3, p. 34; 1950.
[4] H. C. Minnett and N. R. Labrum, "Solar radiation at a wavelength of 3.18 cm," *Aus. Jour. Sci. Res. A.*, vol. 3, p. 60; 1950.
[5] J. H. Piddington and H. C. Minnett, "Radio frequency radiation from the constellation of Cygnus," *Aus. Jour. Sci. Res. A.*, vol. 5, p. 17; 1952.

Several hundred of these sources of radio-frequency energy have been located. The most intense ones, the sources suitable for our measurements, range from one minute to thirty minutes of arc in source width. At a frequency of 100 mc the intensity varies downward from 128×10^{-24} watts/cycles/second/steradian.[6]

A few of the sources, viewed at a low angle of elevation above the medium have fluctuating components due to ionospheric twinkling. One of these, Cyngus A, had maximum variations of about 10 per cent at 200 mc.

Antenna Patterns' Measurement Technique

Calculation of Beam Angle

Taking measurements with this technique must be done with the aid of radio astronomy celestial maps. At meter wavelengths the cosmic sources cannot be used if the milky way is in the background. The same holds true for the sun. Those sources which are near the galactic center will seem to stretch out the antenna pattern. Since we now have a gaussian distribution to contend with, the corrected antenna angle is

$$\theta_t = \sqrt{\theta_m^2 - \theta_s^2}\,. \qquad (2)$$

θ_t = the true antenna angle
θ_m = the measured antenna angle
θ_s = the gaussian curve of the contours through which the antenna is sweeping.

For the case of the point source which is sharp and above the intensity of neighboring contours, the particular curve must be considered in subtracting source size from measured size in order to obtain the true antenna angle. In the case of the sun, the intensity may be considered rectangular for a good approximation at the shorter wavelengths but only a fair approximation at the longer wavelengths. The sides are sloping, but except for an element of brightening at the edges (at microwave frequencies) the top is flat. In this case, then, the angle subtended by the sun may be subtracted from the measured antenna angle and the true angle obtained.

If the stars are considered as points on the inner surface of a rotating sphere, there lies within, a stationary globe, the earth. The axes of the two spheres are inclined towards one another; they are not parallel. This failure to "line up" will produce an incorrect pattern for an antenna which is fixed in azimuth and the sources allowed to sweep across it. These patterns can be measured reasonably well if the beam angle is small and the source, therefore, takes only a small amount of time to cross the antenna pattern.

After taking the pattern, the correction factor of sky contours must be applied. Isothermal lines at various frequencies have roughly the same shape. At 18.3 mc the contours are similar to 100 mc.[7] A field technique which might be put to use is to make measurements on a particular antenna under laboratory-controlled conditions, reducing losses to a minimum, and achieving an adequate test pattern. If an antenna test pattern is then obtained of the sun or a cosmic source, a standard can be set up for pattern and sensitivity comparison.

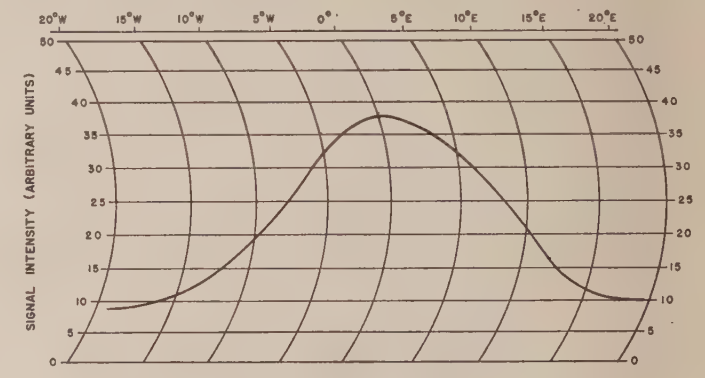

Fig. 2—Sweep across sun on October 16, 1952; main lobe 200-mc antenna.

Measurement at Cornell University

Examples of the measurements of antenna beam angles are simple to analyze for a first approximation to the pattern. Fig. 2 illustrates a swing of the antenna across the sun on October 16, 1952, on 200 mc, by the Cornell University solar noise equipment.[8] A peak value of 37.5 units was subtracted from a background value of 7 units. Therefore, the half-power points are 18 degrees apart. If this figure is reduced by the value of the width of the source (38 minutes at 90 per cent down from the emergent intensity) the resultant value is 17 degrees 22 minutes, assuming the solar source is a rectangle.

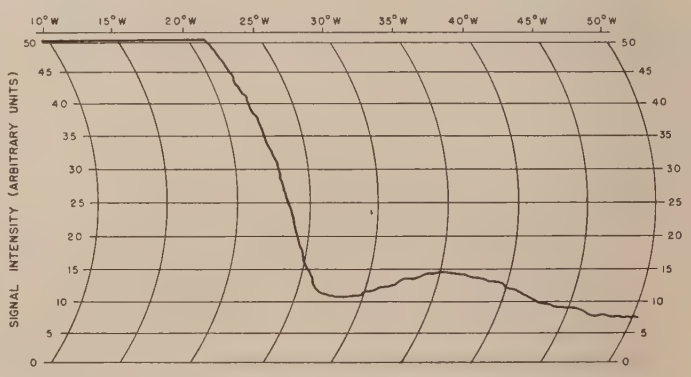

Fig. 3—Sweep across sun on October 16, 1952; side lobe 200-mc antenna.

Fig. 3 illustrates measurements taken of the side lobe patterns. The technique used was to increase the gain of the equipment and concentrate on the side lobe.

Patterns can be taken with an antenna fixed in azimuth. In one minute of time, 15 minutes of arc times the cos of the declination is traversed.

[6] G. J. Stanley and O. B. Slee, "Galactic radiation at radio frequencies," *Aus. Jour. Sci. Res. A.*, vol. 3, p. 234; 1950.

[7] C. A. Shain, "Galactic radiation at 18.3 mc," *Aus. Jour. Sci. Res. A.*, vol. 4, p. 258; 1951.

[8] All measurements used in this paper were made by the Cornell University Radio Astronomy Laboratory. S. M. Colbert made the measurements under the general supervision of Dr. C. Burrows.

Sensitivity

Over-all Sensitivity Measurements

A possibility exists for measuring changes in the sensitivity of the receiving equipment from the antenna to the recorder. Since the essential criterion of the proposed system will rely on its ability to distinguish between noise levels, let us first consider some of the noise equations.

The available noise from a resistor[9] is:

$$P_r = \frac{e_n^2}{4R} = KT_aB \tag{3}$$

P_r = available noise power in watts
e_n = RMS voltage
K = Boltzmann's constant (1.37×10^{-23} watt-sec/degree K)
T_a = "ambient" temperature in degrees Kelvin usually taken at Friis' standard of 290 degrees K
$KT_a \approx 4\times10^{-21}$ watt-seconds.

The noise figure (F) of a receiver under test is the ratio of the actual noise power output to the noise power output of an ideal network.

The component of the noise figure due to internally generated noise (all noise other than the input resistor noise) is equal to $F-1$.

In order to detect a small amount of energy being radiated from external sources, the power absorbed by the antenna and delivered to the receiver must be of the same magnitude or larger than the fluctuations in the internally generated noise of the receiver. If the internally generated noise were perfectly steady in its amplitude, it would be theoretically possible to detect any increase caused by external power being absorbed. The variable component of the noise is a function of the bandwidth and the time constant of the circuit, provided the supplies are stable and no outside interference takes place.

Equation (1) shows that for a bandwidth of one megacycle and a time constant of one second in the recording apparatus it is theoretically possible to detect signals of the order of 1/1,000 of the internal noise.

$$P_{\min} = \frac{KT_aB(F-1)}{\sqrt{BC}} \tag{4}$$

where $KT_aB(F-1)$ or P_n is the power generated in the receiver and delivered to the receiver.

The absorption of energy by the antenna for the case of the homogeneous areas of emission (Θ energy source $>\Theta$ antenna is:

$$P_{abs} = KT_cB. \tag{5}$$

If the antenna absorbs just enough energy from the radio noise source so that a small signal can be seen above the theoretical noise minimum of the receiver $P_{\min}$ can be set equal to P_{abs}. Eliminating the constants we obtain

$$T_c = \frac{T_a(F-1)}{\sqrt{BC}}. \tag{6}$$

P_{abs} = power in watts absorbed by the radiation resistance of the antenna and delivered to the receiver

T_c = sky temperature.

The measurement of $P_{abs\ min}$ (the minimum amount of power which the antenna must absorb before the power output rises above the internal noise power), and with the measurement of the noise figure, the bandwidth and the time constant, a new criterion for receiver evaluation is available, i.e., that of stability.

The ratio of measured power absorbed (with a perceptible change of level noted) to theoretical power ($P_{\min}$) yields the stability factor M. A value of one, just as in the case of the noise figure, is obtained when the unit is perfectly stable.

$$\frac{T_{cm}}{T_{ct}} = M. \tag{7}$$

T_{cm} is the measured temperature of the homogeneous sky source which yields a level which can be distinguished above the fluctuations of receiver noise. By setting P_{abs} equal to $P_{\min}$ in (6) we had actually obtained:

$$T_{ct} = \frac{T_a(F-1)}{\sqrt{B_c}} \tag{8}$$

Where T_{ct} is the theoretical temperature of the antenna plus the receiver.

A change of the stability factor M as evidenced by inability to detect certain radiating areas of the sky will be caused by an increase in F, a change in the system stability due to power supply fluctuations (assuming bandwidth and time constant are fixed), or a deterioration in the antenna system (transmission line, mismatching, etc.). Antenna-gain considerations do not enter into the equation.

In the case of receiving systems dependent on comparison techniques, the stability can be an important consideration. A simple means of detecting increase in the noise factor or other parameters of the receiver is the technique of pointing the antenna at the sky.

As a routine check on the stability factor (M), one technique is to attempt to distinguish between two sets of contours or two areas. Instead of using the absolute power from one area, the difference in power between two areas with a temperature difference of ΔT can be utilized. The receiver should be able to distinguish between two areas which have

$$\Delta T_{cm} = \Delta T_{ct}M. \tag{9}$$

[9] H. T. Friis, "Noise figures of radio receivers," Proc. I.R.E., vol. 32, p. 419; July, 1944.

Fig. 4 represents Cornell University's 200-mc stabilization data which is taken daily. Here a reference resistor at 290 degrees represents a zero point. The levels

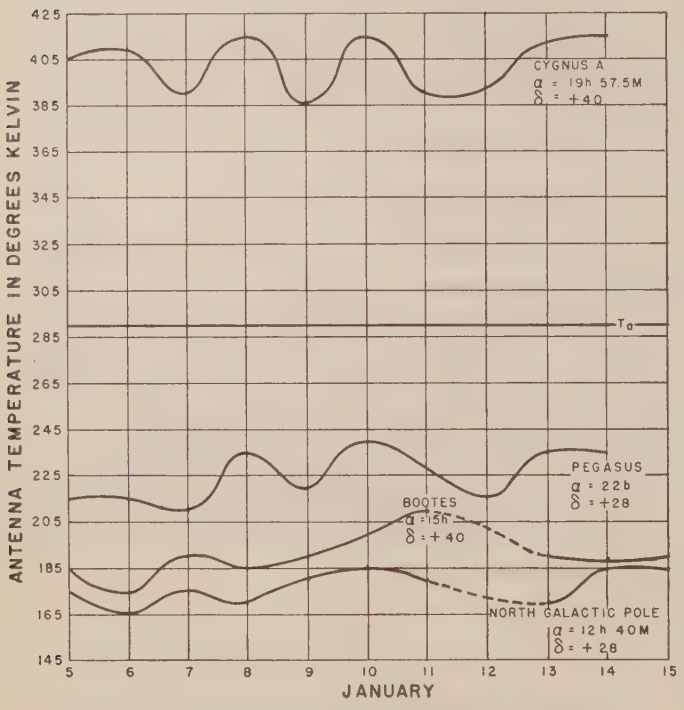

Fig. 4—Stability measurements by means of sky noise January 5–15, 1953.

above and below this are for portions of the sky centered at the location given. Utilizing (8),

$F = 5.9$

$B = 5$ mc

$C = \frac{1}{2}$ second

a ΔT_{ct} of 0.9 degrees K is obtained.

In one radio astronomy project designed for somewhat different purposes, an M of 2 was obtained. For normally stabilized radio receivers it is estimated that an M of 100 would result.

Θ *Antenna* $>$ Θ *Energy Sources*

Summing the Rayleigh Jeans radiation law into total emissive power and converting the equation into frequency form, the power absorbed by an antenna with the energy source smaller than the beam angle is[10]

$$P_{abs} = \frac{G_0 K \Omega_s T_s B}{4\pi} \tag{10}$$

Ω_s = solid angle subtended by the optical disk of the sun

T_s = the equivalent black body temperature of the sun in degrees Kelvin

G = antenna gain in the maximum direction.

Fig. 5 graphs the value of $K\Omega_s T_s$ for the sun as a function of frequency.

Θ *Energy Source* $>$ Θ *Antenna*

If the entire area of the antenna's main lobe is homogeneously filled with radiation, directional characteristics of the antenna do not enter into the calculations.

$$P_{abs} = KTB.$$

For a large antenna, the measurements involving portions of the sky which emit a uniform signal level can be made to obtain a field intensity in microvolts per

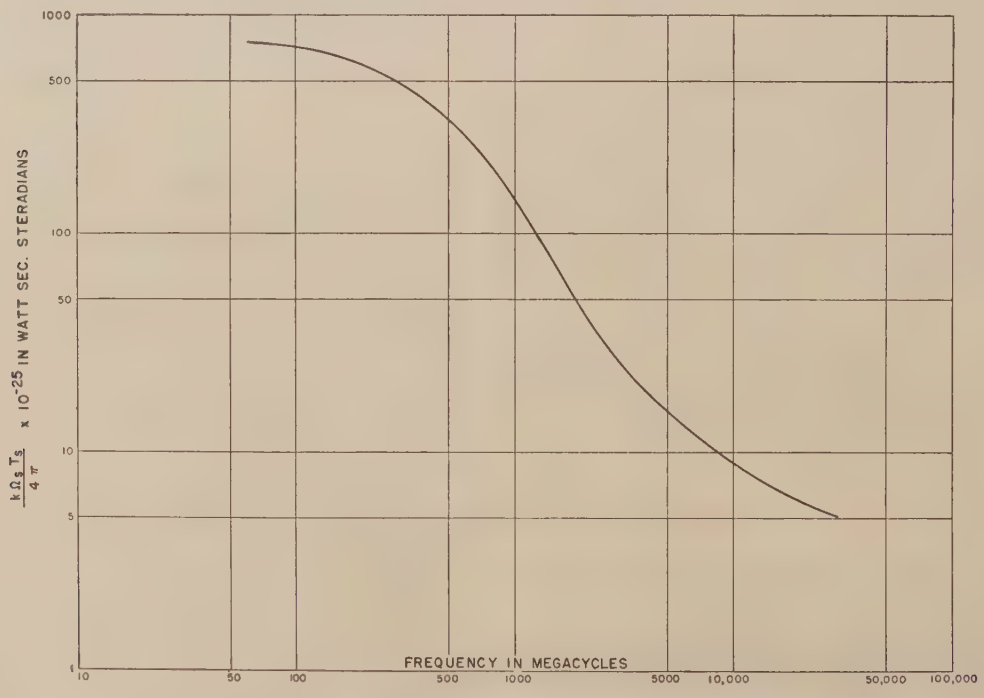

Fig. 5—Effective solar temperature as a function of frequency (Hagen).[11]

[10] H. C. Van de Hulst, "A Course in Radio Astronomy," Leiden Observatory, Holland; 1951.

[11] J. P. Hagen, "Temperature gradient in the sun's atmosphere measured at radio frequencies," *Astron. Jour.*, vol. 113, p. 547; 1951.

meter. References 1 and 7 contain maps plotted for various frequency regions.

No contours are available for the microwave regions due to the low-temperature differences involved. For the galaxy, the effective temperature varies approximately as λ^2. At 4.7 meters, Hey, Parsons and Phillips (1948) obtained a maximum of 18,000 degrees K.[12] For the same region, 30 degrees K was obtained from Reber's data on 60 cm.

Conclusion

The possibilities for the use of this technique lie as a field tool for checking patterns and sensitivity. The inaccuracies involved in evaluating background radiation, contours, and solar energy distribution are such that a very accurate laboratory field set-up would be a better tool. However, for a check on field instruments as to their noise level, stability, antenna pattern, and terrain fluctuations, the technique can be of use.

[12] J. S. Hey, S. S. Parsons, and J. W. Phillips, "An investigation of galactic radiation in the radio spectrum," *Proc. Royal Society*, vol. 192, p. 425; 1948.

Acknowledgment

I would like to thank S. M. Colbert of the Cornell University Radio Astronomy Laboratory for his assistance in obtaining the data for this report. I would also like to thank Prof. Charles Burrows for his co-operation with the work. Miss Martha Henissart, of Wentworth Institute, assisted in preparing the graphs and the data and in many of the problems of interpretation.

IRE Standards on American Recommended Practice for Volume Measurements of Electrical Speech and Program Waves, 1953*

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[1] Chairman, 50–52.

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C. A. Cady, *Chairman*, 50–52

L. B. Cooke, 50–52	W. W. Dean, 50–52	J. W. Miller, 50–52	R. B. Monroe, 50–52

1. General

These standards apply to the methods of and a device for measuring the dynamic magnitude of complex audio-frequency electrical waves such as occur in speech and music.

The measurement of the complex and nonperiodic waves encountered in electrical communication cannot be expressed in simple fashion in the ordinary electrical terms of current, voltage, and power. The concept of "volume" furnishes a practical method of great utility

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to the communications engineer for assigning a numerical value to the magnitude of electrical speech and program waves.

Volumes are read by noting the more extreme meter deflections of a device known as a volume indicator. Since the response of the meter of such an instrument to the rapidly varying waves is greatly dependent upon its dynamic characteristics, a standard for volume measurements must, therefore, include a specification of these characteristics.

This standard uses the term "vu" to express volume in terms of vu above or below a specified reference level, defined herein.

2. Definitions

The following terms apply specifically to the standard volume indicator and its use.

2.1 Volume. The magnitude of a complex audio-frequency wave in an electrical circuit as measured on a standard volume indicator.

2.2 Standard Volume Indicator. A device for the indication of volume, and having the characteristics described in Section 4.

Note—A standard volume indicator consists of at least two parts:

(a) A meter.
(b) An attenuator (adjustable loss) or pad (fixed loss).

2.3 vu (pronounced "vee-you" and customarily written with lower case letters). A quantitative expression for volume in an electric circuit.

Note 1—The volume in vu is numerically equal to the number of db which expresses the ratio of the magnitude of the waves to the magnitude of reference volume.

Note 2—The term vu should not be used to express results of measurements of complex waves made with devices having characteristics differing from those of the standard volume indicator.

2.4 Reference Volume. The level which gives a reading of 0 vu on a standard volume indicator.

Note 1—The methods of reading and calibration are described in Section 4.

Note 2—The "reading of 0 vu" is the algebraic sum of the meter and attenuator readings on the standard volume indicator.

2.5 Reference Deflection. The deflection to the meter-scale point marked 0 vu, 100, or both.

Note 1—This is the deflection at which the meter should be used.

3. Standard Volume Indicator

Volume is measured by means of a volume indicator. This device must conform to the following specifications and must be used in the manner described below.

3.1 Circuit Impedance. The magnitude of the circuit impedance across which the volume indicator is intended to be connected is designated herein as "R" ohms, which is the resistive impedance used for calibration. When so used, the requirements of Sections 4.2 and 4.5 must be met.

Note 1—The impedance presented to the volume indicator by a circuit impedance of "R" ohms is nominally $R/2$ ohms; resulting from the parallel combination of the "R" ohm source and the "R" ohm load.

3.2 Dynamic Characteristics. If a sinusoidal voltage between 35 and 10,000 cycles per second, of such amplitude as to give reference deflection under steady-state conditions is suddenly applied, the meter pointer shall reach 99 per cent of reference deflection in 0.3 second, ± 10 per cent, and shall then overswing reference deflection by at least 1.0 per cent and not more than 1.5 per cent. The time required for the meter pointer to reach its position of rest on the removal of the sinusoidal voltage shall not be greatly different from the time of response.

3.3 Response-Versus-Frequency Characteristic. The response of the volume indicator shall not depart from that at 1,000 cycles per second by more than 0.2 decibel between 35 and 10,000 cycles per second nor more than 0.5 decibel between 25 and 16,000 cycles per second.

3.4 Response to Complex Waves. The response to complex waves of such amplitude as to give reference deflection when read, as described in 4.9, shall be that equivalent to the response with a direct-current meter and a rectifier, the exponent of whose characteristic is 1.2 ± 0.2.

3.5 Reversibility. The response when measuring unsymmetrical waves must be independent of the poling of the volume indicator. Such a characteristic may be obtained by the use of a direct-current meter in conjunction with a full-wave rectifier.

3.6 Graduation of Meter Scale. The point of reference deflection shall be definitely indicated in some suitable manner. The remainder of the scale shall be graduated in vu above and below reference deflection. (See also 4.4)

3.7 Attenuator Marking. The attenuator shall be marked in vu.

3.8 Calibration. A correctly calibrated volume indicator with its attenuator set at 0 vu will give reference deflection when connected to a source of a sinusoidal voltage adjusted to develop 1 milliwatt in a resistance of R ohms, or with the attenuator set at n vu when the calibrating voltage is adjusted to develop a power n decibels above 1 milliwatt.

3.9 Method of Reading Volume Indicator. The reading is determined by the greatest deflections occurring in a period of about a minute for program waves, or a shorter period (e.g., 5 to 10 seconds) for message telephone speech waves, excluding not more than one or

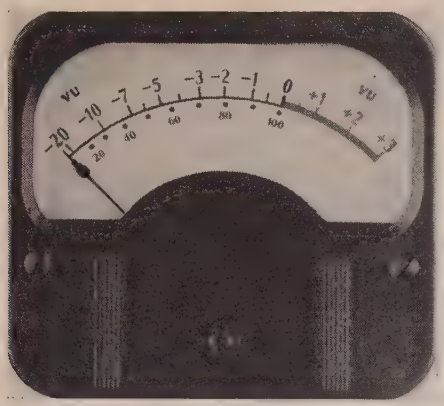

Fig. 1

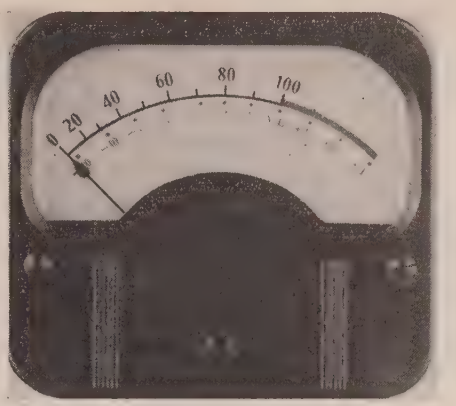

Fig. 2

two occasional deflections of unusual amplitude. The volume indicator is usually connected across the circuit at a point where the impedance is R ohms and the attenuator is adjusted until the deflections, read as described above, just reach the scale point corresponding to reference deflection. The volume in vu is determined by the markings on the attenuator at the setting thus obtained. If for any reason the deflections reach some other scale point than that corresponding to reference deflection, the volume is given by the algebraic addition of the attenuator setting and the actual deflection as read on the meter scale. When volume delivered to an impedance differing from R ohms is to be measured, the volume indicated must be corrected to correspond to this difference. Correction (to be added algebraically) in vu $= 10 \log_{10} R/|Z|$, where $|Z|$ = magnitude of actual impedance.

4. Good Engineering Practice

The following items are not fundamental to this standard but are matters of good engineering practice.

4.1 Impedance. The preferred circuit impedances (R ohms) for calibration of the instrument are 600 and 150 ohms. The volume indicator is normally used as a bridging instrument and when so used its impedance must be sufficiently high so as not to influence unduly the waves in the circuit with which it is used. It is good practice to make the value of impedance not less than 12.5 R ohms, for use on an R-ohm circuit. This corresponds to an insertion loss of approximately 0.4 decibel.

4.2 Harmonic Distortion. The root-mean-square value of the harmonic distortion produced when the volume indicator is bridged across a resistive circuit impedance through which a sinusoidal wave between 25 and 8,000 cycles per second is being transmitted should not exceed 0.2 per cent of the fundamental.

4.3 Ability to Withstand Overload. The instrument should withstand without injury or effect on calibration, a momentary overload of ten times the voltage corresponding to reference deflection, and a continuous overload of five times that voltage, because of the great variation in amplitude which this indicator may encounter.

4.4 Scale. The point of reference deflection should be located within a sector between $\frac{2}{3}$ and $\frac{3}{4}$ of full scale. In addition to the vu scale, a 0-to-100 scale proportioned to voltage should be provided. Samples of the two types of scales in general use are shown in Figs. 1 and 2. The color of the scale card, expressed according to the Munsell system of color identification, should be 2.93Y(9.18/4.61). No markings other than those associated with the scale should be visible on the scale plate when the indicating meter is viewed directly from the front.

4.5 Marking. The nominal input impedance of the device and the impedance for which it is calibrated should is marked on the panel of the volume indicator assembly. This designation may be omitted when the volume indicator is a bridging device calibrated for 600 ohms. The marking of the meter scale does not uniquely determine the calibration. For example, the meter sensitivity might be such that 0 vu on the meter scale corresponds to a volume of +4 vu. In this case, the zero-loss position of the attenuator would be marked +4 vu.

A Gas-Discharge Noise Source for Eight-Millimeter Waves*

T. J. BRIDGES†

Summary—A noise source of the gas-discharge type for use at eight millimeters wavelength is described. Noise temperatures of 19,300 degrees K (18.2 db above room temperature) using neon, and 11,500 degrees K (15.9 db) using argon, have been measured by means of a microwave radiometer. These figures agree well with values calculated from gas-discharge theory.

Introduction

NOISE SOURCES of the gas discharge type have been described for use at cm wavelengths.[1,2] It is now established[2–4] that the noise radiated by such sources is associated with the high-electron temperature in the positive column of the discharge. This is usually of the order of 10^3–10^4 degrees K.

In order to scale such a discharge to shorter wavelengths, it must still be possible to match it to a waveguide, and the electron temperature must remain the same (or higher) at the shorter wavelength. By reducing the dimensions of the gas tube and increasing the pressure, both linearly with wavelength, it can be shown[5] that the electron temperature remains constant. Furthermore, if the scaling has been done in the above way, the tube current required to give a match is unchanged.

Tube and Mount for 8 mm

Using these principles, a tube and waveguide mount have been designed for the 8 mm wavelength band. These are illustrated in Fig. 1.

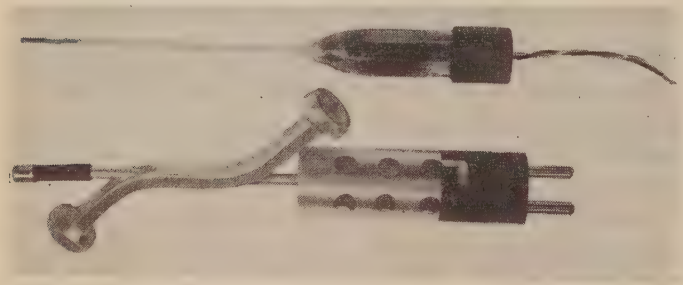

Fig. 1—8 mm noise tube and mount.

The part of the discharge tube that lies in the waveguide is 0.1 inch diameter, which opens out at one end into a bulb, containing a hot cathode. Mounting is in the *E*-plane, at an angle to the waveguide. To give room for the cathode end of the tube, and the anode connection, the waveguide is made with a double bend.

The chief difficulty experienced in the design lay in keeping the glass losses of the tube small. By using a thin-walled tube (<0.01 inch) however, it was possible to keep the loss down to about 0.5 db. Both neon and argon fillings have been used. With neon at a pressure of 30 mm Hg, and argon at 20 mm Hg, good matches were obtained with current above about 40 mA. Some of the measured characteristics of the tubes are given in the following table.

TABLE I

Tube No.	1	3
Gas	Ne	A
Gas pressure mm Hg	30	20
Insertion loss (discharge off) db	0.7	0.3
VSWR (discharge off: match behind)	1.1	1.13
VSWR (discharge on: short behind)	1.15	1.1
Tube current mA	45	35
Measured noise temperature, degrees K	19,300	11,500
Noise figure db	18.2	15.9
Calculated noise temperature, degrees K	21,000	10,000

Measurement of Noise Temperature

The amount of noise radiated from these tubes was measured on a microwave radiometer.[6] Comparison was made with a hot load at 500 degrees C., giving a direct measurement of noise temperature. The tubes containing neon gave a noise temperature of 19,300 degrees K, and those with argon 11,500 degrees K. Calculated values[5] from gas discharge theory are 21,000 degrees K and 10,000 degrees K, respectively.

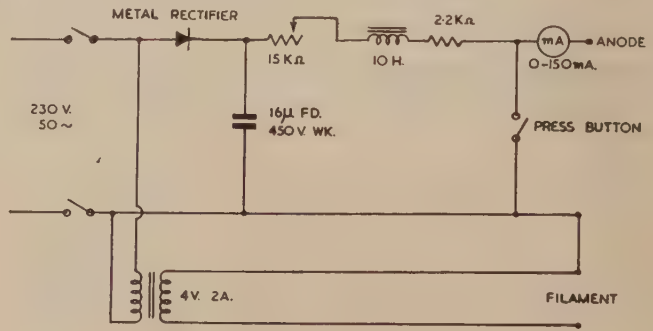

Fig. 2—Power unit for noise tube.

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Power Supply

A suitable power supply for the tube is shown in Fig. 2. The running voltage of the tubes is about 120, but 600–700 volts are required to strike the discharge. In order to supply this momentary over-voltage, a choke, carrying current, is suddenly open-circuited by means of a push-button switch.

Acknowledgment

The author is indebted to G. R. Nicoll, Radar Research Establishment, Malvern, England, for use of the 8 mm radiometer designed by him, and for assistance in making the noise-temperature measurements. He also wishes to thank P. O. Hawkins for helpful advice, and the Admiralty, for permission to publish this paper.

Transient Response in FM*

IGOR GUMOWSKI†

Summary—**The FM response to the unit impulse and the unit step function is calculated for a network whose transfer function is known. A meaning is assigned to the associated Fourier integral, which diverge in the Riemann sense. The method is generalized to any input function which vanishes for negative t. As an illustration of the method the impulse and the step FM responses of a single tuned circuit are calculated.**

The General Method

THE CALCULATION of transient response in FM is no more difficult than the calculation of transient response in AM. In fact, the method to be used is essentially the same.

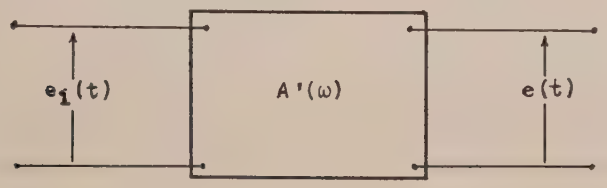

Fig. 1

If $A'(\omega)$ is the transfer function (complex system function) of the network whose response to the signal $e_i(t)$ is to be calculated (Fig. 1), the solution $e(t)$ can be obtained in all cases from the following relation:

$$e(t) = \frac{1}{2\pi}\int_{-\infty}^{\infty} e^{j\omega t}A'(\omega)\int_{-\infty}^{\infty} e^{-j\omega x}e_i(x)\,dx\,d\omega. \qquad (1)$$

From the physical point of view the above solution represents the well known principle of harmonic balance, and it is thus equally valid for FM as for AM. From the mathematical point of view, (1) represents simply the inverse Fourier transform of the product of the network transfer function by the transform of the input signal. If $f(t)$ is the signal to be transmitted, then the input voltage $e_i(t)$ becomes

$$e_i(t) = f(t)e^{j\omega_0 t} \qquad (2a)$$

for AM (suppressed carrier modulation), and

$$e_i(t) = e^{jF(t)}e^{j\omega_0 t} \qquad (2b)$$

for FM, where

$$F(t) = \int_{-\infty}^{t} f(x)\,dx \qquad (2c)$$

and ω_0 is the angular frequency of the carrier.

However, if this well known solution is used in a practical case, there appears a serious difficulty. When such popular functions as the unit step $Y(t)$ or the unit impulse $\delta(t)$ are chosen for the signal to be transmitted, it is found that the Fourier transform of $e_i(t)$ does not exist in the Riemann sense. This difficulty can be eliminated by using mathematically more complex signals to insure the existence of the integrals in the Riemann sense, but unfortunately the mathematical complication thus added renders the resulting solution quite unwieldy.

For theoretical as well as practical purposes, it is more advantageous to keep the divergent integrals and consider these integrals in the sense of the distribution theory (1), instead of considering them in the Riemann sense. In such a case it is easily verified that the Fourier transforms of $e_i(t)$ for $f(t)=a\delta(t)$, $-\pi<a<+\pi$, and $f(t)=aY(t)$, $-\infty<a<+\infty$, become respectively

$$\int_{-\infty}^{\infty} e^{-j\omega t}\cdot e^{jaY(t)}\cdot e^{j\omega_0 t}dt$$

and

$$= \pi(e^{ja}+1)\delta(\omega-\omega_0) + j\frac{e^{ja}-1}{\omega-\omega_0} \qquad (3)$$

$$\int_{-\infty}^{\infty} e^{-j\omega t}\cdot e^{jatY(t)}\cdot e^{j\omega_0 t}dt$$

$$= \pi\delta(\omega-\omega_0) + \pi\delta(\omega-\omega_0-a) + j\frac{a}{(\omega-\omega_0)(\omega-\omega_0-a)}. \qquad (4)$$

If the transfer function $A'(\omega)$ of the network is known, the transient response $e(t)$ can be calculated in a straightforward manner. In fact, letting $A'(\omega)=A(\omega-\omega_0)$ the FM impulse response of the network becomes

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$$e(t) = \frac{1}{2\pi}\int_{-\infty}^{\infty} e^{j\omega t}A'(\omega)\pi(e^{ja}+1)\delta(\omega-\omega_0)d\omega + \frac{j}{2\pi}\int_{-\infty}^{\infty} e^{j\omega t}A'(\omega)\frac{e^{ja}-1}{\omega-\omega_0}d\omega$$

$$= \frac{1}{2}(e^{ja}+1)e^{j\omega_0 t}\int_{-\infty}^{\infty} e^{jxt}A(x)dx - \frac{e^{ja}-1}{2\pi}e^{j\omega_0 t}\int_{-\infty}^{\infty} e^{jxt}A(x)\frac{dx}{jx}$$

$$= e^{j\omega_0 t}\left[A(0)e^{ja} - (e^{ja}-1)\int_{-\infty}^{t} u(x)dx\right], \quad (5a)$$

where

$$u(t) = \frac{1}{2\pi}\int_{-\infty}^{\infty} e^{jxt}A(x)dx \quad (5b)$$

is the impulse response (Green's function, weighting function) of the associated low-pass network. It is important to note that $u(t)$ does not have to be purely real, i.e., the associated low-pass network does not have to satisfy any conditions of "physical realizability."

If the FM impulse response is applied to a perfect discriminator, the detected response will be given by the derivative of the phase angle of the expression inside the brackets. In particular, if $u(t)$ is purely real, i.e. if the network of Fig. 1 was obtained from a low-pass network by means of a frequency translation, then the detected response is given by

$$e_d(t) = \frac{d}{dt}\operatorname{arc\,tg}\frac{A(0)\sin a - \sin a\int_{-\infty}^{t} u(x)dx}{A(0)\cos a + (1-\cos a)\int_{-\infty}^{t} u(x)dx}. \quad (6)$$

The FM step-function response is found in a similar manner:

$$E(t) = \frac{1}{2\pi}\int_{-\infty}^{\infty} e^{j\omega t}A'(\omega)[\pi\delta(\omega-\omega_0)+\pi\delta(\omega-\omega_0-a)]d\omega + \frac{1}{2\pi}\int_{-\infty}^{\infty} e^{j\omega t}A'(\omega)\frac{ja}{(\omega-\omega_0)(\omega-\omega_0-a)}d\omega$$

$$= e^{j\omega_0 t}\left[\frac{1}{2}A(0)+\frac{1}{2}A(a)e^{jat} + \frac{ja}{2\pi}\int_{-\infty}^{\infty} e^{jxt}\frac{A(x)}{x(x-a)}dx\right]$$

$$= e^{j\omega_0 t}\left[A(0) + e^{jat}\int_{-\infty}^{t} e^{-jax}u(x)dx - \int_{-\infty}^{t} u(x)dx\right], \quad (7a)$$

where again

$$u(t) = \frac{1}{2\pi}\int_{-\infty}^{\infty} e^{jxt}A(x)dx. \quad (7b)$$

If $u(t)$ is purely real, the detected step-function response will be given by

$$E_d(t) = \frac{d}{dt}\operatorname{arc\,tg}\frac{\int_{-\infty}^{t} u(x)\sin a(t-x)dx}{A(0)+\int_{-\infty}^{t} u(x)\cos a(t-x)dx - \int_{-\infty}^{t}\mu(x)dx}. \quad (8)$$

If the input signal $e_i(t)$ is of a more general form,

$$e_i(t) = e^{jF(t)\cdot Y(t)}\cdot e^{j\omega_0 t}, \quad (9)$$

where

$$F(t) = \int_{-\infty}^{t} f(x)dx$$

is an arbitrary function, the FM transient response becomes simply

$$e(t) = e^{j\omega_0 t}\left[A(0) + \int_{-\infty}^{t} e^{jF(t-x)}u(x)dx - \int_{-\infty}^{t} u(x)dx\right]. \quad (10)$$

Examining the calculated FM transient responses we notice that they are determined completely by the impulse response $u(t)$ of the associated low-pass network. Since $u(t)$ depends both on the amplitude and the phase characteristic (obtained by steady state sine wave analysis) of the network of Fig. 1, the complete and the detected FM transient responses do also. It is thus necessary to consider both the bandwidth and the phase-shift when designing a network for a given FM transient response.

FM Transient Response of a Single-tuned Circuit

As an illustration of the general method described previously let us determine the FM transient response of a single-tuned circuit (Fig. 2) when the input signal

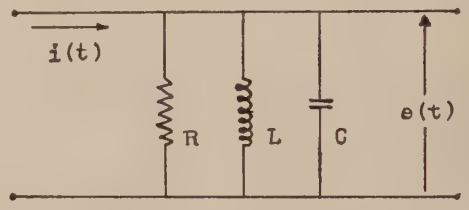

Fig. 2

is fed from a constant current source. By elementary analysis it is found that the transfer function of this network is given by

$$A'(\omega) = R\frac{d}{d + j\left(\frac{\omega}{\omega_0}-\frac{\omega_0}{\omega}\right)}$$

where

$$\omega_0^2 = \frac{1}{LC} \quad \text{and} \quad d = \frac{1}{\omega_0 RC}. \tag{11a}$$

If the bandwidth is small compared with the resonant frequency $\omega_0/2\pi$, then this transfer function simplifies to

$$A'(\omega) = R \frac{\frac{1}{2}d}{\frac{1}{2}d + j\left(\frac{\omega - \omega_0}{\omega_0}\right)}. \tag{11b}$$

Hence, the transfer function $A(x)$ of the associated low-pass network is

$$A(x) = R \frac{\frac{1}{2}d}{\frac{1}{2}d + j\frac{x}{\omega_0}}, \tag{12}$$

and its impulse response is

$$u(t) = \frac{1}{2\pi}\int_{-\infty}^{\infty} e^{jxt}A(x)dx = \frac{Rd}{4\pi}\int_{-\infty}^{\infty} \frac{e^{jxt}dx}{\frac{1}{2}d + j\frac{x}{\omega_0}}$$

$$= \tfrac{1}{2}Rd\omega_0 Y(t)e^{-(d\omega_0/2)t}. \tag{13}$$

Therefore, if $f(t) = a\delta(t)$, i.e. $i(t) = e^{jaY(t)} \cdot e^{j\omega_0 t}$, the FM impulse response will be given by

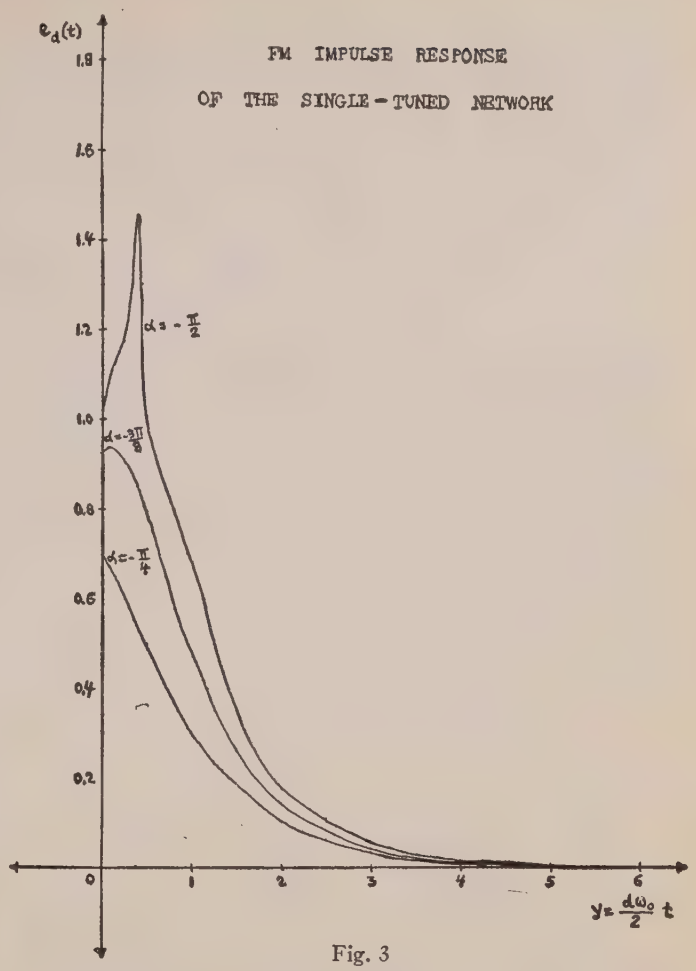

Fig. 3

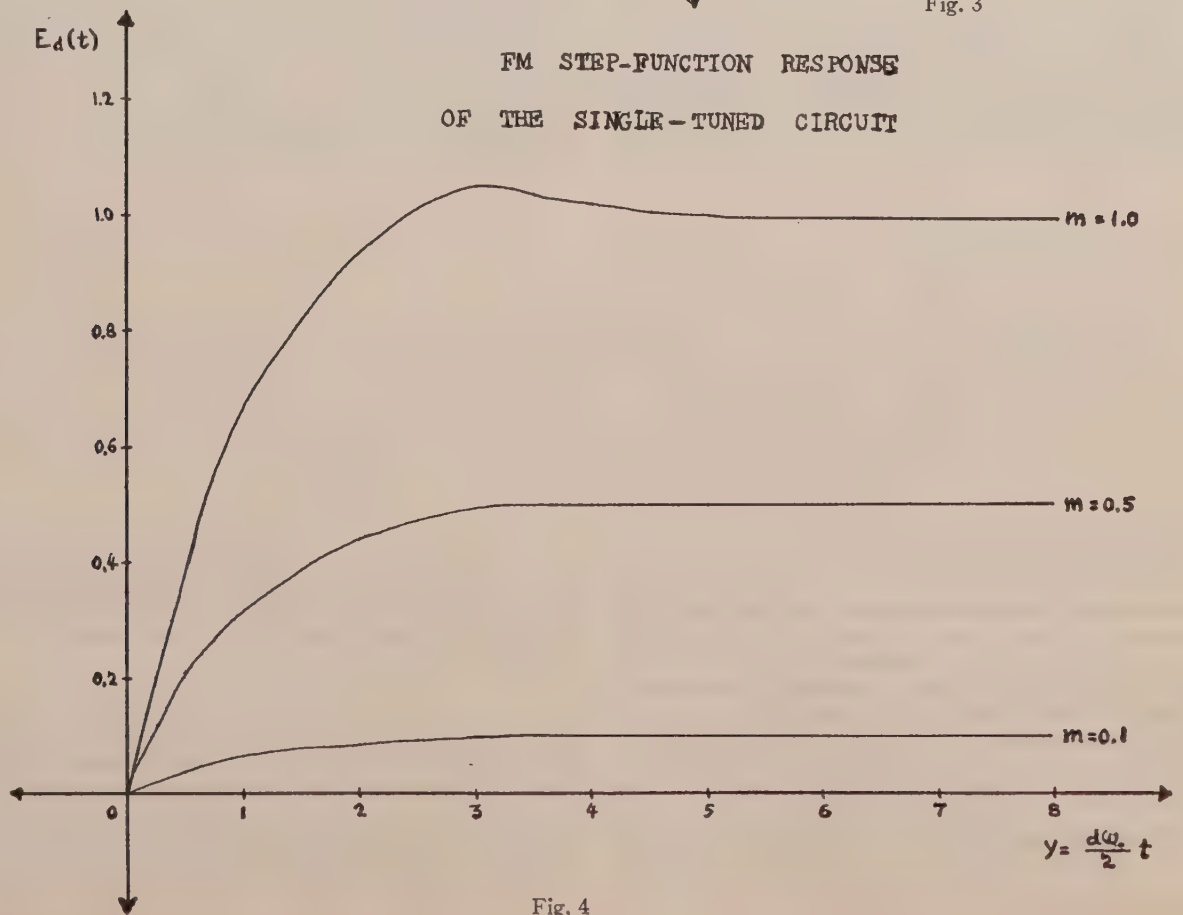

Fig. 4

$$e(t) = Re^{j\omega_0 t}\left[e^{ja} - (e^{ja}-1)\frac{d\omega_0}{2}\int_{-\infty}^{t} Y(x)\epsilon^{-(d\omega_0/2)x}dx\right]$$

$$e(t) = \begin{cases} Re^{j\omega_0 t}\cdot e^{ja} & \text{for } t<0 \\ Re^{j\omega_0 t}\left[1+e^{ja}\cdot e^{-(d\omega_0/2)t} - e^{-(d\omega_0/2)t}\right] & \text{for } t>0, \end{cases} \quad (14)$$

and the detected response becomes

$$e_d(t) = \frac{d}{dt}\operatorname{arc\,tg}\frac{\sin a e^{-(d\omega_0/2)t}}{1+(\cos a - 1)e^{-(d\omega_0/2)t}}\cdot \quad (15)$$

To examine how the detected impulse response varies with the phase sweep a, it is only necessary to plot $e_d(t)$ vs $\frac{1}{2}d\omega_0 t$ for various values of $-\pi<a<+\pi$ (Fig. 3).

However, usually it is more convenient to determine how the transient response is related to the frequency sweep instead of to the phase sweep. This can be easily accomplished by considering the FM response to the step function. In fact,

$$E(t) = e^{j\omega_0 t}\left[R + e^{jat}\int_{-\infty}^{t} e^{-jax}\cdot\frac{Rd\omega_0}{2}\cdot e^{-(d\omega_0/2)x}Y(x)dx - \int_{-\infty}^{t}\frac{Rd\omega_0}{2}e^{-(d\omega_0/2)x}\cdot Y(x)dx\right]$$

$$= \begin{cases} Re^{j\omega_0 t} & \text{for } t<0 \\ Re^{j\omega_0 t}\left[\dfrac{1}{1+j\dfrac{2a}{d\omega_0}}\left(e^{jat}+j\dfrac{2a}{d\omega_0}e^{-(d\omega_0/2)t}\right)\right] & \text{for } t>0. \end{cases} \quad (16)$$

Since $d\omega_0/2\pi$ is the bandwidth of the single-tuned circuit and $2a/d\omega_0$ can be considered as the "sweep index," the response $E(t)$ can be easily written in a per-unit form.

Letting $y=\frac{1}{2}d\omega_0 t$ and $m=2a/d\omega_0$ we obtain for $t>0$

$$E(t) = Re^{j\omega_0 t}\left[\frac{1}{1+jm}(e^{jmy}+jme^{-y})\right] \quad (17)$$

and

$$E_d(t) = \frac{d}{dt}\operatorname{arc\,tg}\frac{me^{-y}+\sin my - m\cos my}{me^{-y}+\cos my + m\sin my}\cdot \quad (18)$$

Fig. 4 shows the step-function responses for $m=0.1$, $m=0.5$ and $m=1$.

Examining the plotted responses we notice that the quality of the FM transmission decreases as the frequency or the phase sweep increases. Especially, when the frequency sweep is equal or greater than the bandwidth of the single-tuned circuit, or when the phase sweep is close to 180 degrees (a close to $\pm\pi$), the FM transient response is quite poor.

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Unit Construction Practice (UCP) Applied to Line-Transmission Equipment*

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Summary—**The equipment practice described and illustrated in this paper and referred to as Unit Construction Practice is based on the conception that the number of different circuit functions used in line-transmission technique is small and that these circuit functions can be performed by relatively few standardized components. Once these facts are recognized it is an obvious step to adopt an equipment practice in which a series of self-contained units, each comprising an assembly of standardized components and designed to perform a distinct circuit function, become the "bricks" which can be used in different combinations to build different types of systems.**

It is claimed that the new practice results in equipment which offers a degree of standardization combined with flexibility and a ready adoption to civil or military purposes not possible using earlier practices. Considerable space-saving is also effected.

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Introduction

THE RAPID EXPANSION in long-distance telegraphy and telephony, which took place during the five or six years prior to the second World War, was brought about by employing carrier-transmission methods to multiply the message-handling capacity of existing line facilities. The net result was that the

volume and complexity of inside plant materially increased and, already, at the outbreak of war the provision of building accommodation and skilled maintenance staff was becoming critical; at the end of the war it was even more so.

However, the aftermath of war brought certain compensations. For the first time large-scale application of carrier-transmission methods had been applied to the provision of military circuits. In the early part of the war extensive use was made of normal pre-war equipment. Experience showed that this equipment was too heavy and bulky for military operations, and too complex for relatively unskilled personnel to service satisfactorily. Efforts were directed towards the production of smaller and lighter components; new circuit techniques were developed, and, around these, new smaller and lighter equipment was designed and constructed. The problems of maintaining and operating equipment with semi-skilled personnel received special study. Maintenance and operating routines were simplified. In fact a new era in transmission technique began. However it was inevitable that the design of this wartime equipment was subservient to the overriding needs of the armed forces and only scant attention was paid to probable peace-time requirements.

Towards the end of 1948 the repercussions on line-transmission technique brought about by the war were becoming apparent. The manufacturers and armed forces had become acutely aware of the desirability of integrating the design of post-war equipment so that, if the need arose again, it could quickly be adapted to military requirements. All concerned had become conscious of the need for greater standardization to facilitate production, supply, and maintenance of equipment. All concerned had become conscious also of the need for smaller equipment so as to permit more efficient utilization of the limited building accommodation now available in the post-war world.

Early in 1949, the Company with which the writer is associated initiated a development project to provide a more rational approach towards the design of line-transmission equipment, and the following pages of this article outline the design principles underlying the Unit Construction Practice (UCP) which emerged as a result.

Design Considerations

The principal consideration was to reconcile the conflicting claims of reliable and efficient performance, ease of maintenance, economy of mounting space, suitability alike for large and small administrations and ready adaptation for military purposes.

To ensure reliable and efficient performance it was decided, wherever possible, to avoid the use of abnormally miniatured components, and to reduce the bulk of the equipment by efficient utilization of the available mounting space using normal approved and accepted types of components.

Experience during the war years had shown that although the mounting of groups of components in hermatically sealed containers provided excellent protection against the effects of humidity, maintenance, and stocking of spares became more costly. In view of this it was decided, whenever possible, to protect components individually against the ingress of moisture.

In earlier equipment practice the apparatus panels were mounted on both sides of a 3-inch channel iron bay framework with the units projecting outward, thus wasting the space between front and rear panels. It was obvious that an economy of mounting space could be effected by a re-design of the bay framework. After consideration it was decided to mount the new equipment on two bay-side frameworks which when placed back-to-back would occupy the same floor space as an existing fully equipped double-sided bay.

Having decided what general category of components would be used and having, by consideration of the new bay-side frameworks, fixed the space available for mounting them, the next step was to study line-transmission circuit technique with the object of providing a greater degree of standardization.

Line-transmission systems comprise principally an assembly of entities, each of which performs a different but distinct circuit function, for example, amplification, modulation, filtration and so forth, all of which contribute to the performance of the system as a whole. A study of earlier practice showed that the assemblies of components required to perform these functions were rarely designed as self-contained units. Furthermore, assemblies of components used in one particular type of equipment were usually different from assemblies used in another type of equipment even though these assemblies were designed to perform essentially similar functions.

Further study revealed to what remarkable degree the same, or basically the same entities could be used without change to build different types of transmission systems. At this stage it was a logical step to conceive the idea of a series of standardized self-contained units, of a size and shape fully to utilize the available bay-side framework mounting space, which would become the "bricks" that could be used in different combinations to build different systems. In general each unit would be designed to perform only one circuit function and care would be taken to ensure that its electrical performance was suitable for as many different applications as possible. For obvious reasons these "bricks" became styled Functional Units.

Attention was next directed towards the best method of mounting these "bricks" in various combinations to form appropriate panel assemblies. A universal type of panel frame held special appeal and compared favorably with earlier equipment practice where it was necessary to drill and tap each mounting plate in a manner peculiar to the particular panel assembly required.

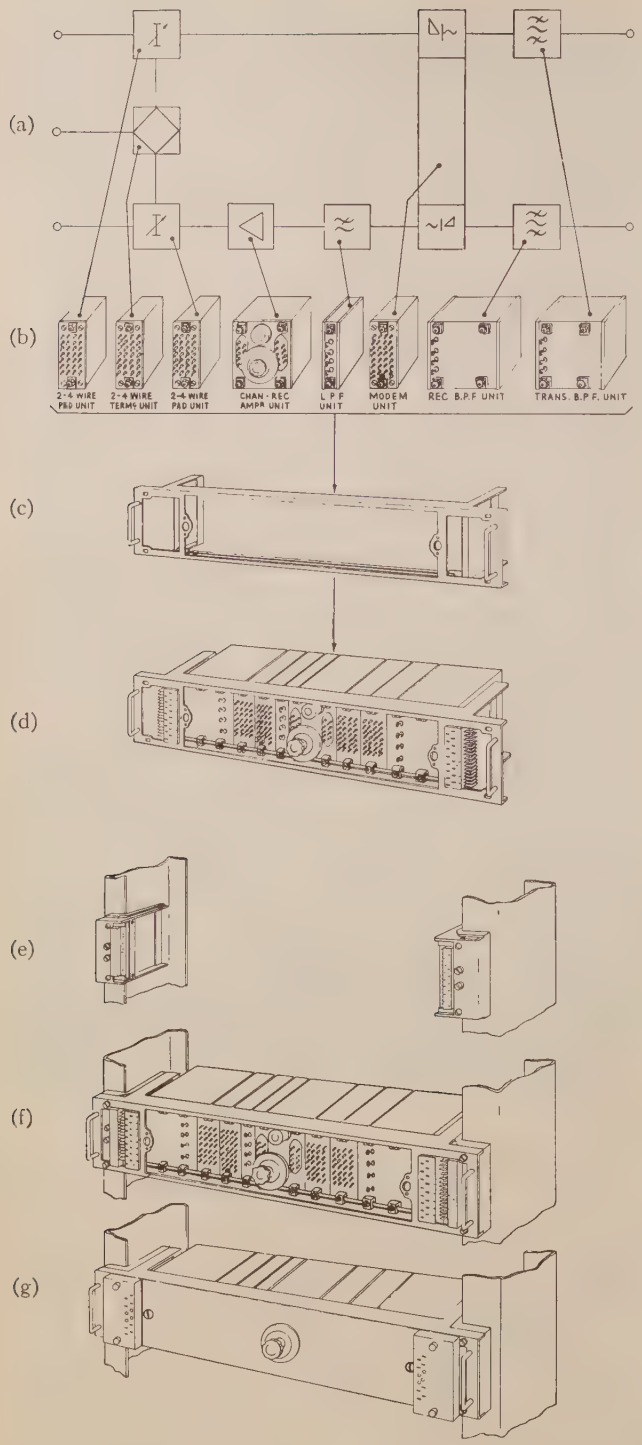


Fig. 1(a) to (g)—Schematic showing the principles of UCP applied to the design of a typical "channel panel" for a carrier-telephone system.

Thought next turned to the problem of making the bay-to-panel connections. It seemed that the advantages of quickly detachable jack-in panels could not be ignored. However, the prejudices against the use of dry contacts in circuits carrying small currents was still considerable. Eventually it was decided to arrange for the panel to jack into the bay-side but also to make provision so that, if desired, the jack contacts could be bridged by short soldered straps.

Principles of UCP

The principles of the new practice are illustrated by Figs. 1(a) to 1(g) which gives an example of a typical "channel panel" for use in a carrier-telephone equipment.

Fig. 1(a) shows a block schematic of the panel. Each distinct circuit function, indicated by separate blocks on the schematic, is performed by a Functional Unit. Fig. 1(b) illustrates the Functional Units required.

Fig. 1(c) illustrates the universal panel frame, and Fig. 1(d) shows the Functional Units mounted on it.

Fig. 1(e) shows a section through the side members of the bay-side framework, and Fig. 1(f) the bay-side framework with the panel mounted on it. The connections between bay and panel wiring are made by plug and socket assemblies which also provide test points for routine measurements.

Fig. 1(g) shows the panel fitted with its dust cover. This is a flat plate attached at each end by quick release fasteners.

From this example it will be clear that, with the exception of the band-pass filter units, the same Functional Units can be used over and over again to build "channel panels" for many different carrier systems. Furthermore, if each unit is designed to have a high return loss at its input and output, then the over-all frequency response characteristic of any combination of units will be substantially the algebraic sum of the characteristics of the individual units. This is of considerable advantage when designing systems to meet specific performance requirements since it is possible for the designer to predetermine with reasonable precision what over-all performance can be expected from a given combination of units.

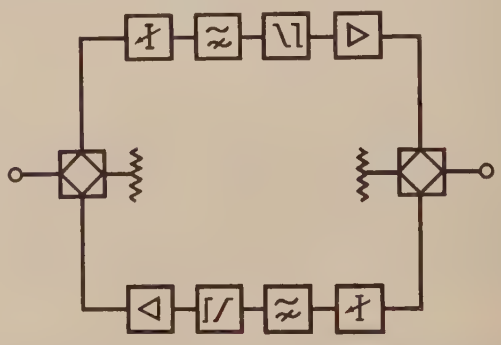


Fig. 2—Block schematic of a VF 2-wire repeater panel.

Consider next the typical 2-wire voice-frequency repeater panel shown in block schematic form in Fig. 2. It can readily be shown that the 2/4-wire terminating units, pad units and voice-frequency amplifier units can be identical to those used to build up the channel panel referred to in the previous example. A study of line-transmission technique will reveal many other examples where some or all of these typical Functional Units can be used in different combinations to build different types of systems.

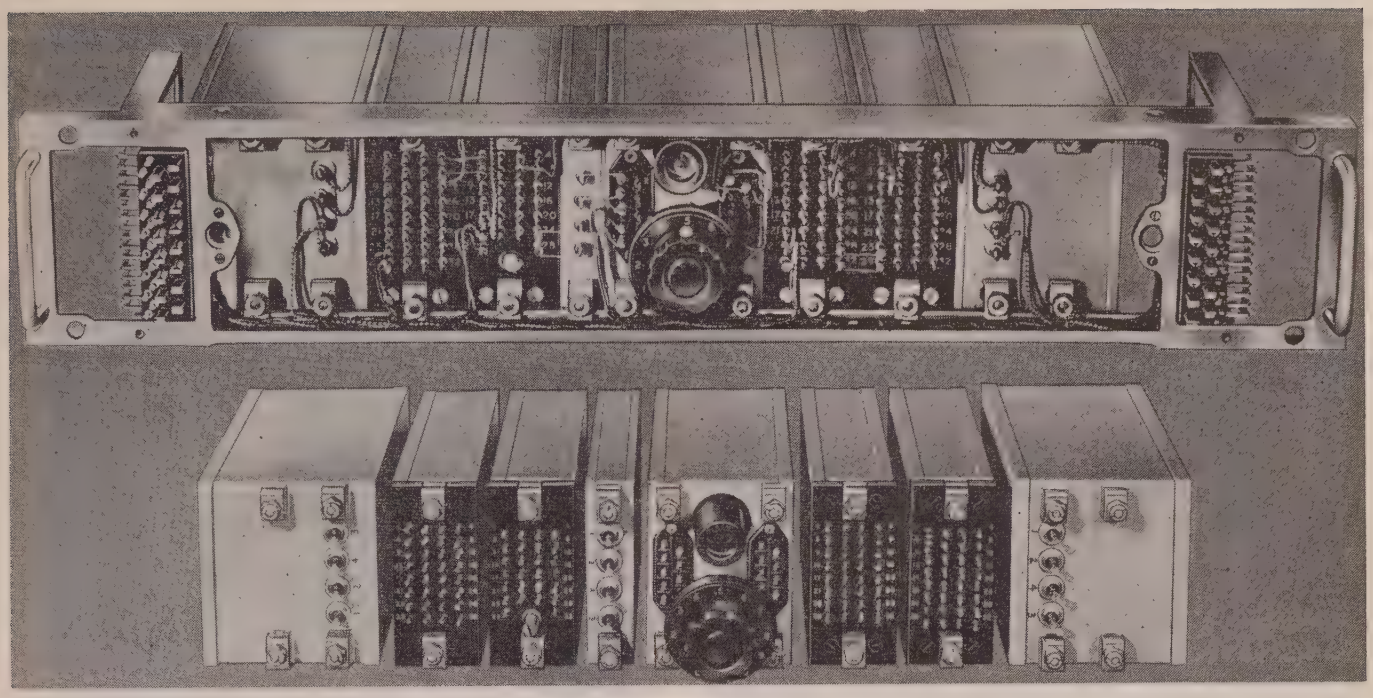


Fig. 3—Typical UCP carrier-telephone "channel panel."

It will be seen that the principles of UCP are based on the conception that the number of different circuit functions employed in line transmission technique is small and that these circuit functions can be performed by relatively few standardized components. Once these facts are recognized it is an obvious and logical step to adopt an equipment practice in which a minimum number of self-contained units, each comprising an assembly of standardized components and designed to perform a specific circuit function, are used in different combinations to build different types of equipment.

Design Features

It has been shown that the fundamental entities of UCP are Functional Units and that these units are assembled in various combinations on universal panel frames which are jacked into the bay-side. The following pages describe certain design features of the Functional Units, universal panel frames and bay-side frameworks used in the new equipment practice.

A. Functional Units

The majority of the Functional Units comprise a simple metal framework in which the various components are mounted, a top plate mounting terminal tags and any necessary controls, and a metal cover which encloses the whole assembly. Each unit is so designed that the whole of the mounting space within the framework is fully utilized. In the larger units components and wiring are made more accessible by hinging the frameworks.

The mechanical design of filter and equalizer network units remains substantially unchanged from earlier equipment practice. However, a considerable size reduction has been achieved by the use of silvered mica and/or polystyrene capacitors, and smaller powder core inductors having improved intermodulation characteristics.

Fig. 3 shows the Functional Units required to build up the typical channel panel shown diagrammatically in Fig. 1. All these units are designed to mount on to a $3\frac{1}{2}$-inch panel frame. It has been found that with few exceptions this one height of unit is sufficient for all purposes. Figs. 4 (and 5, following page), show exposed views of two typical units.

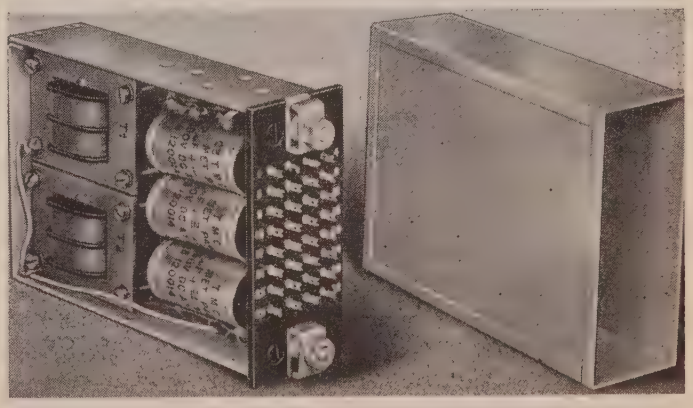
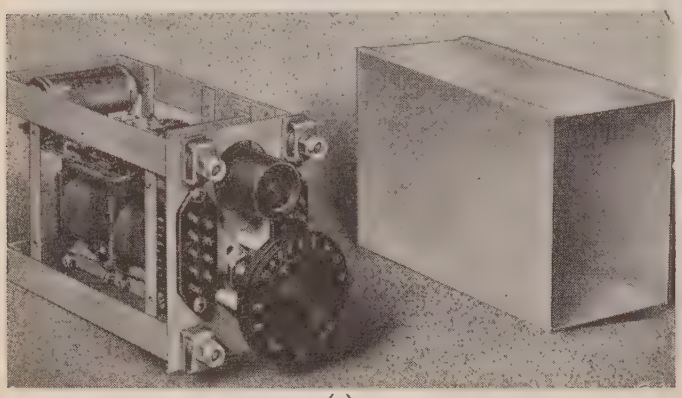


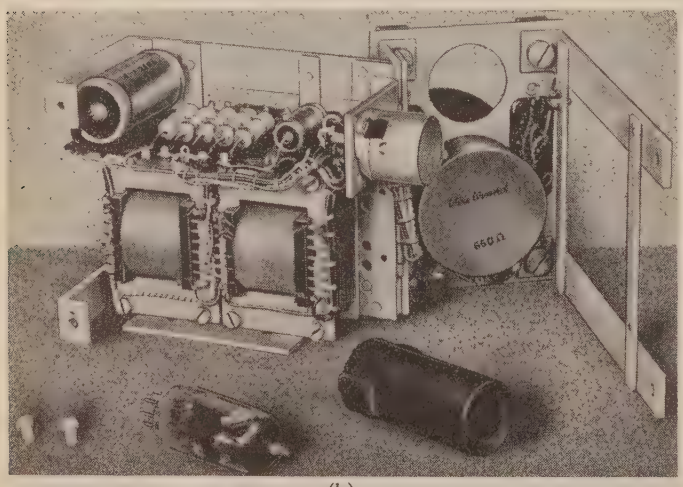
Fig. 4—2/4-wire Functional Unit.

It is interesting to consider the design of the single tube voice-frequency amplifier shown in Fig. 5. This amplifier is so designed that it may be used as a channel-receive amplifier in all types of carrier-telephone systems, and as a line amplifier in voice-frequency telegraph systems, 2-wire and 4-wire voice-frequency repeater systems, 2-band systems, privacy systems, etc.

Study of these applications showed the amplifier required to have maximum working gain of 44 db, about 10 db of negative feedback, substantially flat frequency response and high input and output return loss over frequency range of 300 cps to 3,400 cps, power handling capacity of +17 dbm with harmonic distortion not greater than 30 db, and to operate from HT supply of 150v. Using 6*AK*5 tube and modern circuit technique it was relatively easy to meet this specification.



(a)



(b)

Fig. 5(a) and (b)—VF amplifier Functional Unit showing how unit framework is hinged to give accessibility to wiring and components.

B. Panel Frames

The structural form of the panel frame is shown in the line drawings of Fig. 1 and the method of attaching Functional Units to it is shown in Fig. 6. To provide light weight combined with adequate strength to withstand vibration and shock effects the frame is constructed throughout of sheet steel.

When the Functional Units are mounted on the panel frame the wiring between units lies all in one plane and the terminal tags of the units are accessible for making voltage and level measurements. This has particular advantage from a point of view of servicing and maintenance. If fault develops on any panel, it is possible, by making level measurement at input and output of each unit in turn, quickly to locate the faulty unit which may then be disconnected and replaced without difficulty.

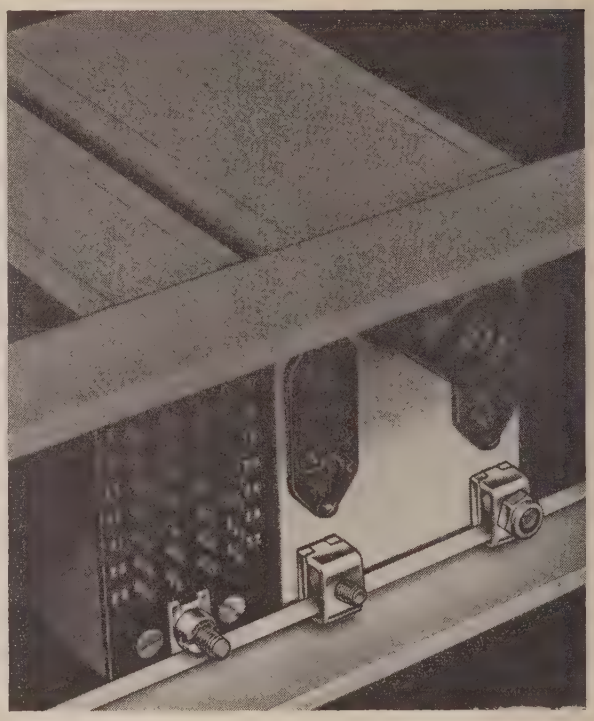


Fig. 6—Method of attaching Functional Units to panel frame.

The end members of the panel frame in combination with details fitted to the bay-side framework form what are known as "panel slides." These slides provide means for guiding the panel into position and taking the weight while it is being fitted to or removed from the bay-side. They also provide means for mounting the plug and socket assemblies for making connections between panel and bay wiring.

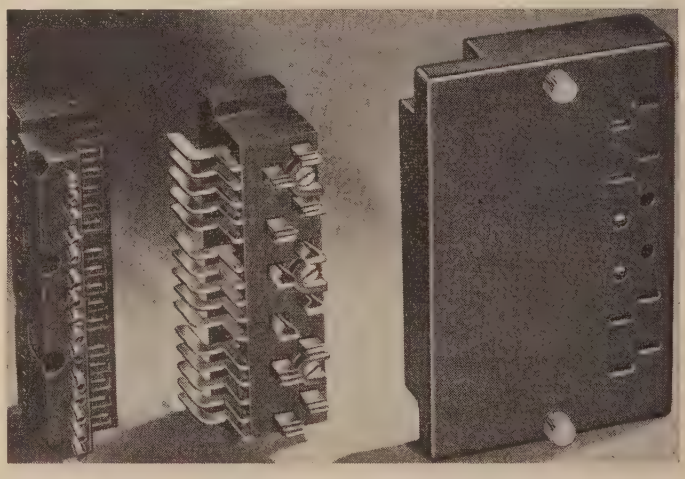


Fig. 7—Plug and socket assembly used in the new equipment practice.

C. Panel-to-Bay Connections

The design of the plug and socket assemblies is evident from Fig. 7. Accurate registration of blades and sockets is ensured by the panel slides. Provision is made for wiring across the contacts if desired, and sockets for routine measurements are provided on the panel plugs.

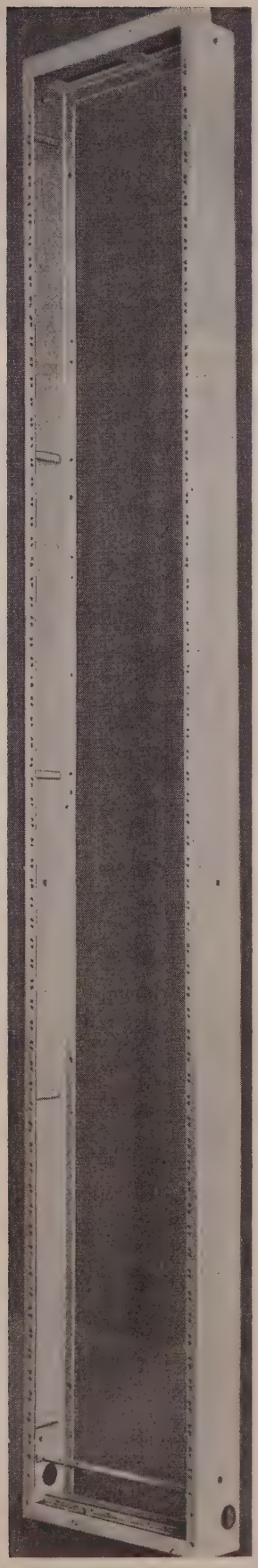


Fig. 8—Bay-side framework. The framework is constructed throughout of lightweight folded sheet stel.

D. Bay-side Framework

Fig. 8 shows new bay-side framework, is constructed throughout of folded sheet steel, and how a step is formed at top to carry office and inter-bay cabling.

E. Miscellaneous Panels

The number of bay-to-panel connections which can be provided by the plug and socket assemblies is of necessity restricted, and a maximum of thirty outlets is available on a $3\frac{1}{2}$-inch jack-in panel frame. This number is adequate for most purposes; occasionally, however, a greater number is required, and when this is necessary, the panel is hinged at one end and the bay cable wired direct to the panel units.

Sometimes instances arise where it is more convenient to depart from the Functional Unit principle, and either to combine more than one circuit function in the same unit, or to provide the required circuit function by a combination of units. A typical example of the latter is seen in Fig. 9, which shows the line amplifier for a 12-channel open-wire carrier-telephone equipment. The amplifier is built up from four separate units, namely, input attenuator, input stage, intermediate stage, and output stage. These units are mechanically similar to Functional Units used in other parts of the equipment.

Conclusions

An attempt has been made to describe and illustrate the design principles of UCP and to outline the considerations which led up to its development. The new practice is based on the conception that the number of different circuit functions employed in line-transmission technique is small, and that these functions can be performed by relatively few standardized components. Advantage is taken of these facts by designing a self-contained unit series, each comprising an assembly of standardized components and in general performs a distinct circuit function. These Functional Units are used in different combinations to build different types of systems.

It is considered that the new practice offers a more rational approach to the design of line-transmission equipment than is possible using earlier practice. From the manufacturing viewpoint, the production advantages of building equipment from a limited number of standardized units are obvious. So far as the user is concerned, the advantages are in terms of ease of maintenance and testing routines, and the stocking of spare parts. Both user and manufacturer alike benefit because quantities of panel frames, bay-side frameworks, and Functional Units can be held in stock, and different types of systems built up with a minimum of delay.

The first costs of UCP equipment compare favorably with earlier equipments designed to meet similar performance requirements. This is to be expected since, by restricting the different types of unit to a minimum, a greater control over the quality and uniformity of the product is possible, and standardized production and testing methods can be adopted.

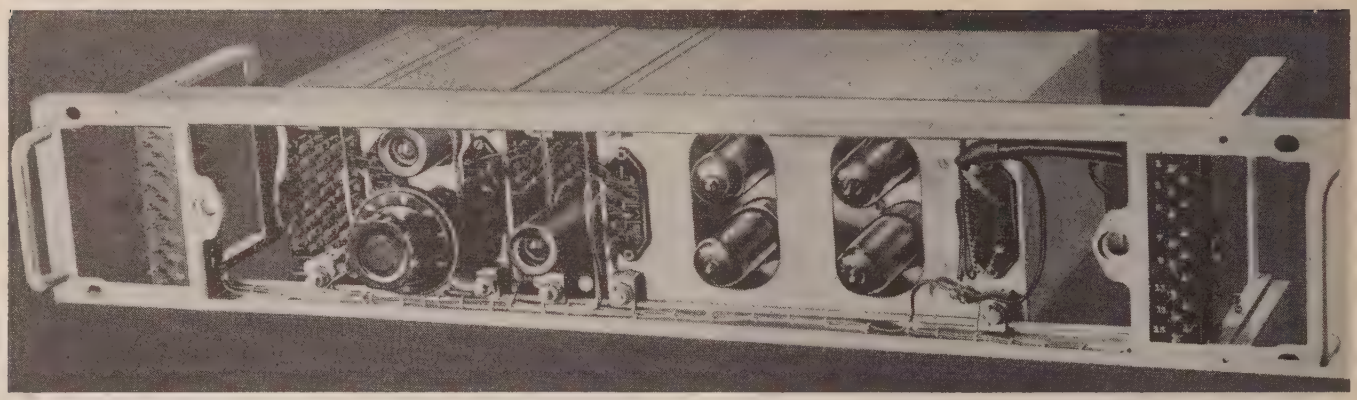


Fig. 9—Line amplifier for a 12-channel open-wire carrier-telephone system.

UCP equipment is not "miniaturized" in the sense that the term is often used. The reduction in bulk, evident from the group of 12-channel carrier-telephone terminal equipments shown in Fig. 10, in comparison with earlier equipment is as much as two- or three-fold. This has been achieved by efficient utilization of the available mounting space rather than by the use of abnormally miniaturized components. In spite of the considerable space-saving it will be seen from the illustrations that accessibility has been improved over that provided by earlier equipment practices.

Because the various Functional Units are so designed that they may be used in a variety of system applications, it is necessary for only a limited number of different types of units to be held in stock to serve as spares for a wide range of systems. Furthermore, since the panels can readily be detached from the bay-side frameworks and the units easily removed from the panel frames, in the event of a fault occurring, a spare unit or panel can be fitted with a minimum waste of time and effort.

In the new practice the heavy double-sided channel iron racks used in earlier equipment practice have been abandoned in favor of lightweight folded sheet steel bay-side frameworks. This feature, coupled with the fact that the panels may readily be detached from the bay-side frameworks, is particularly useful during installation or transportation since the bay-side frameworks and the individual panels can be handled separately.

Summarizing it will be seen that the principal features of UCP equipment are a bay-side framework which permits the maximum use of the available mounting space, a universal jack-in panel frame, and a series of Functional Units which can be mounted in any combination on the panel frame. Each of these features benefits one or other aspect of line-transmission practice. Space-saving reduces the bulk of equipment required to provide a given service, thereby easing the problems associated with the provision of building accommodation. The jack-in panel enables service to be restored quickly in the event of a fault developing in the equipment. The provision of a minimum number of different types of Functional Units to satisfy all probable requirements not only eases the problems concerned with maintenance routines and the stocking of spare

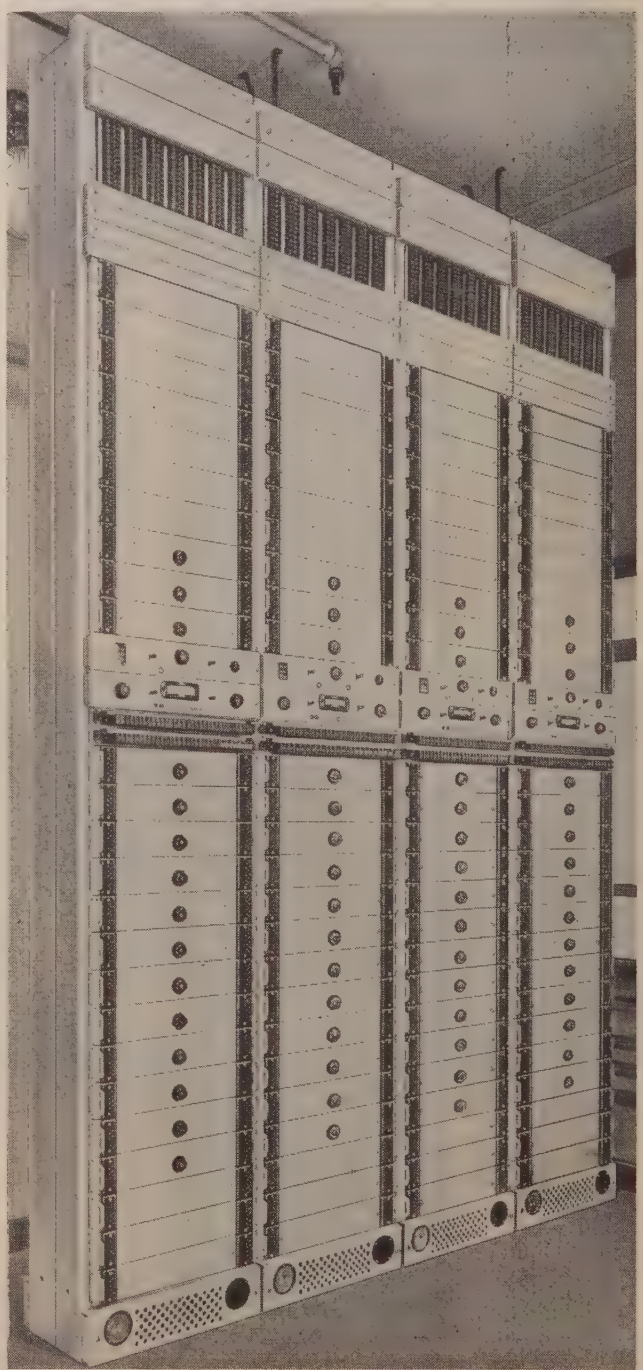


Fig. 10—Group of 12-channel carrier-telephone equipments for use on cable circuits. Each bay-side comprises a complete 12-channel terminal equipment.

parts, but also offers a degree of standardization combined with flexibility and a ready adaptation to civil or military purposes not possible using earlier equipment practice.

Although this article has dealt specifically with the application of UCP to line-transmission equipment, it is thought that the design principles of the new practice could be applied with advantage in other fields of electronic engineering. Essentially the design approach is to analyze the particular application envisaged, by means of functional block schematics, to make sure that the various circuit functions occur sufficiently frequently to justify the method of attack, and then to design a minimum number of self-contained functional units which can be used over and over again as the "bricks" to build different equipments.

Acknowledgments

The writer wishes to acknowledge his indebtedness to J. G. Flint, Technical Director, Telephone Manufacturing Co. Ltd., for permission to publish the information contained in the article, and to Messrs. Automatic Telephone and Electric Co. Ltd., for the use of certain photographs to illustrate the test.

Switching Time in Junction Diodes and Junction Transistors*

ROBERT H. KINGSTON†, ASSOCIATE MEMBER, IRE

Summary—**The time in which a junction diode may be switched from forward to reverse conduction is of great importance in computing networks. By considering the behavior of the minority carriers in a diode in a representative switching circuit an approximate solution for the switching transient may be derived. The transient is separated into two phases: first, one of constant current, where the flow is limited by the external resistance, and second, a "collection" phase, where the current decays at a rate determined by the minority carrier lifetime and the dimensions of the diode. A critical parameter in the solution is the ratio of the short-circuit reverse current to the forward current before switching. The mathematical treatment is a boundary value solution of the minority carrier diffusion equations which is accomplished by the use of Laplace transformations. The duration of the two phases of current flow is determined for a planar junction, a hemispherical junction, and for a planar junction with junction-to-contact distance small compared to a diffusion length. The last treatment is extended to the junction transistor and the behavior of the collector current is calculated. The general results indicate that for a given minority carrier lifetime the last two of the three diode structures will give the smallest switching times. In addition it is found generally that the time is minimized by decreasing lifetime and increasing the ratio of reverse to forward current.**

Introduction

IN THE USE of junction diodes in computer-switching circuits the switching time of the diode from forward to reverse voltage is of great importance. The problem may be represented as follows: In Fig. 1, a forward current, I_f, is flowing in the diode; at time, $t=0$, the switch, S, is thrown to the right. The switching time may now be defined as that time in which the voltage, V_D, reaches 90 per cent of the battery voltage. (The final static resistance of the diode is assumed to be much greater than R_0.)

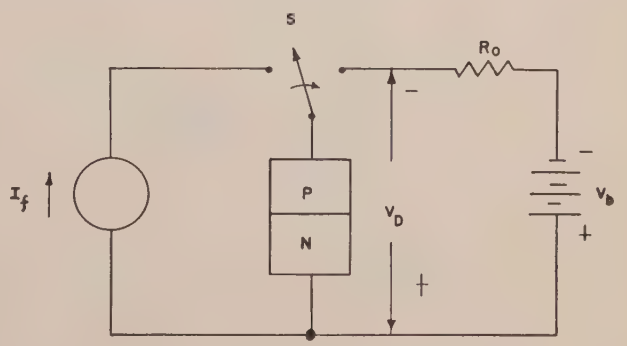


Fig. 1—Representative switching circuit.

Physical Model[1]

For simplicity, the treatment to follow is applied to a p-n junction where the conductivity of the p-type material is much greater than that of the n-type material. Practical embodiments of such a model are the familiar bonded n-type diodes and indium alloy-process devices. This limitation gives the qualitative picture of the carrier and potential distribution shown in Fig. 2. Here, the energy-band representation shows a forward current, I_f, proportional to the negative gradient of the hole concentration at the right of the barrier. The electron density in the p-type material is negligible since $\sigma_p \gg \sigma_n$ and is therefore neglected.

Given the above initial picture of the junction, it is now necessary to consider the behavior of the system after $t=0$, when the diode is switched into the reverse-biasing circuit. To anticipate the results it may be stated immediately that the initial transient in the circuit will be a constant current, given by V_b/R_0. That is, for a reasonable length of time after switching, the voltage

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† Lincoln Laboratory, MIT, Cambridge, Mass.

[1] The physical description follows the methods and notation of W. Shockley, "Electrons and Holes in Semiconductors," D. Van Nostrand Co., Inc., New York, N. Y.; 1950, or "The theory of p-n junctions in semiconductors and p-n junction transistors," *Bell Sys. Tech. Jour.*, vol. 28, p. 435; 1949.

across the junction will be small compared to the battery voltage; therefore, the current will be completely determined by the series resistance. (A practical circuit might have a battery voltage of 20 v and a series resistance of 20,000 ohms.) It should be emphasized that the voltage across the junction (the "space charge" voltage) does not change abruptly at $t=0$; actually, it is a dependent function of the hole density at the barrier which in turn is determined by the flow of diffusion

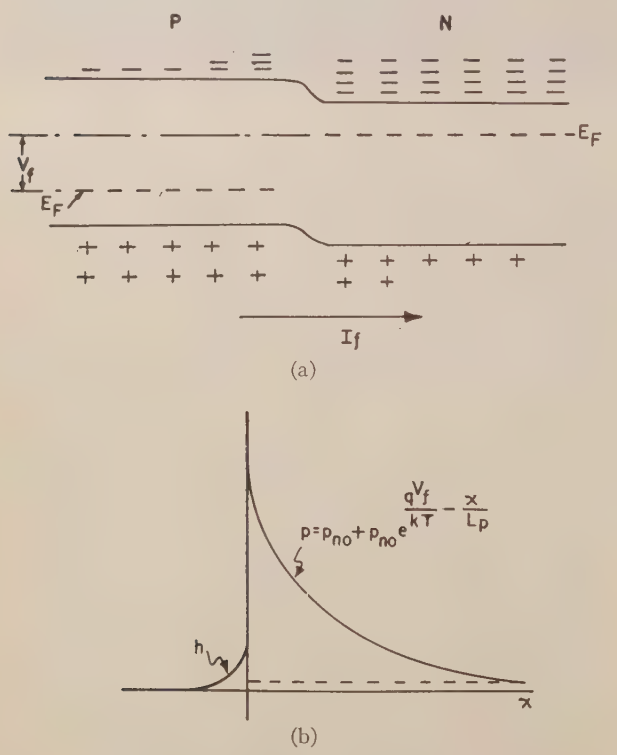

Fig. 2—(a) Energy bands; and (b) Minority-carrier densities for forward biased p-n junction.

current across the barrier. If the term p_{n0} in Fig. 2, b is neglected, the behavior of the hole density and space-charge potential may be represented as in Fig. 3.

Between $t=0$ and T_I, a constant current, $I_r=V_b/R_0$, flows to the left across the barrier; the junction voltage is given by kT/q ln p_0/p_{n0} plus a small ohmic contribution which is neglected. The hole density, p_0, at the boundary may be determined as a function of time from a solution of the diffusion equation. (The dashed lines represent the functions used in the solutions below.) At or near T_I, the hole density approaches zero, V tends toward minus infinity and a new set of conditions now holds. These are that the hole density at the barrier is practically zero, therefore V is now much larger than kT/q and the diffusion current is no longer constant. The behavior of the hole density is now calculated according to this condition, $(p(0)=0)$, and $J=-Ddp/dx$ gives the current as a function of time. It might be noted that the electron current at the barrier is neglected throughout the treatment. This assumption is found to be valid since any electron flow must come either from the space charge region or the p-type material. The former is unlikely since the space-charge current, $I_n=CdV/dt$, is found to be negligible compared to the diffusion current, and the latter is ruled out by the high conductivity of the p-type material. The remainder of the treatment is concerned with the mathematical solution of the diffusion equations subject to the boundary conditions.

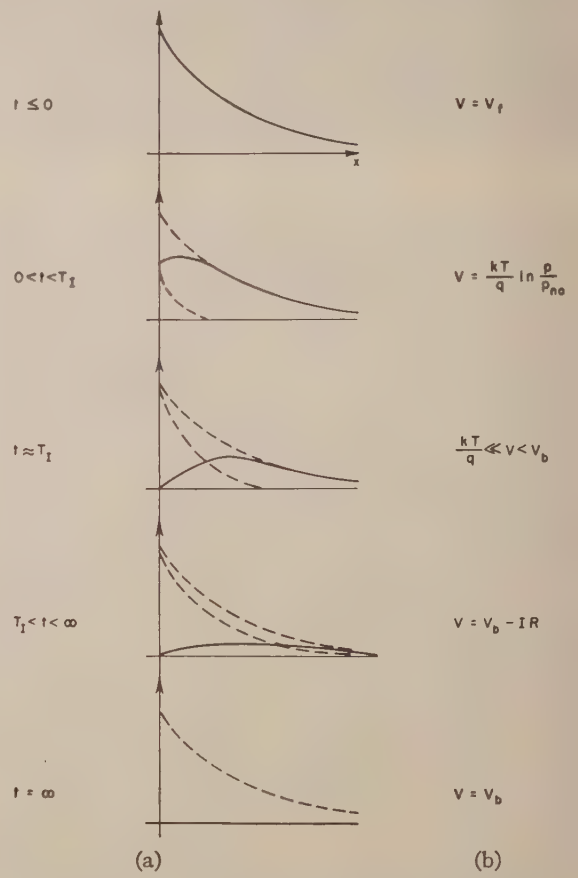

Fig. 3—(a) Hole density, and (b) junction potential, during switching process.

Mathematical Treatment[2]

Three cases are to be treated mathematically subject to the above criteria. They are the planar diode, the hemispherical diode, and the narrow-base diode. These are represented in Fig. 4 and discussed as treated.

A. Planar Diode

Planar junction with length of n-type material much greater than a diffusion length. (Fig. 4a. $W \gg L_p$.) The diffusion equation

$$\frac{\partial p}{\partial t} = D_p \frac{\partial^2 p}{\partial x^2} - \frac{p}{\tau_p} \tag{1}$$

may be rewritten to read

$$\frac{\partial p}{\partial T} = \frac{\partial^2 p}{\partial X^2} - p \tag{2}$$

[2] B. Lax and S. F. Neustadter (to be published) have treated the planar junction more rigorously than the treatment to follow, and their results are in good agreement with the approximations used in this paper.

where

$$T = \frac{t}{\tau_p} \quad \text{and} \quad X = \frac{x}{L_p}.$$

Since the initial distribution of holes at $t=0$ is a steady-state solution to the diffusion equation; a complete solution may be found by adding to this a new transient solution. The over-all solution for the constant-current phase should satisfy the boundary condition that $J(X=0)$ is constant and equal to $(-J_r)=(-V_b/AR_0)$,

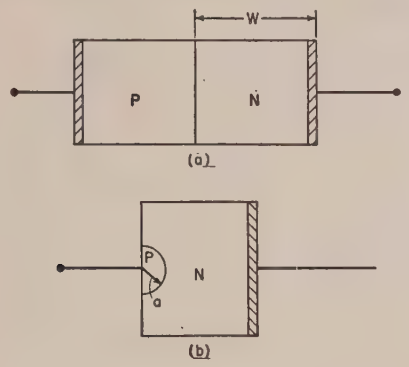

Fig. 4—Junction models used for calculation.

where A is the junction area. Now the current due to the steady state solution is J_f; therefore, the transient solution should satisfy the diffusion equation and the boundary condition that $J(X=0)=(J_f+J_r)$. When this solution, which we call $p_I(X, T)$ is subtracted from the initial steady-state solution, sketched in Fig. 3, required answer is obtained, since now $J(X=0)=J_f-(J_f+J_r)=-J_r$. By Laplace transformation method[3]

$$P_I = \frac{(J_f + J_r)L_p}{D_p} I(X, \sqrt{T})$$

where

$$I(x, z) \overset{\Delta}{=} \frac{2}{\sqrt{\pi}} \int_0^z e^{-u^2} e^{-x^2/4u^2} du, \tag{3}$$

$$I(0, z) = \text{erf } z$$

and

$$I(x, \infty) = e^{-x}.$$

The diode voltage for $0<T<T_I$ is thus given by

$$V = \frac{kT}{q} \ln \frac{p_0 - p_I(0, T)}{p_{n0}} \tag{4}$$

which is plotted in Fig. 5 for several values of J_r/J_f. As may be seen from the graph the voltage across the space charge region remains of the order of kT/q (0.026 v at room temperature) until a time very near T_I when it decreases very rapidly to minus infinity. It is this low value of voltage up to T_I which validates the assumption of constant current during the first phase of the diode transient. The time, T_I, is found by equating the transient solution at $X=0$, to the steady-state value of

[3] See Appendix for a discussion of the mathematical treatment.

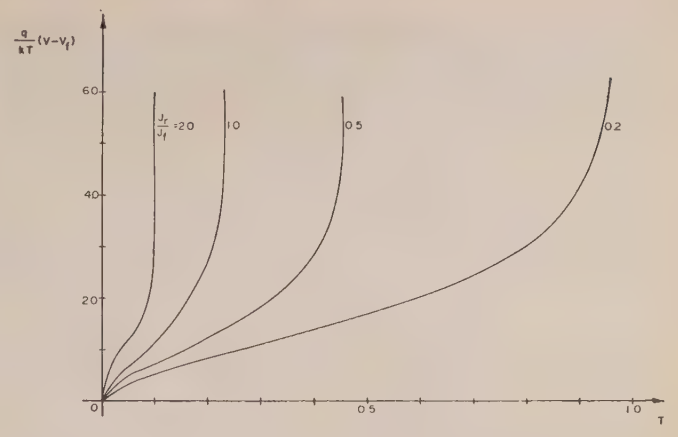

Fig. 5—Junction voltage during constant-current phase.

the hole density, p_0, giving

$$\text{erf } \sqrt{T_I} = \frac{1}{1 + J_r/J_f} \tag{5}$$

since

$$J_f = \frac{D_p P_0}{L_p}. \tag{6}$$

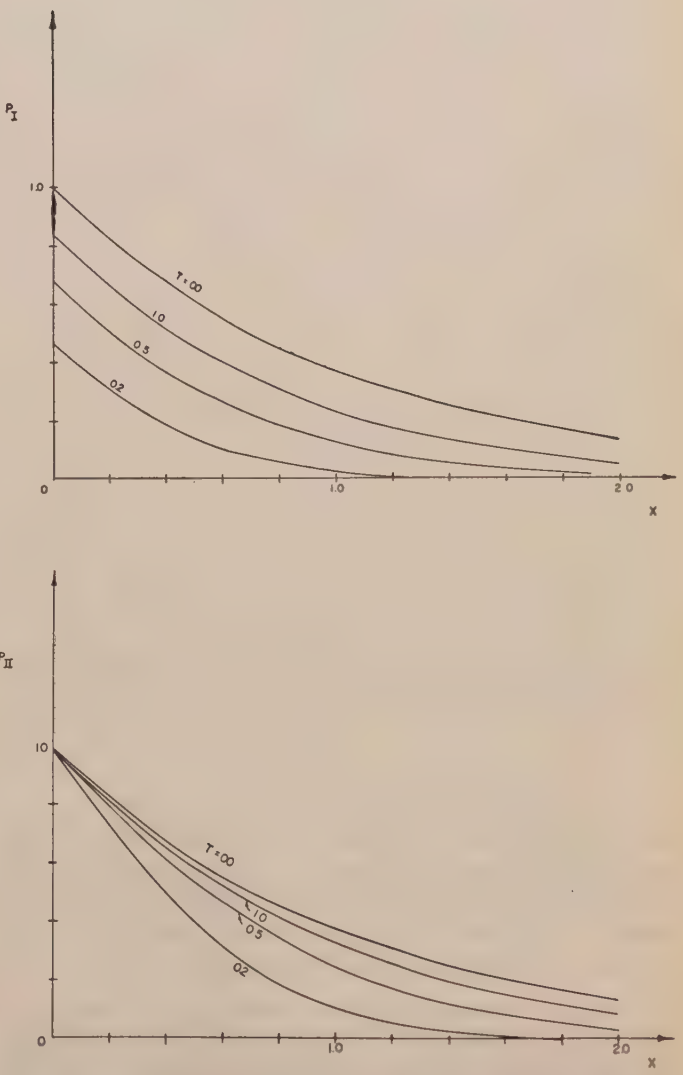

Fig. 6—The functions, p_I and p_{II}, for the planar junction.

For the remainder of the solution, the diffusion equation is solved subject to the boundary condition that $p(X=0)=0$ for all T, or specifically, $p_{II}(0, T)$ is constant and equal to p_0. This gives

$$p_{II} = \frac{J_f L_p}{D_p}\left[e^{-X} - I\left(X, \frac{X}{2\sqrt{T}}\right)\right] \tag{7}$$

and

$$J_{II} = -D_p\left(\frac{\partial p_{II}}{\partial X}\right)_{X=0} = J_f\left[\operatorname{erf}\sqrt{T} + \frac{e^{-T}}{\sqrt{\pi T}}\right]. \tag{8}$$

For the complete solution of the problem, p_{II} when $J_{II}=(J_f+J_r)$ is assumed to equal p_I at $T=T_I$. Actually, the magnitude and slope of p_I and p_{II} at $(X=0)$ are set equal at the end of the constant-current phase. Fig. 6 shows the normalized functions from which it may be

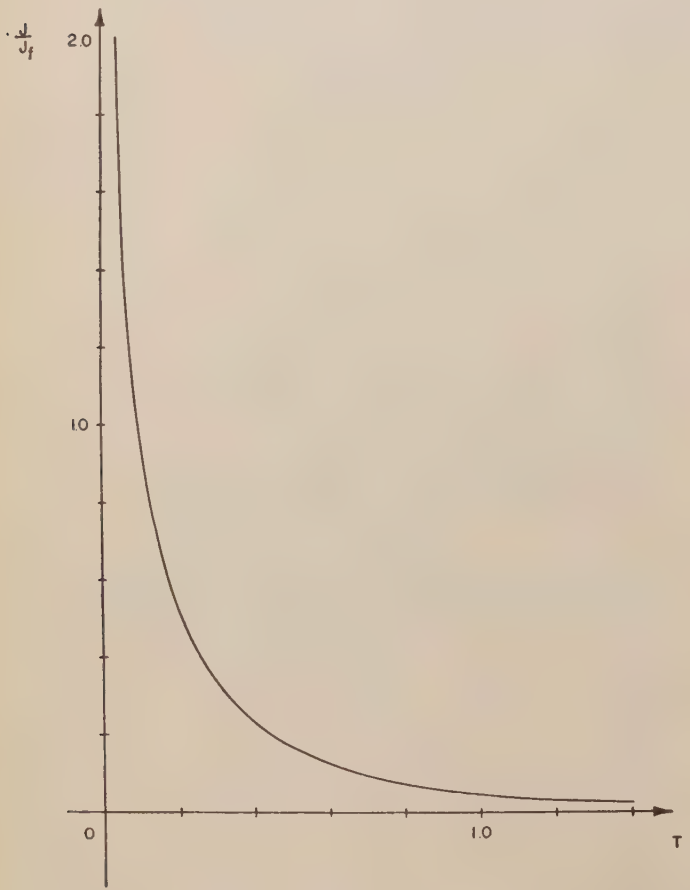

Fig. 7.—Junction reverse current after $T=T_I$.

seen that the above approximation is reasonable. If the true diffusion current $(J_{II}-J_f)$ is now plotted[4] as in Fig. 7, the reverse current after $T=T_I$ is given by this curve starting at $J=J_r$. The current previous to this time is, of course, J_r, from $T=0$ to $T=T_I$. (Note that the time scale may be shifted in the two functions p_I and p_{II}.)

B. Hemispherical Diode

Semi-infinite n-type material. (Fig. 4b) The diffusion equation in spherical co-ordinates becomes

$$\frac{\partial u}{\partial T} = \frac{\partial^2 u}{\partial R^2} - u \tag{9}$$

where $u=rp$, $R=r/L_p$, and $T=t/\tau_p$. The solution for the first phase is then

$$p_I = \frac{(J_f+J_r)}{D_p}\left\{\frac{A^2}{A^2-1} I(R-A, \sqrt{T}\,)\right.$$
$$+\frac{A}{A^2-1}\left[e^{T(1/A^2-1)}e^{(R-A/A)}\operatorname{erfc}\left(\frac{R-A}{2\sqrt{T}}+\frac{\sqrt{T}}{A}\right)\right.$$
$$\left.\left.+I\left(R-A, \frac{R-A}{2\sqrt{T}}\right) - I(R-A, \infty)\right]\right\} \tag{10}$$

where $A=a/L_p$. Now the forward current is given by

$$J_f = \frac{Dp_0}{L_p}\left(\frac{L_p}{a}+1\right) \tag{11}$$

therefore, at $R=A$, the hole density becomes zero when $p_{II}=p_0$ or

$$\frac{1}{1+J_r/J_f} = \frac{A}{A-1}\operatorname{erf}\sqrt{T}$$
$$+\frac{1}{1-A}\left[1-e^{T(1-A^2/A^2)}\operatorname{erfc}\left(\frac{\sqrt{T}}{A}\right)\right]. \tag{12}$$

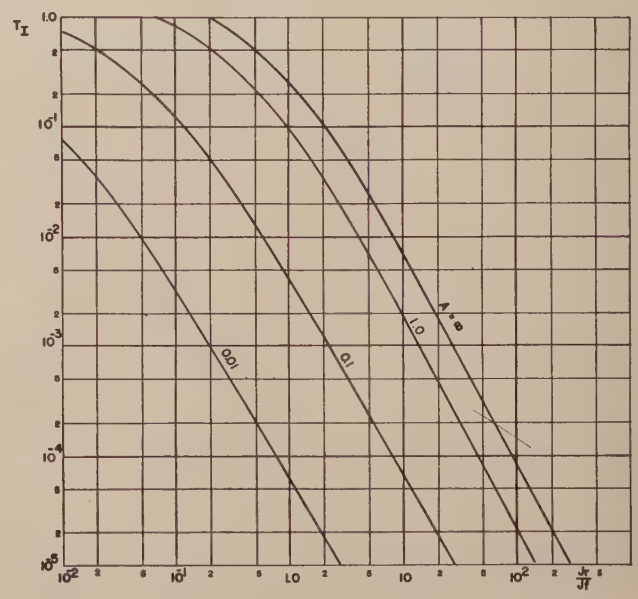

Fig. 8—Storage time, T_I, as a function of A and J_r/J_f.

This relationship is plotted in Fig. 8, which also gives the storage time for the planar case if A is set equal to infinity. The transient solution for the decaying phase is determined as in the linear case giving

[4] This particular equation has been derived previously by Lax and Neustadter, op. cit.; R. G. Shulman and M. E. McMahon, "Recovery currents in germanium p-n junction diodes," paper presented at the AIEE Winter General Meeting, New York, N. Y., January 20, 1953, also *Jour. Appl. Phys.*, vol. 24, p. 1267, 1953; and E. M. Pell, "Recombination rate in germanium by observation of pulsed reverse characteristics," *Phys. Rev.*, vol. 90, p. 228; 1953.

$$p_{II} = p_0 \frac{a}{r}\left[I(R-A,\infty) - I\left(R-A, \frac{R-A}{2\sqrt{T}}\right)\right] \quad (13)$$

and

$$J_{II} = \frac{D_p p_0}{a} + \frac{D_p p_0}{L_p}\left[\text{erf}\ \sqrt{T} + \frac{1}{\sqrt{\pi T}}\, e^{-T}\right]. \quad (14)$$

Subtracting J_f from the above equation gives the net diode reverse current

$$J_{II} - J_f = \frac{A}{A+1}\left[\text{erf}\ \sqrt{T} - 1 + \frac{1}{\sqrt{\pi T}}\, e^{-T}\right] \quad (15)$$

after use of (11). This is exactly the same equation as that plotted in Fig. 7, after multiplication by the factor $A/(A+1)$, and the method of determining the decay current is as in the previous solution. Fig. 9 shows the general curve for the decay time during the second phase as a function of the current ratio, J_r/J_f, and the radius of the contact, $A = a/L_p$. Here the decay time is defined as that in which the current falls from J_r to 10 per cent of J_r. This is the same time as that required for the diode-voltage to reach 90 per cent of the battery value.

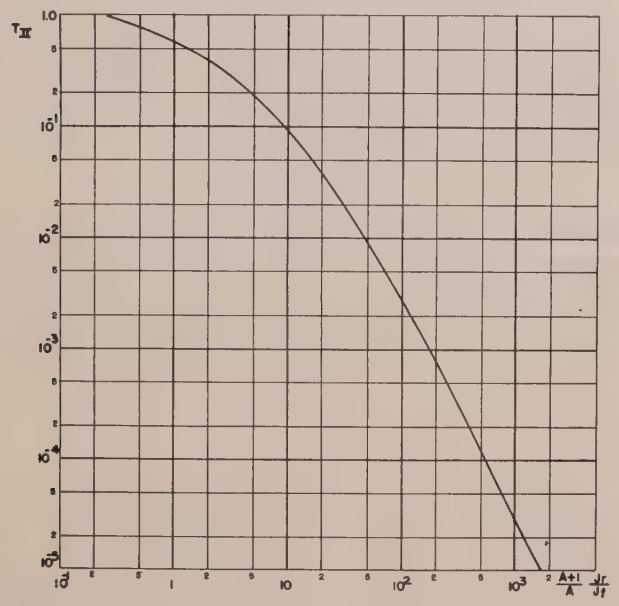

Fig. 9—Decay time, T_{II}, as a function of A and J_r/J_f.

C. Narrow-Base Diode

Width of n-type material, W; lifetime assumed infinite. (Fig. 4a. $W \ll L_p$.)

Here the boundary conditions will be slightly different than the two preceding treatments. If the ohmic contact to the diode is assumed to be a sink for holes the hole density at $x = W$ must be zero at all times. This is not generally true for all diode contacts but the calculation will be made on this basis since it is directly applicable to the junction transistor where the collector junction is, almost by definition, a sink for holes.

The diffusion equation may be modified to read:

$$\frac{\partial p}{\partial T} = \frac{\partial^2 p}{\partial X^2} \quad (16)$$

where $T = D_p t/W^2$ and $X = x/W$. By the previous methods the constant-current phase solution is found to be

$$p_I = \frac{(J_f + J_r)W}{D}\left(\frac{2}{\pi^2}\right)\sum_{n=0}^{\infty} \frac{(-1)^n(1 - e^{-\pi^2 T(n+1/2)^2})}{(n+\frac{1}{2})^2} \cdot \sin \pi(n+\tfrac{1}{2})(1-X) \quad (17)$$

and the solution for the decaying phase with the hole-density constant at $X = 0$, is

$$P_{II} = \frac{J_f W}{D}\left[1 - X + \frac{2}{\pi}\sum_{n=1}^{\infty} \frac{(-1)^n e^{-\pi^2 n^2 t}}{n} \cdot \sin \pi n(1-X)\right]. \quad (18)$$

These solutions are similar to those plotted in Fig. 6 except that they vanish at $X = 1$. The net decay current at the barrier and at the collector or ohmic contact, $(X = W)$, is plotted in Fig. 10.

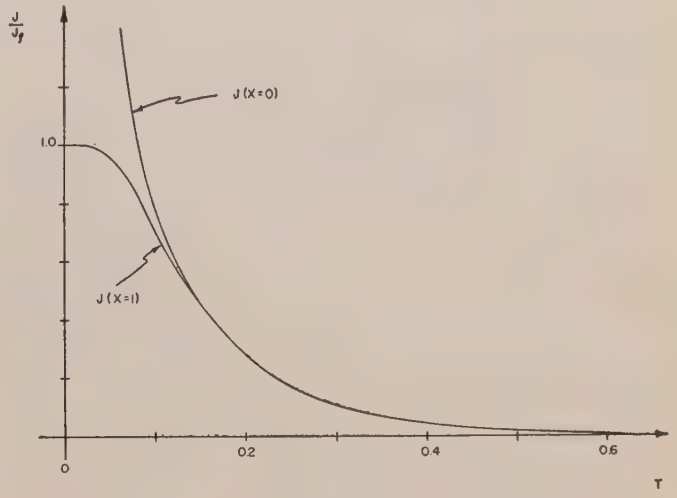

Fig. 10—Net current flow during decaying-current phase (narrow-base).

The decay time T_{II}, defined as before, is plotted in Fig. 11 with the constant-current phase time, T_I. Also shown on the same plot is T_C, which is the time required for the collector current in a junction transistor, $(J_{II}(W) - J_f)$, to fall to 10 per cent of J_f, the initial collector current, after the emitter is switched from forward to reverse. The time, T_C, includes the time, T_I, for the constant-current phase. For $J_r/J_f \ll 1$, T_C is obtained from the equation

$$J_I(W) = (J_f + J_r)\left[1 - \frac{2}{\pi}\sum_{n=0}^{\infty}(-1)^n \frac{e^{-\pi^2 T(n+1/2)^2}}{(n+\frac{1}{2})}\right] \quad (19)$$

which is the transient current at $X = W$ for the constant-current phase.

Discussion of Results

The calculations for the planar diode are seen to be a special case of the hemispherical diode with an infinite radius of curvature. Most interesting in the resultant curves is the rapid decrease of the storage times, T_I and T_{II}, with decreasing junction radius. It should be

emphasized that these calculations really give an upper limit on the switching times, since, as the radius decreases, the current density, in practical cases becomes so large that the recombination near the junction is no longer linear. The net result of this conductivity modulation should be a decrease in the effective lifetime caused by the large increase in majority carrier density required to neutralize the excess minority carriers. This does not mean that the optimum over-all design results are obtained with the smallest possible junction radius. It should be remembered that forward direction spreading resistance will increase not only with decreasing radius but also with decreasing lifetime, since effective radius for spreading resistance calculation is of the order of the junction radius plus diffusion length.

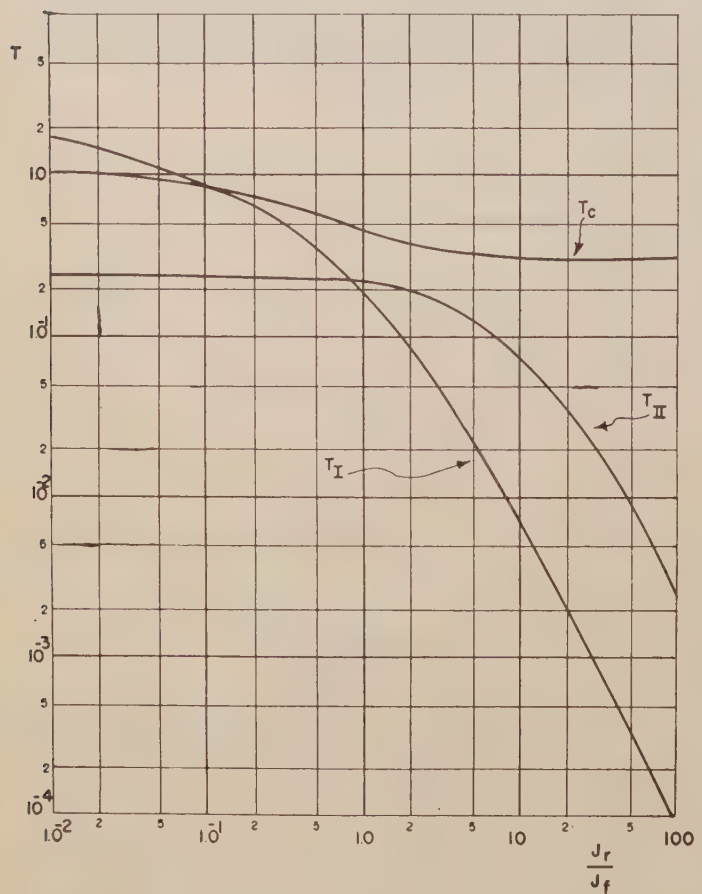

Fig. 11—Time constants, T_I, T_{II}, and T_C, for the narrow-base diode and junction transistor.

The resultant curves for the narrow-base diode are quite similar to the linear diode except for the scale. The normalized parameter, T, which was based on the lifetime in the first two cases is now based on the width of the semiconductor body. In fact, when considering the junction transistor application, the time variable may be written as

$$T = \frac{Dt}{W^2} = \left(\frac{D}{\pi W^2}\right)\pi t = \pi f_{c\cdot o}\cdot t \qquad (20)$$

where f_{co} is the frequency cutoff for α as calculated from the diffusion equation.[5] In this connection, it should be remembered that the T_C curves give the behavior of the current source αI_e while the actual collector voltage will also be a function of the collector impedance. Similarly, if the time constant of reverse emitter capacitance and circuit resistance is comparable to calculated switching time of emitter, solution is no longer accurate.

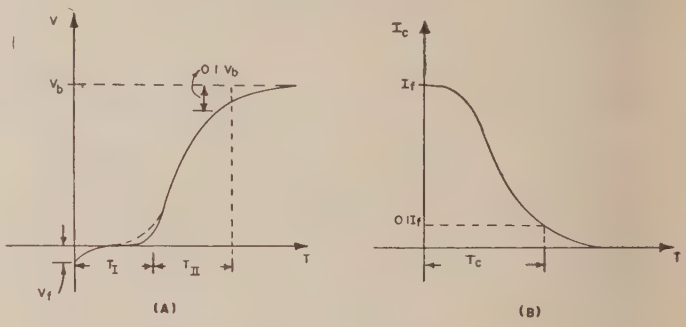

Fig. 12—Switching transients in junction diodes and transistors.

Conclusion

The voltage, V_D, across the diode in Fig. 1 will behave as shown in Fig. 12(a), where the times indicated are shown in Figs. 8, 9, and 11. If the diode is the emitter of a junction transistor, then the collector current will behave as in Fig. 11(b) where T_C is given in Fig. 10. The ratio J_r/J_f is also equal to V_B/I_fR_0.

Acknowledgment

The mathematical treatment was aided by the contributions of J. Keilson, S. F. Neustadter, G. C. Hunt, and the computing group of this laboratory. Of especial value were the suggestions and encouragement of E. Rawson, B. Lax, J. E. Thomas, and R. B. Adler.

Appendix

Mathematical Methods

The solution of the diffusion equation for the different boundary conditions was obtained by Laplace transformation in the time domain and solution of the resultant equation by standard techniques. For the linear and hemispherical diodes the appropriate inverse transforms were found in Magnus and Oberhettinger,[6] and Hameister.[7] The transforms for the narrow-base diode are again in Magnus and Oberhettinger.[6]

The integral of (3) of the text may be written

$$I(x, z) = \frac{\sqrt{\pi}}{4}\left[e^{-z}\,\mathrm{erfc}\left(\frac{x}{2z} - z\right) - e^{x}\,\mathrm{erfc}\left(\frac{x}{2z} + z\right)\right].$$

S. F. Neustadter of this laboratory is responsible for this form.

[5] W. Shockley, M. Sparks and G. K. Teal, "*P-n* junction transistors," *Phys. Rev.*, vol. 83, p. 151; 1951.

[6] W. Magnus and F. Oberhettinger, "Formulas and Theorems for the Special Functions of Mathematical Physics," Chelsea Pub. Co., New York, N. Y., p. 129 and p. 136; 1949.

[7] E. Hameister, "Laplace Transformation," Verlag von R. Oldenbourg, Munich and Berlin, Germany, p. 136; 1943.

Bounds Existing on the Time and Frequency Responses of Various Types of Networks*

ARMEN H. ZEMANIAN†, ASSOCIATE, IRE

Summary—**The transient response of fixed, lumped, linear, and stable networks is investigated and many bounds are shown to exist on the impulse and step responses of various classes of system functions. Conversely, if the impulse response is restricted in certain ways, bounds must then exist on the frequency response. These bounds, which are obtained by manipulating the Fourier transforms, have many practical implications. For instance, the response due to a unit impulse of current on a passive driving point impedance which has a shunting capacity, C, across its input terminals is bounded by $\pm(1/C)$ and the rise time for a low pass system of this form must be greater than rC where r equals the value of the impedance under dc conditions. Similar statements may be made for transfer functions satisfying certain mathematical restrictions. As other examples, more severe lower bounds have been found on the settling time for such systems. Also the overshoot or undershoot of the response to a step input of current for an RC driving point impedance cannot be greater than one hundred per cent.**

Introduction

IT IS A WELL KNOWN fact that linear networks cannot have arbitrary frequency characteristics.[1] For the case of minimum reactance or susceptance networks, many relations have been found between the real and imaginary parts of the system immittance and similar relations exist between the amplitude and phase functions in the minimum phase case. The positive real characteristic of a driving point immittance is another such example. The object of this investigation has been to discover an analogous set of restrictions on the time response of various types of networks.

The results in general indicate that the impulse and unit step responses of various classes of networks are bounded in many ways and, due to the reciprocal nature of the direct and inverse Fourier transforms, similar restrictions which may be placed on the time responses fix bounds on the corresponding frequency responses. These bounds which are defined in the forthcoming theorems, are not approximate relations but definite restrictions that the responses of the appropriate systems must obey.

A special case of theorem one has appeared previously. M. Nadler[2] formulates a similar problem mathematically and points out that, for a two-terminal, lowpass, passive network which has a certain amount of shunting capacity, C, between its terminals, the rise time from zero to the final value cannot be less than rC where r is the final value. It has been found in this investigation that such a restriction can be applied to any system function which satisfies a few mathematical conditions and furthermore a number of theorems of this type are proven. Finally it should be pointed out that some of these results and methods have analogies in probability theory and, in particular, to the conditions on the characteristic functions.[3,4]

Some Theorems Concerning Time and Frequency Responses

The Applicability of the Theorems

In the forthcoming discussion it is assumed that a unit step or impulse of current is impressed upon the network and a corresponding voltage response is taken as the output so that the system function is an impedance. However these theorems apply just as well if the input is a voltage, the output a current, and the system function an admittance. Similarly both input and output may be voltages or currents and the system function a transfer function.

Furthermore it is assumed that the input is applied at time, $t=0$, so that the unit impulse response, $W(t)$, and the unit step response, $A(t)$, are identically equal to zero for negative values of time. Also the systems are taken to be fixed, lumped, linear, and stable networks such that the system impedance, $Z(s)=N(s)/D(s)$, has no poles on the real frequency axis or in the right half plane. The real and imaginary parts of $Z(j\omega)$ are represented by the symbols $R(\omega)$ and $I(\omega)$, respectively.

Finally only those system impedances are considered for which the polynomial, $D(s)$, is of degree equal to or greater than the degree of the polynomial, $N(s)$. This restriction is needed to insure the convergence of certain integrals which will arise in the forthcoming proofs. Thus the system impedance may be expanded into the following two infinite series:

$$Z(s) = K + \frac{1}{Cs} + \frac{K_2}{s^2} + \frac{K_3}{s^3} + \cdots ; \qquad |s| \geqq p$$

$$Z(s) = r + k_1 s + k_2 s^2 + k_3 s^3 + \cdots ; \qquad |s| \leqq q.$$

Since these expansions will only be of interest around the origin and as s approaches infinity, let p be a number

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[1] H. W. Bode, "Network Analysis and Feedback Amplifier Design," D. Van Nostrand Co., New York, N. Y., 1st ed., chap. 13; 1945.

[2] M. Nadler, "Letter to the editor," PROC. I.R.E., vol. 38, p. 441; April, 1950.

[3] H. Cramér, "Random Variables and Probability Distributions," Cambridge Tracts in Mathematics and Mathematical Physics, Cambridge University, London, Eng., no. 36; 1937.

[4] H. Cramér, "Mathematical Methods of Statistics," Princeton University Press, 1st ed., chap. 10; 1946.

greater than one and also greater than the distance from the origin to the pole of $Z(s)$ furthest from the origin and let q be a number less than one and also less than the distance from the origin to the pole nearest the origin. In the forthcoming theorems, designating a system impedance by the symbol, $Z(s)$, or its corresponding unit impulse response by the symbol, $W(t)$, will mean that the system satisfies the foregoing restrictions.

The Development of the Theorems

Theorem 1: If $Z(s)$ is such that $K=0$ and $R(\omega)\geqq 0$ for all ω, then $W(t)\leqq(1/C)$.

Note: For a passive driving point impedance, the symbol, C, would represent the shunting capacity between the input terminals if there were no other purely capacitive path through the network. Also K is stipulated to be zero so that $\int_0^\infty R(\omega)d\omega$ converges.

Proof: The Fourier transform which relates the impulse response to the real part of the system function is given as follows:[5]

$$W(t) = \frac{2}{\pi}\int_0^\infty R(\omega)\cos\omega t d\omega. \tag{1}$$

Since $R(\omega)$ is never negative, by hypothesis:

$$|W(t)| = \left|\frac{2}{\pi}\int_0^\infty R(\omega)\cos\omega t d\omega\right| \leqq \frac{2}{\pi}\int_0^\infty R(\omega)d\omega.$$

Now by integrating $Z(s)$ around the right half s plane and making use of the fact that there are no poles in the right half s plane nor on the $j\omega$ axis, a value may be obtained for $\int_0^\infty R(\omega)d\omega$.[1] Thus:

$$\frac{2}{\pi}\int_0^\infty R(\omega)d\omega = \frac{1}{C}.$$

This completes the proof.

Theorem 2: If $Z(s)$ is such that $K=0$ and $(dR/d\omega)\leqq 0$ for $0\leqq\omega$ then $|W(t)|<(2r/\pi t)$ and $|W(t)|\leqq(1/C)$.

Proof: Since $\int_0^\infty R(\omega)d\omega=\pi/2C$, $R(\omega)$ vanishes as ω increases beyond all bound. Thus the condition that $(dR/d\omega)\leqq 0$ for $0\leqq\omega$ requires that $r\geqq R(\omega)\geqq 0$ for $0\leqq\omega$, the system being a lumped, linear, fixed, and stable one. Thus theorem one may be invoked.

The second part of the proof is based upon expression obtained by integrating transform (1) by parts.

$$W(t) = -\frac{2}{\pi t}\int_0^\infty \frac{dR}{d\omega}\sin\omega t d\omega.$$

Since $(dR/d\omega)\leqq 0$ for $0\leqq\omega$:

$$|W(t)| < -\frac{2}{\pi t}\int_0^\infty \frac{dR}{d\omega}d\omega = \frac{2r}{\pi t}.$$

This completes the proof.

Theorem 3: If $Z(s)$ is such that $I(\omega)\leqq 0$ for $0\leqq\omega$, then:

$$K \leqq A(t) < 2r - K.$$

[5] E. A. Guillemin, "Communications Networks," John Wiley & Sons, New York, N. Y., 1st ed., vol. II, chap. 11; 1935.

Note: If $r=0$ and if $I(\omega)\geqq 0$ then a similar theorem may be proved which states that $|A(t)|\leqq K$.

Proof: Use is made of the following Fourier transform:

$$A(t) = \frac{2}{\pi}\int_0^\infty \frac{I(\omega)}{\omega}\cos\omega t d\omega + r. \tag{2}$$

As before:

$$\left|\int_0^\infty \frac{I(\omega)}{\omega}\cos\omega t d\omega\right| \leqq \int_0^\infty \frac{|I(\omega)|}{\omega}d\omega.$$

Again by integrating $Z(s)/s$ around the right half s plane, a value may be obtained for $\int_0^\infty[|I(\omega)|/\omega]d\omega$. In this case, however, the path of integrating must be indented slightly to the right of the pole that exists at the origin.[1] Thus:

$$\int_0^\infty \frac{|I(\omega)|}{\omega}d\omega = \frac{\pi}{2}(r-K).$$

Furthermore since $\int_0^\infty[|I(\omega)|/\omega]d\omega$ is a positive quantity, $r\geqq K$. Thus:

$$K \leqq A(t) < 2r - K$$

and the proof is thereby complete.

By placing similar restrictions on the impulse response, bounds may be found on the frequency functions as follows.

Theorem 4: If $W(t)$ is such that $\int_0^\infty W(t)dt=r$ and $W(t)\geqq 0$ for all t, then $|Z(j\omega)|\leqq r$, $|R(\omega)|\leqq r$ and $|I(\omega)|<r$ for all ω.

Note: Stating that $\int_0^\infty W(t)dt=r$ is the same as stating that $K=0$. For if K were not zero, the impulse response would have a singularity at $t=0$ and this integral would not exist.

The proof of this theorem is the same as that of theorem one wherein use is made of the following Fourier integrals.[5]

$$Z(j\omega) = \int_{-\infty}^\infty W(t)e^{-j\omega t}dt. \tag{3}$$

$$R(\omega) = \int_0^\infty W(t)\cos\omega t dt. \tag{4}$$

$$I(\omega) = -\int_0^\infty W(t)\sin\omega t dt. \tag{5}$$

If the first derivative of the impulse response is never positive then a theorem similar to theorem two may be constructed for these frequency functions. Similarly if the second derivative is never negative then a set of three bounds may be found on each of these functions. The general case where the mth derivative of $W(t)$ is never negative or never positive is stated by Theorem 5. The proof of this theorem makes use of the expressions which are obtained by repeatedly integrating the expressions (3), (4), and (5) by parts, and inserting the appropriate initial and final values bearing in mind that $W(t)$ must vanish as t goes to infinity for these expressions to hold. The derivatives of $W(t)$ at $t=0$ are taken

to be the limiting values that these derivatives approach as t goes to zero through positive values. From expression 3:

$$Z(j\omega) = \sum_{k=0}^{n} \frac{1}{(j\omega)^{k+1}} \cdot \frac{d^k W}{dt^k}\Big]_{t=0} + \frac{1}{(j\omega)^{n+1}} \int_0^\infty \frac{d^{n+1}W}{dt^{n+1}} e^{-j\omega t} dt. \tag{6}$$

From expression 4:

$$R(\omega) = \sum_{k=1}^{n} \frac{(-1)^k}{\omega^{2k}} \cdot \frac{d^{2k-1}W}{dt^{2k-1}}\Big]_{t=0} + \text{remainder} \tag{7}$$

$$\text{where the remainder} = \frac{(-1)^n}{\omega^{2n}} \int_0^\infty \frac{d^{2n}W}{dt^{2n}} \cos \omega t dt = \frac{(-1)^{n+1}}{\omega^{2n+1}} \int_0^\infty \frac{d^{2n+1}W}{dt^{2n+1}} \sin \omega t dt.$$

From expression 5:

$$I(\omega) = \sum_{k=0}^{n} \frac{(-1)^{k+1}}{\omega^{2k+1}} \cdot \frac{d^{2k}W}{dt^{2k}}\Big]_{t=0} + \text{remainder} \tag{8}$$

$$\text{where the remainder} = \frac{(-1)^{n+1}}{\omega^{2n+1}} \int_0^\infty \frac{d^{2n+1}W}{dt^{2n+1}} \cos \omega t dt = \frac{(-1)^{n+2}}{\omega^{2n+2}} \int_0^\infty \frac{d^{2n+2}W}{dt^{2n+2}} \sin \omega t dt.$$

Theorem 5: If $W(t)$ is such that $\int_0^\infty W(t)dt = r$ then, when the mth derivative is never positive or never negative, the following set of bounds hold where $n = 1, 2, 3, \cdots, m$.

$$|Z(j\omega)| \leqq \left| \sum_{k=0}^{n-1} \frac{1}{(j\omega)^{k+1}} \cdot \frac{d^k W}{dt^k}\Big]_{t=0} \right| + \left| \frac{2}{(j\omega)^{n+1}} \cdot \frac{d^n W}{dt^n}\Big]_{t=0} \right|. \tag{9}$$

If the $2m$th derivative of $W(t)$ is never negative then the following set of bounds hold where $n = 1, 2, 3, \cdots, m$.

$$\sum_{k=1}^{n} \frac{(-1)^k}{\omega^{2k}} \cdot \frac{d^{2k-1}W}{dt^{2k-1}}\Big]_{t=0} - \left| \frac{1}{\omega^{2n}} \cdot \frac{d^{2n-1}W}{dt^{2n-1}}\Big]_{t=0} \right| \leqq R(\omega) \leqq \sum_{k=1}^{n} \frac{(-1)^k}{\omega^{2k}} \cdot \frac{d^{2k-1}W}{dt^{2k-1}}\Big]_{t=0} + \left| \frac{1}{\omega^{2n}} \cdot \frac{d^{2n-1}W}{dt^{2n-1}}\Big]_{t=0} \right|. \tag{10}$$

$$\sum_{k=0}^{n} \frac{(-1)^{k+1}}{\omega^{2k+1}} \cdot \frac{d^{2k}W}{dt^{2k}}\Big]_{t=0} - \frac{1}{\omega^{2n+1}} \cdot \frac{d^{2n}W}{dt^{2n}}\Big]_{t=0} \leqq I(\omega) \leqq \sum_{k=0}^{n} \frac{(-1)^{k+1}}{\omega^{2k+1}} \cdot \frac{d^{2k}W}{dt^{2k}}\Big]_{t=0} + \frac{1}{\omega^{2n+1}} \cdot \frac{d^{2n}W}{dt^{2n}}\Big]_{t=0}. \tag{11}$$

And if the $(2m+1)$th derivative of $W(t)$ is never positive then the following set of bounds hold where $n = 1, 2, 3, \cdots, m$.

$$\sum_{k=1}^{n} \frac{(-1)^k}{\omega^{2k}} \cdot \frac{d^{2k-1}W}{dt^{2k-1}}\Big]_{t=0} - \frac{1}{\omega^{2n+1}} \cdot \frac{d^{2n}W}{dt^{2n}}\Big]_{t=0} < R(\omega) < \sum_{k=1}^{n} \frac{(-1)^k}{\omega^{2k}} \cdot \frac{d^{2k-1}W}{dt^{2k-1}}\Big]_{t=0} + \frac{1}{\omega^{2n+1}} \cdot \frac{d^{2n}W}{dt^{2n}}\Big]_{t=0}. \tag{12}$$

$$\sum_{k=0}^{n} \frac{(-1)^{k+1}}{\omega^{2k+1}} \cdot \frac{d^{2k}W}{dt^{2k}}\Big]_{t=0} - \left| \frac{1}{\omega^{2n+2}} \cdot \frac{d^{2n+1}W}{dt^{2n+1}}\Big]_{t=0} \right| < I(\omega) < \sum_{k=0}^{n} \frac{(-1)^{k+1}}{\omega^{2k+1}} \cdot \frac{d^{2k}W}{dt^{2k}}\Big]_{t=0} + \left| \frac{1}{\omega^{2n+2}} \cdot \frac{d^{2n+1}W}{dt^{2n+1}}\Big]_{t=0} \right|. \tag{13}$$

Proof: The proof of this theorem proceeds in the same way as the proofs of the previous ones. Expression (9) is obtained from (6), expressions (10) and (12) are obtained from (7), and expressions (11) and (13) are obtained from (8). Then the specification of any even derivative as never negative or any odd derivative as never positive designates all the lower even derivatives as never negative and all the lower odd derivatives as never positive. Thus all the lower ordered bounds also hold and the proof is therefore complete.

These derivatives need not be never-negative or never-positive for bounds of this nature to hold. For instance, assuming that the magnitude of the nth derivative of $R(\omega)$ is integrable over the semi-infinite range, then the corresponding bound on the impulse response is of the form:

$$|W(t)| \leqq \frac{\text{constant}}{t^n}.$$

Finally by stipulating $R(\omega)$ as never-negative for $0 \leqq \omega$ and assuming that the impulse response dies out to within certain limits after a period of time, bounds on the impulse response stronger than the one given by theorem one may be found.

Theorem 6: If $R(\omega) \geqq 0$ for all ω and if $|W(t)| \leqq (\epsilon/C) < (1/C)$ for $t_\epsilon \leqq t$ and $W(0) = 1/C$, then $|W(t)|$ lies within the following set of bounds.

For

$$\frac{t_\epsilon}{2} \leqq t < t_\epsilon, \quad |W(t)| < (.637 + .638\epsilon)\frac{1}{C}; \tag{14}$$

for

$$\frac{t_\epsilon}{3} \leqq t < \frac{t_\epsilon}{2}, \quad |W(t)| < (.785 + 1.916\epsilon)\frac{1}{C}; \tag{15}$$

for

$$\frac{t_\epsilon}{4} \leqq t < \frac{t_\epsilon}{3}, \quad |W(t)| < (.907 + .493\epsilon)\frac{1}{C}; \tag{16}$$

and for

$$0 \leqq t < \frac{t_\epsilon}{4}, \quad |W(t)| \leqq \left[1 - \frac{t^2}{4t_\epsilon^2}(1 - \epsilon)\right]\frac{1}{C}. \tag{17}$$

Note: These bounds are illustrated in Fig. 1, following page, for a negligible ϵ.

Proof: Consider the following inequality which holds since $R(\omega) \geqq 0$ by hypothesis:

$$|W(t)| \leqq \frac{2}{\pi} \int_0^\infty R(\omega) |\cos \omega t| d\omega. \tag{18}$$

Inserting the inequality:

$$|\cos \omega t| \leqq \frac{3}{4} + \frac{1}{4} \cos 2\omega t$$

and integrating, the following is obtained:

$$|W(t)| \leqq \frac{3}{4} \cdot \frac{1}{C} + \frac{1}{4} |W(2t)|.$$

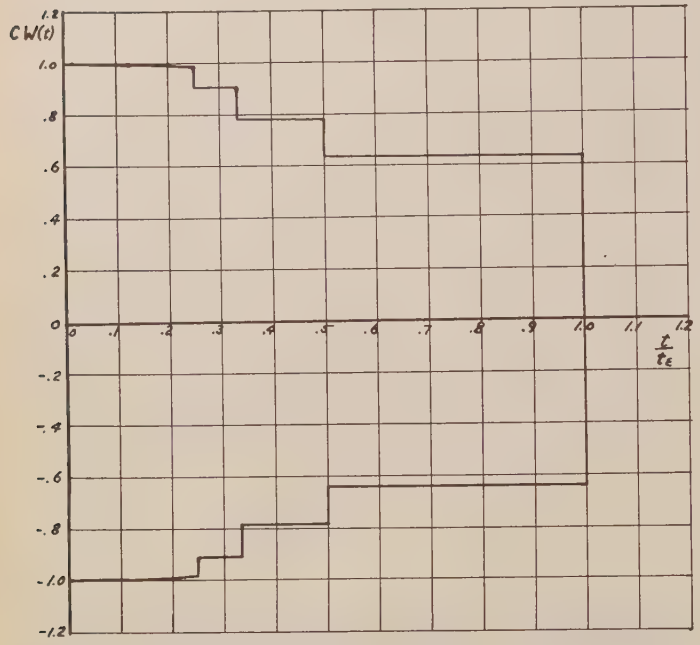

Fig. 1—Illustration of the bounds defined in Theorem 6 for a negligible ϵ.

However, for

$$\frac{t_\epsilon}{2} \leqq t < t_\epsilon, \qquad |W(2t)| \leqq \frac{\epsilon}{C}.$$

So:

$$|W(t)| < \frac{1}{C}\left[1 - \frac{1}{4}(1-\epsilon)\right].$$

Repeating the argument n times:
For

$$\frac{t_\epsilon}{2^n} \leqq t < \frac{t_\epsilon}{2^{n-1}}, \quad |W(t)| < \frac{1}{C}\left[1 - \frac{1}{4^n}(1-\epsilon)\right]. \tag{19}$$

But for

$$t < \frac{t_\epsilon}{2^{n-1}}, \qquad \frac{1}{4} \cdot \frac{t^2}{t_\epsilon^2} < \frac{1}{4^n}$$

so the inequality (19) may be replaced by (20):

$$|W(t)| < \frac{1}{C}\left[1 - \frac{t^2}{4t_\epsilon^2}(1-\epsilon)\right]. \tag{20}$$

Since n is an arbitrary positive integer, this holds for the interval $0 < t < t_\epsilon$. Furthermore $W(0) = 1/C$, and so the proof for the bound (17) is complete.

For bounds (14) and (16), the $|\cos \omega t|$ is expanded in a Fourier series and inserted into expression (18).

$$|W(t)| \leqq \frac{2}{\pi}\int_0^\infty \left[\frac{2}{\pi} R(\omega) - \frac{4}{\pi}\sum_{n=1}^{\infty} \frac{(-1)^n}{4n^2-1} R(\omega) \cos 2n\omega t\right] d\omega. \tag{21}$$

Since the infinite series in expression (21) converges uniformly it may be integrated term by term. Then the following inequality may be obtained.

$$|W(t)| < \frac{2}{\pi}\left[\frac{1}{C} + 2\sum_{n=1}^{\infty} \frac{1}{4n^2-1} |W(2nt)|\right].$$

Now for $(t_\epsilon/2) \leqq t < t_\epsilon$ and for $n = 1, 2, 3, \cdots$, $|W(2nt)| \leqq (\epsilon/C) < (1/C)$. Furthermore:

$$\sum_{n=1}^{\infty} \frac{1}{4n^2-1} < .5010.$$

Thus:

$$|W(t)| < \frac{1}{C}(.637 + .638\epsilon). \tag{14}$$

Similarly it can be shown that, for $(t_\epsilon/4) \leqq t < (t_\epsilon/3)$,

$$|W(t)| < \frac{1}{C}(.907 + .493\epsilon). \tag{16}$$

Finally to develop the bound (15), the $\cos \omega t$ in expression (1) is replaced by $f_1(\omega t)$.

$$f_1(\omega t) = \begin{cases} \cos \omega t & \text{for} \quad 0 \leqq \omega t \leqq \dfrac{\pi}{2} \\ -\dfrac{2}{\pi} x + 1 & \text{for} \quad \dfrac{\pi}{2} \leqq \omega t \leqq \pi \end{cases}$$

$$f_1(\omega t) = f_1(-\omega t); \qquad f_1(\omega t) = f_1(\omega t + 2\pi).$$

Expanding this into a Fourier series and proceeding exactly as before, the following is obtained:
For

$$\frac{t_\epsilon}{3} \leqq t < \frac{t_\epsilon}{2}, \qquad W(t) < \frac{1}{C}(.785 + 1.916\epsilon).$$

To develop a symmetrical set of lower bounds for this interval, the function $g_1(\omega t) = -f_1(\omega t + \pi)$ may be applied in exactly the same way. This completes the proof.

An analogous set of theorems may be constructed on $R(\omega)$, $Z(j\omega)$ and $A(t)$. This has been done and the whole matter is discussed in much greater detail elsewhere.[6]

Some Physical Interpretations of These Theorems

These theorems have practical implications some of which are listed below.

1. The condition that $R(\omega) \geqq 0$ for all ω is the characteristic which differentiates passive driving-point impedances from active ones. So the impulse responses of such systems are limited by the bounds in Theorems 1 and 6, for any two terminal network must have some

[6] A. H. Zemanian, "Investigation of the Transient Response of Linear Systems," New York Univ., Dept. of Elec. Eng., doctoral dissertation, chaps. 3–4; 1953.

stray shunting capacity across its terminals. Furthermore this result means that the slope of the unit-step response cannot be greater than $1/C$ and so the rise time from zero to the final value, r, cannot be less than rC. Nor can the rise time from $0.1r$ to $0.9r$ be less than $0.8rC$. These restrictions hold for any system function which satisfies the conditions of these theorems.

2. From Theorem 3, it can be seen that any RC two terminal network cannot have a percentage overshoot or undershoot in its voltage response to a step input of current greater than one hundred per cent. Furthermore, if the driving point impedance of this RC network assumes a value of K as ω goes to infinity then the percentage overshoot or undershoot must be less than $100\ (r-K)/r$. It could also be said for the current in an RL network when a voltage source is applied.

3. Theorem 4 yields the result that if $|Z(j\omega)|$ or $|R(\omega)|$ at any angular velocity, $\omega \neq 0$, is greater than its value at $\omega=0$, then $W(t)$ must be negative at some time. In other words, the unit step response cannot be monotonic increasing. Similarly these frequency responses must satisfy the bounds of Theorem 6 if the unit step response is to be monotonic increasing.

4. Inserting the value, $m=1$, into expression (10) of Theorem 5, it can be seen that any impulse response that vanishes with time, has a finite, nonzero initial value and is always concave upward will have a Fourier transform whose real part is never negative. This fact is of use if one is interested in designing a passive driving-point impedance from prescribed transient behavior.

5. System functions which satisfy the conditions of Theorem 6 will have corresponding step responses whose rise times are greater than the lower bounds mentioned in part one of this section. Similarly if the settling time is defined as follows then even stronger bounds may be constructed for this settling time.

Definition: The "settling time to ϵ," $\tau_{s\epsilon}$, is the least time beyond which the step response remains within $r(1-\epsilon)$ of its final value, r, and the impulse response remains within $\epsilon(1/C)$ of its initial value, $1/C$, where ϵ is a positive quantity less than unity. This assumes that the input functions are impressed at time, $t=0$.

It has been shown that for appropriate systems the unit impulse response must remain within the bounds indicated in Theorem 6. Up to time, t_ϵ, the maximum value that the unit step response may attain is the area between the upper set of bounds and the time axis. Moreover, if the time, t_ϵ, from which these bounds are constructed is so taken that the unit step response is always within $|1-\epsilon|r$ of its final value (that is, $t_\epsilon=\tau_{s\epsilon}$) then the maximum final value for a given C and $\tau_{s\epsilon}$ may be calculated. This, in turn, yields a lower bound on $\tau_{s\epsilon}$ for a given r and C.

$$\tau_{s\epsilon} > \frac{1-\epsilon}{.774 + .681\epsilon} rC. \tag{22}$$

By extending bound (16) over the interval of bound (15) (this is permissible as can be seen from the proof of Theorem 6) a stronger lower bound on $\tau_{s\epsilon}$ is obtained for larger ϵ.

$$\tau_{s\epsilon} > \frac{1-\epsilon}{.794 + .444\epsilon} rC. \tag{23}$$

Similarly extending bound (14) up to $\tau_{s\epsilon}/2$:

$$\tau_{s\epsilon} > \frac{1-\epsilon}{.808 + .329\epsilon} rC. \tag{24}$$

And finally if the bound (14) is extended up to $\tau_{s\epsilon}$:

$$\tau_{s\epsilon} > \frac{1-\epsilon}{.916 + .0833\epsilon} rC. \tag{25}$$

The strongest lower bound on $\tau_{s\epsilon}$ obtained by combining (22) through (25) in sections is indicated in Fig. 2.

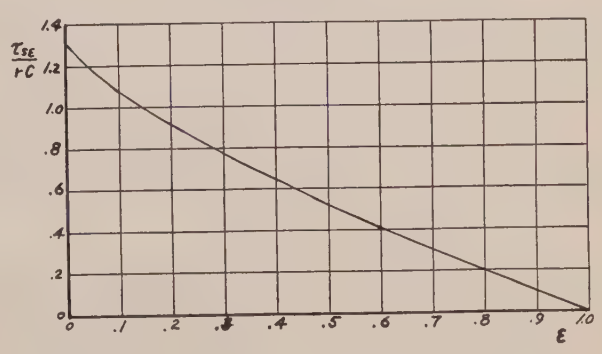

Fig. 2—A lower bound on the settling time to ϵ.

Acknowledgment

The writer is indebted to Professors J. H. Mulligan, Jr., and C. F. Rehberg, of New York University, for the many valuable suggestions and criticisms extended throughout the development of this work.

Appendix

Notations and Symbols

$A(t)$ = the response to a unit step function applied at time, $t=0$.
C = the capacity existing between the two terminals of a passive driving-point impedance.
ϵ = a positive real number less than one.
$I(\omega)$ = the imaginary part of a system function for real frequencies.
K = the constant multiplier of a system function having an equal number of poles and zeros.
$r=Z(0)$ = the resistance of a system function under dc conditions.
$R(\omega)$ = the real part of a system function for real frequencies.
s = a complex variable.
t = time.
τ = time.
$\tau_{s\epsilon}$ = the settling time to ϵ of the time response of a system as defined in the paper.
$W(t)$ = the response to a unit impulse function applied at time, $t=0$.
ω = angular frequency $=(2\pi)$ (frequency).
$Z(s)$ = a system impedance.

Matrix Analysis of Multi-Terminal Transducers*

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Summary—A multi-terminal transducer (m.t.t.) is a network with a set of n input terminals and a set of n output terminals. The properties of m.t.t.'s and their representation by matrices are a generalization of four-pole theory.

The m.t.t. is represented by its impedance, admittance or transfer matrices. The special properties of m.t.t.'s that exhibit symmetry, reciprocity or both properties are discussed. The representation is *definite* or *indefinite* according to whether the voltage reference terminals are specified or not.

Methods are derived to write down the transfer matrices of given m.t.t.'s. Some simple networks may be analyzed by inspection; more complicated networks may be treated as a cascade of simpler ones

Introduction

MATRIX ANALYSIS OF linear four-poles has received extensive treatment in the literature,[1–5] and the references cited are just a few out of a much longer list. The treatment, originally restricted to passive networks obeying the reciprocity relation, was later extended to networks containing vacuum tubes[6–8] and gyrators.[9,10]

A four-pole being a network with one pair of input terminals and one pair of output terminals, there are two obvious ways of generalizing four-pole treatment: first, networks having many *pairs* of terminals, and second, networks having a set of n input terminals and a set of n output terminals, with $n>2$. The first generalization occurs in the treatment of branches or junctions in transmission lines.[11,12] The second generalization is the subject of this paper.

A network with n input terminals and n output terminals, with a one-one correspondence between both sets of terminals, will be termed an *n-terminal transducer*. A four-pole is thus a 2-terminal transducer. Some examples of multi-terminal transducers (henceforth abbreviated to m.t.t.) are: sections of multiwire transmission lines, sections of a distributed amplifier, or sections of coupled transmission lines.

M.t.t.'s have been treated by Rice,[13] using matrix theory. Our approach is different from Rice's in the following points, which make the discussion more general.

1. The discussion is not restricted to networks that obey the reciprocity relation.
2. In the paper by Rice, "the leads marked 0 play a special role in that all the voltages are measured with respect to them, and the currents which they carry are the sum of the currents flowing into or out of the remaining terminals."[14] The treatment outlined in this paper enables any terminal to be specified as ground terminal, or even to leave the ground terminal unspecified.
3. Methods are derived to write down the matrices of given networks by inspection, similar to those of ordinary four-poles.[5]

For the sake of completeness some of the results due to Rice are included.

The reader is referred to any standard textbook[15] for the explanation of matrix manipulations. The following notations will be employed:

Square or rectangular matrices will be denoted by bold-face capitals: $\mathbf{Y}$; column matrices by bold-face lower case letters: $\mathbf{v}$. The submatrices of partitioned matrices will be denoted by subscripts

$$\mathbf{Y} = \left\| \begin{matrix} \mathbf{Y}_{11} & \mathbf{Y}_{12} \\ \mathbf{Y}_{21} & \mathbf{Y}_{22} \end{matrix} \right\|, \qquad \mathbf{v} = \left\| \begin{matrix} \mathbf{v}_1 \\ \mathbf{v}_2 \end{matrix} \right\|,$$

except in one case where, because of their repeated separate occurrence, and in analogy with the notation for four-poles, the submatrices will be denoted by different capitals, as in (2).

Single elements of matrices or submatrices will be denoted by superscripts:

$$\mathbf{v}_1 = \left\| \begin{matrix} v_1{}^1 \\ v_1{}^2 \\ \cdot \\ \cdot \\ v_1{}^n \end{matrix} \right\|.$$

* Decimal classification: R143. Original manuscript received by the IRE, July 7, 1953.

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[3] L. Brillouin, "Wave Propagation in Periodic Structures," chap. IX, "Matrices and the propagation of waves along an electric line," McGraw-Hill Book Co., Inc., New York, N. Y.; 1946.

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[5] P. I. Richards, "Applications of matrix algebra to filter theory," PROC. I.R.E., vol. 34, p. 145P; March, 1946.

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[9] B. D. H. Tellegen, "The Synthesis of Passive, Resistanceless Four-Poles that May Violate the Reciprocity Relation," Phillips Res. Rep., Eindhoven, Netherlands, vol. 3, p. 321; October, 1948.

[10] B. D. H. Tellegen and E. Klauss, "The Parameters of a Passive Four-Pole that May Violate the Reciprocity Relation," Phillips Res. Rep., Eindhoven, Netherlands, vol. 5, p. 81; April, 1950.

[11] C. G. Montgomery, R. H. Dicke, and E. M. Purcell, "Principles of Microwave Circuits," Radiation Laboratory Series No. 8, McGraw-Hill Book Co., Inc., New York, N. Y.; 1948.

[12] N. Marcuvitz, "Waveguide Handbook," Radiation Laboratory Series No. 10, McGraw-Hill Book Co., Inc., New York, N. Y.; 1951.

[13] S. O. Rice, "Steady state solutions of transmission line equations," *Bell Sys. Tech. Jour.*, vol. 20, p. 131; April, 1941.

[14] *Ibid.*, p. 151.

[15] E. A. Guillemin, "The Mathematics of Circuit Analysis," John Wiley and Sons, New York, N. Y.; 1949.

In order to conserve space, a column matrix will sometimes be written in a row, with special brackets:

$$\mathbf{v} = \{\mathbf{v}_1, \mathbf{v}_2\} = \{v_1^1, v_1^2 \cdots v_1^n, v_2^1, v_2^2 \cdots v_2^n\}.$$

$\mathbf{A}'$ denotes the transpose of $\mathbf{A}$. The unit matrix, of any order, will be denoted by $\mathbf{I}$, and its order will either be obvious from the context, or explicitly specified. All other notations, if not self-explanatory, will be defined on their first occurrence.

The Linear Multi-Terminal Transducer

Consider a network with $2n$ terminals (Fig. 1), of which n terminals are specified as input terminals (conventionally placed on the left) and n output terminals (on the right). Furthermore, let there be a one-one correspondence between the set of input terminals and the set of output terminals, so that a numbering $1, 2, \cdots, n$ at the input would unambiguously define the numbering at the output. Let v_1^p denote the voltage of the p-th input terminal, referred to any arbitrary input terminal, and i_1^p the current *into* it; let v_2^q be the voltage of the q-th output terminal, referred to an arbitrary output terminal, and i_2^q the current *out* of it.

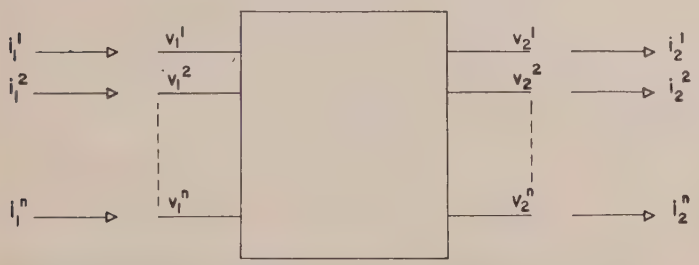

Fig. 1—Multi-terminal transducer.

In a *linear* network, provided there is any relation between $\mathbf{v}_1$, $\mathbf{v}_2$, $\mathbf{i}_1$, and $\mathbf{i}_2$, it may be expressed by

$$\left\| \begin{matrix} \mathbf{v}_1 \\ \mathbf{i}_1 \end{matrix} \right\| = \mathbf{K} \left\| \begin{matrix} \mathbf{v}_2 \\ \mathbf{i}_2 \end{matrix} \right\|. \tag{1}$$

$\mathbf{K}$ is a square matrix of order $2n \times 2n$, whose elements are independent of the $\mathbf{v}$ and $\mathbf{i}$. Nothing is assumed, *a priori*, about the *transfer matrix* $\mathbf{K}$, except its existence; it may be singular, or not unique, or exhibit any irregular properties. This would not be surprising, as the voltages are not completely and uniquely defined.

Make the following assumptions about $\mathbf{v}$ and $\mathbf{i}$:

(*a*) The sum of all the input currents is zero; the sum of all the output currents is zero.

(*b*) The currents do not depend on the voltages of the various terminals, only on voltage *differences* between pairs of terminals belonging to the same set. Hence, an arbitrary voltage v_1^0 (or v_2^0) may be added to all the input (output) voltages, without changing the currents.

These assumptions will put some restrictions on the elements of $\mathbf{K}$. Further analysis will be facilitated by partitioning $\mathbf{K}$ into four $n \times n$ submatrices,

$$\mathbf{K} = \left\| \begin{matrix} \mathbf{A} & \mathbf{B} \\ \mathbf{C} & \mathbf{D} \end{matrix} \right\|, \tag{2}$$

so that (1) is equivalent to

$$\begin{aligned} \mathbf{v}_1 &= \mathbf{A}\mathbf{v}_2 + \mathbf{B}\mathbf{i}_2 \\ \mathbf{i}_1 &= \mathbf{C}\mathbf{v}_2 + \mathbf{D}\mathbf{i}_2. \end{aligned} \tag{3}$$

The elements of $\mathbf{A}$ and $\mathbf{D}$ are dimensionless numbers, $\mathbf{B}$ is an impedance matrix, and $\mathbf{C}$ is an admittance matrix.

The current into the p-th input terminal is

$$i_1^p = \sum_r C^{pr} v_2^r + \sum_s D^{ps} i_2^s \quad (r, s = 1, 2, \cdots, n). \tag{4}$$

(In this section, all summations are over the range 1 to n of the index below the summation sign.)

Let the output terminals be open, so that $i_2^s = 0$ for all s. The sum of all the input currents, according to (4) and assumption (*a*), will be

$$\sum_p i_1^p = \sum_p \sum_r C^{pr} v_2^r = \sum_r \left(\sum_p C^{pr} \right) v_2^r = 0. \tag{5}$$

If this is to be true for any $\mathbf{v}_2$,

$$\sum_p C^{pr} = 0 \quad \text{for all } r. \tag{6}$$

On the other hand, suppose all the output terminals are shorted together, then all the v_2^r are equal; an appropriate constant may be added to all of them (assumption (*b*)) to make them all zero. The sum of the input currents is then

$$\sum_p i_1^p = \sum_p \sum_s D^{ps} i_2^s = \sum_s \left(\sum_p D^{ps} \right) i_2^s = 0. \tag{7}$$

This however, does not imply that all the $\sum_p D^{ps} = 0$, because $\sum_s i_2^s = 0$. The only necessary implication is that

$$\sum_p D^{ps} = \text{the same for all } s. \tag{8}$$

Now we take the general case (4) and add an arbitrary voltage v_2^0 to all the v_2^r:

$$i_1^p = \sum_r C^{pr}(v_2^r + v_2^0) + \sum_s D^{ps} i_2^s$$

$$i_1^p = \sum_r C^{pr} v_2^r + \left(\sum_r C^{pr} \right) v_2^0 + \sum_s D^{ps} i_2^s. \tag{9}$$

The i_1^p and i_2^s in (9) are the same as in (4), under assumption (*b*). Subtracting (4) from (9),

$$\left(\sum_r C^{pr} \right) v_2^0 = 0, \tag{10}$$

and, as v_2^0 is arbitrary,

$$\sum_r C^{pr} = 0 \text{ for all } r. \tag{11}$$

The voltage of the q-th input terminal, according to (3), is

$$v_1{}^q = \sum_r A^{qr} v_2{}^r + \sum_s B^{qs} i_2{}^s. \tag{12}$$

Add a voltage $v_1{}^0$ to all the imput voltages, and $v_2{}^0$ to all the output voltages,

$$v_1{}^q + v_1{}^0 = \sum_r A^{qr}(v_2{}^r + v_2{}^0) + \sum_s B^{qs} i_2{}^s$$

$$= \sum_r A^{qr} v_2{}^r + \left(\sum_r A^{qr} \right) v_2{}^0 + \sum_s B^{qs} i_2{}^s. \tag{13}$$

Subtracting (12) from (13),

$$v_1{}^0 = \left(\sum_r A^{qr} \right) v_2{}^0 \text{ for all } q, \tag{14}$$

hence,

$$\sum_r A^{qr} = \text{the same for all } q. \tag{15}$$

Equations (6), (8), (11), and (15) are the restrictions imposed on $\mathbf{K}$, due to assumptions (*a*) and (*b*) above.

To summarize the results: If no terminal is specified as the voltage reference point, the voltages and currents have to be restricted according to assumptions (*a*) and (*b*), and this results in the following:

1. The sum of all the elements in any row of $\mathbf{A}$ is the same.
2. The sum of all the elements in any column of $\mathbf{D}$ is the same.
3. The sum of all the elements in any row or column of $\mathbf{C}$ is zero.

This matrix $\mathbf{K}$ will be referred to as the *indefinite transfer matrix*[16] because the voltage reference terminal is left undefined.

Specifying the Voltage Reference Terminal

Let $\mathbf{K}$ be the indefinite transfer matrix of a m.t.t.' and let the n-th input and output terminals be specified as the reference terminals for the input and output voltages, respectively. We now have $(n-1)$ independent input voltages and $(n-1)$ output voltages, which are the differences between the voltages of the various terminals and those of the respective reference terminals. The "new" voltages constitute $(n-1)$-rowed columns $\bar{\mathbf{v}}_1$ and $\bar{\mathbf{v}}_2$, respectively.

In the following discussion assume that $n=4$. The extension of the treatment to any other value of n is obvious.

The voltages $\mathbf{v}$ and $\bar{\mathbf{v}}$ are related by

$$v^r - v^4 = \bar{v}^r \quad (r = 1, 2, 3), \tag{16}$$

but we may add a voltage $-v^4$ to all the components of $\mathbf{v}$, so that

$$\left.\begin{aligned} v^r &= \bar{v}^r \quad (r = 1, 2, 3) \\ v^4 &= 0. \end{aligned}\right\} \tag{17}$$

The last expression, written in matrix form, is

$$\mathbf{v} = \mathbf{P}\bar{\mathbf{v}}, \quad \text{where} \quad \mathbf{P} = \begin{Vmatrix} 1 & 0 & 0 \\ 0 & 1 & 0 \\ 0 & 0 & 1 \\ 0 & 0 & 0 \end{Vmatrix}. \tag{18}$$

The "old" currents $\mathbf{i}$ (4 components) may also be expressed by "new" currents $\bar{\mathbf{i}}$ (3 components),

$$\left.\begin{aligned} i^s &= \bar{i}^s \quad (s = 1, 2, 3) \\ i^4 &= -\bar{i}^1 - \bar{i}^2 - \bar{i}^3 \end{aligned}\right\}, \tag{19}$$

with the last line following from assumption (*b*). In matrix notation

$$\mathbf{i} = \mathbf{Q}\bar{\mathbf{i}}, \quad \text{where} \quad \mathbf{Q} = \begin{Vmatrix} 1 & 0 & 0 \\ 0 & 1 & 0 \\ 0 & 0 & 1 \\ -1 & -1 & -1 \end{Vmatrix}. \tag{20}$$

The problem now is to find a transfer matrix $\bar{\mathbf{K}}$ that relates the input and output quantities defined above.

The column $\{\mathbf{v}, \mathbf{i}\}$ may be transformed by

$$\begin{Vmatrix} \mathbf{v} \\ \mathbf{i} \end{Vmatrix} = \begin{Vmatrix} \mathbf{P} & 0 \\ 0 & \mathbf{Q} \end{Vmatrix} \times \begin{Vmatrix} \bar{\mathbf{v}} \\ \bar{\mathbf{i}} \end{Vmatrix}, \tag{21}$$

which is the same as (18) and (20) written together. Equation (1) then becomes

$$\begin{Vmatrix} \mathbf{P} & 0 \\ 0 & \mathbf{Q} \end{Vmatrix} \times \begin{Vmatrix} \bar{\mathbf{v}}_1 \\ \bar{\mathbf{i}}_1 \end{Vmatrix} = \mathbf{K} \times \begin{Vmatrix} \mathbf{P} & 0 \\ 0 & \mathbf{Q} \end{Vmatrix} \times \begin{Vmatrix} \bar{\mathbf{v}}_2 \\ \bar{\mathbf{i}}_2 \end{Vmatrix}. \tag{22}$$

Although the matrix on the extreme left of (22) is singular, the equation may be solved for $\{\mathbf{v}_1, \mathbf{i}_1\}$. From (18) and (20), it is easy to check that

$$\mathbf{P}'\mathbf{Q} = \mathbf{Q}'\mathbf{P} = \mathbf{I}, \tag{23}$$

a 3×3 unit matrix; then

$$\begin{Vmatrix} \mathbf{Q}' & 0 \\ 0 & \mathbf{P}' \end{Vmatrix} \times \begin{Vmatrix} \mathbf{P} & 0 \\ 0 & \mathbf{Q} \end{Vmatrix} = \begin{Vmatrix} \mathbf{Q}'\mathbf{P} & 0 \\ 0 & \mathbf{P}'\mathbf{Q} \end{Vmatrix} = \mathbf{I}, \tag{24}$$

a 6×6 unit matrix. (22) is therefore premultiplied by

$$\begin{Vmatrix} \mathbf{Q}' & 0 \\ 0 & \mathbf{P}' \end{Vmatrix},$$

yielding

$$\begin{Vmatrix} \bar{\mathbf{v}}_1 \\ \bar{\mathbf{i}}_1 \end{Vmatrix} = \begin{Vmatrix} \mathbf{Q}' & 0 \\ 0 & \mathbf{P}' \end{Vmatrix} \times \mathbf{K} \times \begin{Vmatrix} \mathbf{P} & 0 \\ 0 & \mathbf{Q} \end{Vmatrix} \times \begin{Vmatrix} \bar{\mathbf{v}}_2 \\ \bar{\mathbf{i}}_2 \end{Vmatrix}. \tag{25}$$

The required transfer matrix is, then,

$$\bar{\mathbf{K}} = \begin{Vmatrix} \mathbf{Q}' & 0 \\ 0 & \mathbf{P}' \end{Vmatrix} \times \mathbf{K} \times \begin{Vmatrix} \mathbf{P} & 0 \\ 0 & \mathbf{Q} \end{Vmatrix}. \tag{26}$$

The transformation for the submatrices of K is

$$\bar{\mathbf{K}} = \begin{Vmatrix} \mathbf{Q}' & 0 \\ 0 & \mathbf{P}' \end{Vmatrix} \times \begin{Vmatrix} \mathbf{A} & \mathbf{B} \\ \mathbf{C} & \mathbf{D} \end{Vmatrix} \times \begin{Vmatrix} \mathbf{P} & 0 \\ 0 & \mathbf{Q} \end{Vmatrix}$$

[16] J. Shekel, "Voltage reference node—its transformations in nodal analysis," *Wireless Eng.*, vol. 31, p. 6; January, 1954. See also J. Shekel, "Matrix representation of transistor circuits," Proc. I.R.E., vol. 40, p. 1493; November, 1952.

$$= \left\| \begin{matrix} \mathbf{Q'AP} & \mathbf{Q'BQ} \\ \mathbf{P'CP} & \mathbf{P'DQ} \end{matrix} \right\|, \tag{27}$$

$$\left.\begin{aligned} \overline{\mathbf{A}} &= \mathbf{Q'AP} \quad \overline{\mathbf{B}} = \mathbf{Q'BQ} \\ \overline{\mathbf{C}} &= \mathbf{P'CP} \quad \overline{\mathbf{D}} = \mathbf{P'DQ} \end{aligned}\right\}. \tag{28}$$

When $\mathbf{P}$ and $\mathbf{Q}$ are written in full, and the multiplication carried out, the final result, written for a general n, is

$$\begin{aligned} \overline{A}^{rs} &= A^{rs} - A^{ns} \\ \overline{B}^{rs} &= B^{rs} - B^{ns} - B^{rn} + B^{nn} \\ \overline{C}^{rs} &= C^{rs} \\ \overline{D}^{rs} &= D^{rs} - D^{rn} \quad (r, s = 1, 2, \cdots, n-1). \end{aligned} \tag{29}$$

These, then, are the components of the *definite* transfer matrix $\overline{\mathbf{K}}$, as obtained from those of the indefinite matrix, due to the specification of the n-th input and output terminals as voltage reference points. The numbering of the terminals, of course, is arbitrary, so that any terminal pair (input and output) may be chosen as "ground," and the transfer matrix transformed by (29).

Admittance and Impedance Matrices[17]

An alternative representation of a four-pole, other than by its transfer matrix, is by its admittance or impedance matrix. We will apply this method to the m.t.t., and define its admittance by the matrix $\mathbf{Y}$ in

$$\{\mathbf{i}_1, -\mathbf{i}_2\} = \mathbf{Y}\{\mathbf{v}_1, \mathbf{v}_2\}, \tag{30}$$

and its impedance by $\mathbf{Z}(=\mathbf{Y}^{-1})$ in

$$\{\mathbf{v}_1, \mathbf{v}_2\} = \mathbf{Z}\{\mathbf{i}_1, -\mathbf{i}_2\}, \tag{31}$$

$\mathbf{v}_1$, $\mathbf{v}_2$, $\mathbf{i}_1$ and $\mathbf{i}_2$ having the same meaning as in (1).

It should be noted that in this representation all the currents are treated as entering the terminals, hence the negative sign prefixing $\mathbf{i}_2$.

To derive $\mathbf{Y}$ from $\mathbf{K}$, (1) is written as two equations

$$\begin{aligned} \mathbf{v}_1 &= \mathbf{Av}_2 + \mathbf{Bi}_2 \\ \mathbf{i}_1 &= \mathbf{Cv}_2 + \mathbf{Di}_2 \\ \mathbf{Bi}_2 &= \mathbf{v}_1 - \mathbf{Av}_2 \\ \mathbf{i}_2 &= \mathbf{B}^{-1}\mathbf{v}_1 - \mathbf{B}^{-1}\mathbf{Av}_2. \end{aligned} \tag{32}$$

$$\begin{aligned} \mathbf{i}_1 &= \mathbf{Cv}_2 + \mathbf{D}(\mathbf{B}^{-1}\mathbf{v}_1 - \mathbf{B}^{-1}\mathbf{Av}_2) \\ &= \mathbf{Cv}_2 + \mathbf{DB}^{-1}\mathbf{v}_1 - \mathbf{DB}^{-1}\mathbf{Av}_2 \\ &= \mathbf{DB}^{-1}\mathbf{v}_1 + (\mathbf{C} - \mathbf{DB}^{-1}\mathbf{A})\mathbf{v}_2. \end{aligned} \tag{33}$$

Comparing (32) and (33) with (30),

$$\mathbf{Y} = \left\| \begin{matrix} \mathbf{DB}^{-1} & \mathbf{C} - \mathbf{DB}^{-1}\mathbf{A} \\ -\mathbf{B}^{-1} & \mathbf{B}^{-1}\mathbf{A} \end{matrix} \right\|. \tag{34}$$

By a similar derivation,

$$\mathbf{Z} = \left\| \begin{matrix} \mathbf{AC}^{-1} & \mathbf{AC}^{-1}\mathbf{D} - \mathbf{B} \\ \mathbf{C}^{-1} & \mathbf{C}^{-1}\mathbf{D} \end{matrix} \right\|. \tag{35}$$

[17] The results of this section and the following one are brought by Rice, in an appendix to his paper, "Steady state solutions of transmission line equations," *ibid.*

The inverse of $\mathbf{C}$ appears in all the submatrices of $\mathbf{Z}$. In the indefinite form of $\mathbf{K}$, $\mathbf{C}$ is singular (see (6) and (11)); but when $\mathbf{K}$ is expressed in definite form, one row and one column of $\mathbf{C}$ are removed, so that it is no longer singular. It is therefore possible to use the impedance representation only in the definite form. On the other hand, it is the inverse of $\mathbf{B}$ that appears in $\mathbf{Y}$, and as there are no restrictions on $\mathbf{B}$, the admittance representation may be used in the indefinite form as well as in the definite one.

The inverse transformation, from $\mathbf{Y}$ or $\mathbf{Z}$ to $\mathbf{K}$, is similarly derived:

$$\begin{aligned} \mathbf{i}_1 &= \mathbf{Y}_{11}\mathbf{v}_1 + \mathbf{Y}_{12}\mathbf{v}_2 \\ -\mathbf{i}_2 &= \mathbf{Y}_{21}\mathbf{v}_1 + \mathbf{Y}_{22}\mathbf{v}_2 \\ \mathbf{Y}_{21}\mathbf{v}_1 &= -\mathbf{Y}_{22}\mathbf{v}_2 - \mathbf{i}_2 \\ \mathbf{v}_1 &= -\mathbf{Y}_{21}^{-1}\mathbf{Y}_{22}\mathbf{v}_2 - \mathbf{Y}_{21}^{-1}\mathbf{i}_2, \end{aligned} \tag{36}$$

$$\begin{aligned} \mathbf{i}_1 &= \mathbf{Y}_{11}(-\mathbf{Y}_{21}^{-1}\mathbf{Y}_{22}\mathbf{v}_2 - \mathbf{Y}_{21}^{-1}\mathbf{i}_2) + \mathbf{Y}_{12}\mathbf{v}_2 \\ \mathbf{i}_1 &= (\mathbf{Y}_{12} - \mathbf{Y}_{11}\mathbf{Y}_{21}^{-1}\mathbf{Y}_{22})\mathbf{v}_2 - \mathbf{Y}_{11}\mathbf{Y}_{21}^{-1}\mathbf{i}_2. \end{aligned} \tag{37}$$

Comparing (36) and (37) with (1),

$$\mathbf{K} = \left\| \begin{matrix} -\mathbf{Y}_{21}^{-1}\mathbf{Y}_{22} & -\mathbf{Y}_{21}^{-1} \\ \mathbf{Y}_{12} - \mathbf{Y}_{11}\mathbf{Y}_{21}^{-1}\mathbf{Y}_{22} & -\mathbf{Y}_{11}\mathbf{Y}_{21}^{-1} \end{matrix} \right\|. \tag{38}$$

A similar derivation, starting from $\mathbf{Z}$, results in

$$\mathbf{K} = \left\| \begin{matrix} \mathbf{Z}_{11}\mathbf{Z}_{21}^{-1} & \mathbf{Z}_{11}\mathbf{Z}_{21}^{-1}\mathbf{Z}_{22} - \mathbf{Z}_{12} \\ \mathbf{Z}_{21}^{-1} & \mathbf{Z}_{21}^{-1}\mathbf{Z}_{22} \end{matrix} \right\|. \tag{39}$$

Transformations (34), (35), (38), and (39) are well known in four-pole theory; only there the submatrices are scalars, so that the expressions are written somewhat differently. For example, in (39),

$$\mathbf{B} = \mathbf{Z}_{11}\mathbf{Z}_{21}^{-1}\mathbf{Z}_{22} - \mathbf{Z}_{12},$$

whereas, in a four-pole, it is usually written as

$$B = \frac{Z_{11}Z_{22} - Z_{12}Z_{21}}{Z_{21}} = \frac{|\mathbf{Z}|}{Z_{21}}.$$

Reciprocity

Nothing has been assumed, so far, about reciprocity relations in the network. All that was said in the preceding sections, and will be said in the sections following, applies to networks which may contain vacuum tubes, transistors or gyrators, in addition to bilateral elements, with the only restriction that all elements operate in their linear ranges. In this section we will discuss a special type of the general linear network, viz. the class of networks that exhibit the reciprocity relation. It is to be expected that this restriction on the generality of the network will be reflected in its transfer matrix.

The reciprocity relation is expressed by the symmetry of the $\mathbf{Y}$ and $\mathbf{Z}$ matrices. From the symmetry of $\mathbf{Y}$,

$$\mathbf{Y}_{11} = \mathbf{Y}_{11}' \tag{40}$$

$$\mathbf{Y}_{22} = \mathbf{Y}_{22}' \tag{41}$$

$$\mathbf{Y}_{12} = \mathbf{Y}_{21}'. \tag{42}$$

From (34) and (40),

$$\mathbf{DB}^{-1} = (\mathbf{DB}^{-1})' = (\mathbf{B}^{-1})'\mathbf{D}'$$
$$\mathbf{B}'\mathbf{DB}^{-1} = \mathbf{D}'$$
$$\mathbf{B}'\mathbf{D} = \mathbf{D}'\mathbf{B}, \tag{43}$$

while (34) and (41) yield

$$\mathbf{B}^{-1}\mathbf{A} = (\mathbf{B}^{-1}\mathbf{A})' = \mathbf{A}'(\mathbf{B}^{-1})'$$
$$\mathbf{AB}' = \mathbf{BA}'. \tag{44}$$

(42) is equivalent to

$$-(\mathbf{B}')^{-1} = \mathbf{C} - \mathbf{DB}^{-1}\mathbf{A}$$
$$\mathbf{C} = \mathbf{DB}^{-1}\mathbf{A} - (\mathbf{B}')^{-1}$$
$$= (\mathbf{B}')^{-1}\mathbf{D}'\mathbf{A} - (\mathbf{B}')^{-1}$$
$$= (\mathbf{B}')^{-1}(\mathbf{D}'\mathbf{A} - \mathbf{I})$$
$$\mathbf{B}'\mathbf{C} = \mathbf{D}'\mathbf{A} - \mathbf{I}$$
$$\mathbf{D}'\mathbf{A} - \mathbf{B}'\mathbf{C} = \mathbf{I}, \tag{45}$$

or, proceeding differently,

$$\mathbf{C} = \mathbf{DB}^{-1}\mathbf{A} - (\mathbf{B}')^{-1}$$
$$= \mathbf{DA}'(\mathbf{B}')^{-1} - (\mathbf{B}')^{-1}$$
$$= (\mathbf{DA}' - \mathbf{I})(\mathbf{B}')^{-1}$$
$$\mathbf{CB}' = \mathbf{DA}' - \mathbf{I}$$
$$\mathbf{DA}' - \mathbf{CB}' = \mathbf{I}. \tag{46}$$

A similar derivation, starting with the symmetry of $\mathbf{Z}$, gives four analogous results:

$$\mathbf{C}'\mathbf{A} = \mathbf{A}'\mathbf{C} \tag{47}$$
$$\mathbf{DC}' = \mathbf{CD}' \tag{48}$$
$$\mathbf{A}'\mathbf{D} - \mathbf{C}'\mathbf{B} = \mathbf{I} \tag{49}$$
$$\mathbf{AD}' - \mathbf{BC}' = \mathbf{I}. \tag{50}$$

The eight restrictions, (43) to (50), are equivalent to

$$\begin{Vmatrix} \mathbf{A} & \mathbf{B} \\ \mathbf{C} & \mathbf{D} \end{Vmatrix}^{-1} = \begin{Vmatrix} \mathbf{D}' & -\mathbf{B}' \\ -\mathbf{C}' & \mathbf{A}' \end{Vmatrix}, \tag{51}$$

as may easily be confirmed by direct multiplication of both matrices.

In four-poles, where A, B, C and D are scalars, (43), (44), (47) and (48) are no restrictions, and the other four results merge into one, which is the only consequence of the reciprocity relation:

$$|\mathbf{K}| = AD - BC = 1.$$

Symmetry

A multi-terminal transducer is symmetrical if its transfer matrix is invariant to the interchange of roles between input and output terminals.

The interchange implies reversal of signs of both $\mathbf{i}_1$ and $\mathbf{i}_2$, to conform with the accepted definitions, so that after the interchange

$$\begin{Vmatrix} \mathbf{v}_2 \\ -\mathbf{i}_2 \end{Vmatrix} = \begin{Vmatrix} \mathbf{A} & \mathbf{B} \\ \mathbf{C} & \mathbf{D} \end{Vmatrix} \times \begin{Vmatrix} \mathbf{v}_1 \\ -\mathbf{i}_1 \end{Vmatrix}, \tag{52}$$

or,

$$\begin{Vmatrix} \mathbf{v}_2 \\ \mathbf{i}_2 \end{Vmatrix} = \begin{Vmatrix} \mathbf{A} & -\mathbf{B} \\ -\mathbf{C} & \mathbf{D} \end{Vmatrix} \times \begin{Vmatrix} \mathbf{v}_1 \\ \mathbf{i}_1 \end{Vmatrix}. \tag{53}$$

Comparing (1), (2) and (53),

$$\begin{Vmatrix} \mathbf{A} & \mathbf{B} \\ \mathbf{C} & \mathbf{D} \end{Vmatrix}^{-1} = \begin{Vmatrix} \mathbf{A} & -\mathbf{B} \\ -\mathbf{C} & \mathbf{D} \end{Vmatrix}. \tag{54}$$

By direct multiplication, (54) may be expressed by four relations,

$$\mathbf{A}^2 - \mathbf{BC} = \mathbf{I} \tag{55}$$
$$\mathbf{D}^2 - \mathbf{CB} = \mathbf{I} \tag{56}$$
$$\mathbf{AB} = \mathbf{BD} \tag{57}$$
$$\mathbf{DC} = \mathbf{CA}. \tag{58}$$

If the transducer exhibits both symmetry and reciprocity, (54) and (51) are to be satisfied simultaneously; then

$$\mathbf{A} = \mathbf{D}'; \quad \mathbf{B} = \mathbf{B}'; \quad \mathbf{C} = \mathbf{C}'. \tag{59}$$

In four-poles, (57) and (58) are equivalent to $A = D$; while this and (55) or (56) show that

$$|\mathbf{K}| = AD - BC = A^2 - BC = D^2 - BC = 1,$$

and (59) adds nothing new.

Indeed, in a four-pole, symmetry ($Y_{11} = Y_{22}$, $Y_{12} = Y_{21}$) implies reciprocity ($Y_{12} = Y_{21}$). On the other hand, a m.t.t. may be symmetrical ($\mathbf{Y}_{12} = \mathbf{Y}_{21}$) without exhibiting reciprocity ($\mathbf{Y}_{12} = \mathbf{Y}_{21}'$); see, for example, the network in Fig. 2, which is symmetrical, but does not obey the reciprocity relation, due to the presence of the vacuum tube.

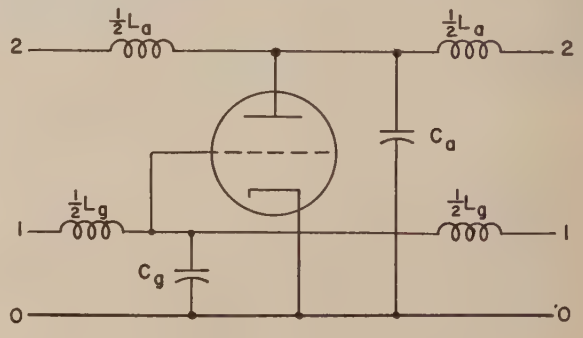

Fig. 2—One section of a distributed amplifier.

Admittance and Impedance Transformations

An n-terminal transducer may be terminated by an n-terminal network. The termination is represented by a square matrix $\mathbf{Y}_2$ or $\mathbf{Z}_2$,

$$\mathbf{i}_2 = \mathbf{Y}_2\mathbf{v}_2, \tag{60}$$
$$\mathbf{v}_2 = \mathbf{Z}_2\mathbf{i}_2, \tag{61}$$

and the *admittance* representation may be definite or indefinite.[16] $\mathbf{Y}_2$ is the admittance matrix obtained from a nodal analysis of the network, $\mathbf{Z}_2$ is the inverse of $\mathbf{Y}_2$, when the latter is expressed in its definite form. (Neither

$\mathbf{Y}_2$ nor $\mathbf{Z}_2$ is connected with the mesh analysis of the terminating network.)

The terminated transducer appears at its input terminals as an n-terminal network of admittance $\mathbf{Y}_1$ or impedance $\mathbf{Z}_1$. The transducer thus *transforms* $\mathbf{Y}_2$ into $\mathbf{Y}_1$ or $\mathbf{Z}_2$ into $\mathbf{Z}_1$.

The transfer equation (1) may be written as

$$\mathbf{v}_1 = \mathbf{A}\mathbf{v}_2 + \mathbf{B}(\mathbf{Y}_2\mathbf{v}_2)$$
$$(\mathbf{Y}_1\mathbf{v}_1) = \mathbf{C}\mathbf{v}_2 + \mathbf{D}(\mathbf{Y}_2\mathbf{v}_2),$$

whence

$$\mathbf{Y}_1(\mathbf{A} + \mathbf{B}\mathbf{Y}_2)\mathbf{v}_2 = (\mathbf{C} + \mathbf{D}\mathbf{Y}_2)\mathbf{v}_2;$$

this is true for any $\mathbf{v}_2$, so that

$$\mathbf{Y}_1 = (\mathbf{C} + \mathbf{D}\mathbf{Y}_2)(\mathbf{A} + \mathbf{B}\mathbf{Y}_2)^{-1}. \tag{62}$$

The impedance transformation is obtained if we write (1) as

$$(\mathbf{Z}_1\mathbf{i}_1) = \mathbf{A}(\mathbf{Z}_2\mathbf{i}_2) + \mathbf{B}\mathbf{i}_2$$
$$\mathbf{i}_1 = \mathbf{C}(\mathbf{Z}_2\mathbf{i}_2) + \mathbf{D}\mathbf{i}_2,$$

then

$$\mathbf{Z}_1(\mathbf{C}\mathbf{Z}_2 + \mathbf{D})\mathbf{i}_2 = (\mathbf{A}\mathbf{Z}_2 + \mathbf{B})\mathbf{i}_2$$
$$\mathbf{Z}_1 = (\mathbf{A}\mathbf{Z}_2 + \mathbf{B})(\mathbf{C}\mathbf{Z}_2 + \mathbf{D})^{-1}. \tag{63}$$

Transformations (62) and (63) are well known in four-pole theory, only there they are usually written as quotients of scalar expressions.

Derivation of the Transfer Matrix of a Given Transducer

Cascading

Two m.t.t.'s may be connected *in cascade*, so that the output terminals of one are the input terminals of the other. Let $\mathbf{K}_1$ and $\mathbf{K}_2$ be the transfer matrices of the transducers, respectively. The transfer from the first set of terminals to the second one is

$$\{\mathbf{v}_1, \mathbf{i}_1\} = \mathbf{K}_1\{\mathbf{v}_2, \mathbf{i}_2\}$$

and from the second to the third

$$\{\mathbf{v}_2, \mathbf{i}_2\} = \mathbf{K}_2\{\mathbf{v}_3, \mathbf{i}_3\}.$$

The total transfer from the first to the third is then

$$\{\mathbf{v}_1, \mathbf{i}_1\} = \mathbf{K}_1\mathbf{K}_2\{\mathbf{v}_3, \mathbf{i}_3\},$$

so that the two cascaded transducers are equivalent to a single m.t.t. whose transfer matrix is

$$\mathbf{K} = \mathbf{K}_1\mathbf{K}_2. \tag{64}$$

Conversely, a complicated m.t.t. may sometimes be split into a cascade of simpler sections, and its transfer matrix derived as a product of simple transfer matrices.

We will now proceed to derive the transfer matrices of some simple sections, which may be used by themselves or as building blocks for more complicated m.t.t.'s.

Series section

In the m.t.t. shown in Fig. 3 it is evident that

$$\mathbf{i}_1 = \mathbf{i}_2, \tag{65}$$

but

$$\mathbf{v}_1 = \mathbf{v}_2 + \mathbf{Z}\mathbf{i}_2. \tag{66}$$

$\mathbf{Z}$ is a square matrix, where Z^{pp} is the impedance of the element between the $p-p$ terminals, and Z^{pq} is the

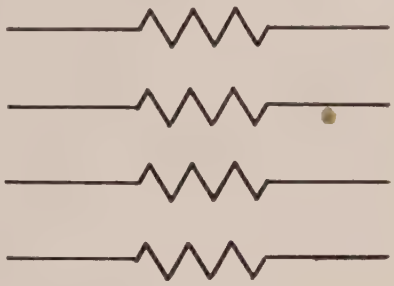

Fig. 3—Series section.

mutual impedance between the elements in the $p-p$ and $q-q$ branches. Comparing (65) and (66) with (1), we see that for the series section

$$\mathbf{K} = \left\| \begin{matrix} \mathbf{I} & \mathbf{Z} \\ 0 & \mathbf{I} \end{matrix} \right\|. \tag{67}$$

This $\mathbf{K}$ is in its indefinite form.

Shunt section

The m.t.t. in Fig. 4 has

$$\mathbf{v}_1 = \mathbf{v}_2, \tag{68}$$

and

$$\mathbf{i}_1 = \mathbf{i}_2 + \mathbf{Y}\mathbf{v}_2. \tag{69}$$

Here $\mathbf{Y}$ is the indefinite admittance matrix of the shunting network. If this network comprises 2-terminal elements only, Y^{pp} is the sum of the admittances of the

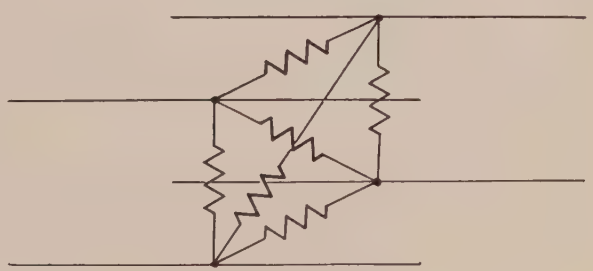

Fig. 4—Shunt section.

elements connected to the p terminals, and Y^{pq} is the negative of the admittance of the element in the branch between the p and q terminals.

(68) and (69) together give

$$\mathbf{K} = \left\| \begin{matrix} \mathbf{I} & 0 \\ \mathbf{Y} & \mathbf{I} \end{matrix} \right\| \tag{70}$$

as the indefinite transfer matrix of the shunt section.

"Twist" section

Fig. 5 shows a m.t.t. whose only effect is to introduce a "twist" between the input and output terminals. The components of $\mathbf{v}_2$ are merely a permutation of those of $\mathbf{v}_1$, and $\mathbf{i}_2$ is similarly derived from $\mathbf{i}_1$ by the same permutation.

$$\begin{aligned} \mathbf{v}_1 &= \mathbf{T}\mathbf{v}_2 \\ \mathbf{i}_1 &= \mathbf{T}\mathbf{i}_2 \end{aligned} \tag{71}$$

Fig. 5—Twist section.

$\mathbf{T}$ is a square matrix, constructed as follows: for each connection between the p-th input terminal and q-th output terminal, $T^{pq}=1$; all other elements of $\mathbf{T}$ are zero. For the network in Fig. 5,

$$\mathbf{T} = \begin{Vmatrix} 0 & 1 & 0 & 0 \\ 0 & 0 & 0 & 1 \\ 0 & 0 & 1 & 0 \\ 1 & 0 & 0 & 0 \end{Vmatrix}.$$

(71) indicates that the indefinite transfer matrix of a twist section is

$$\mathbf{K} = \begin{Vmatrix} \mathbf{T} & 0 \\ 0 & \mathbf{T} \end{Vmatrix}. \tag{72}$$

More complicated sections

Some sections, like lattices or bridged sections, cannot be treated as cascades of simpler sections. Fig. 6 may serve as an example, which, but for the element b, could be treated as a cascade of a series section and a twist section. A m.t.t. like this may easily be analyzed, by first writing its admittance matrix, then passing to the transfer matrix via (38).

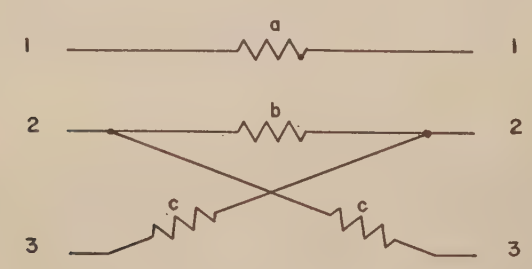

Fig. 6—Section to be analyzed by the admittance matrix.

In the example shown in Fig. 6, a, b and c denote admittance values. The indefinite admittance matrix of the network is

$$\mathbf{Y} = \left\| \begin{array}{ccc:ccc} a & 0 & 0 & -a & 0 & 0 \\ 0 & b+c & 0 & 0 & -b & -c \\ 0 & 0 & c & 0 & -c & 0 \\ \hdashline -a & 0 & 0 & a & 0 & 0 \\ 0 & -b & -c & 0 & b+c & 0 \\ 0 & -c & 0 & 0 & 0 & c \end{array} \right\|.$$

From (38) we get the four submatrices of $\mathbf{K}$:

$$|\mathbf{Y}_{21}| = \begin{vmatrix} -a & 0 & 0 \\ 0 & -b & -c \\ 0 & -c & 0 \end{vmatrix} = ac^2$$

$$\mathbf{B} = -\mathbf{Y}_{21}^{-1} = -\frac{1}{ac^2} \times \begin{Vmatrix} -c^2 & 0 & 0 \\ 0 & 0 & -ac \\ 0 & -ac & ab \end{Vmatrix}$$

$$\mathbf{B} = \begin{Vmatrix} \frac{1}{a} & 0 & 0 \\ 0 & 0 & \frac{1}{c} \\ 0 & \frac{1}{c} & -\frac{b}{c^2} \end{Vmatrix}$$

$$\mathbf{A} = -\mathbf{Y}_{21}^{-1}\mathbf{Y}_{22} = \mathbf{B}\mathbf{Y}_{22}$$

$$= \begin{Vmatrix} \frac{1}{a} & 0 & 0 \\ 0 & 0 & \frac{1}{c} \\ 0 & \frac{1}{c} & -\frac{b}{c^2} \end{Vmatrix} \times \begin{Vmatrix} a & 0 & 0 \\ 0 & b+c & 0 \\ 0 & 0 & c \end{Vmatrix}$$

$$\mathbf{A} = \begin{Vmatrix} 1 & 0 & 0 \\ 0 & 0 & 1 \\ 0 & 1+\frac{b}{c} & -\frac{b}{c} \end{Vmatrix}$$

$$\mathbf{D} = -\mathbf{Y}_{11}\mathbf{Y}_{21}^{-1} = \mathbf{Y}_{11}\mathbf{B}$$

$$= \begin{Vmatrix} a & 0 & 0 \\ 0 & b+c & 0 \\ 0 & 0 & c \end{Vmatrix} \times \begin{Vmatrix} \frac{1}{a} & 0 & 0 \\ 0 & 0 & \frac{1}{c} \\ 0 & \frac{1}{c} & -\frac{b}{c^2} \end{Vmatrix}$$

$$\mathbf{D} = \begin{Vmatrix} 1 & 0 & 0 \\ 0 & 0\,1+\frac{b}{c} & \\ 0 & 1 & -\frac{b}{c} \end{Vmatrix}.$$

$$\mathbf{C} = \mathbf{Y}_{12} - \mathbf{Y}_{11}\mathbf{Y}_{21}^{-1}\mathbf{Y}_{22} = \mathbf{Y}_{12} + \mathbf{D}\mathbf{Y}_{22}$$

$$= \left\| \begin{matrix} -a & 0 & 0 \\ 0 & -b & -c \\ 0 & -c & 0 \end{matrix} \right\| + \left\| \begin{matrix} 1 & 0 & 0 \\ 0 & 0 & 1+\frac{b}{c} \\ 0 & 1 & -\frac{b}{c} \end{matrix} \right\|$$

$$\times \left\| \begin{matrix} a & 0 & 0 \\ 0 & b+c & 0 \\ 0 & 0 & c \end{matrix} \right\|$$

$$= \left\| \begin{matrix} -a & 0 & 0 \\ 0 & -b & -c \\ 0 & -c & 0 \end{matrix} \right\| + \left\| \begin{matrix} a & 0 & 0 \\ 0 & 0 & b+c \\ 0 & b+c & -b \end{matrix} \right\|$$

$$\mathbf{C} = \left\| \begin{matrix} 0 & 0 & 0 \\ 0 & -b & b \\ 0 & b & -b \end{matrix} \right\|.$$

In the distributed amplifier, the series sections are composed of the inductances

$$\mathbf{Z}_1 = \mathbf{Z}_3 = \left\| \begin{matrix} j\omega L_g & 0 \\ 0 & j\omega L_a \end{matrix} \right\|.$$

The shunt section contains the capacitors and the vacuum tube. We shall denote the mutual conductance and the internal conductance of the tube by G_m and G_a, respectively. The electrode-to-cathode capacitances may be assimilated in C_g and C_a.

$$\mathbf{Y}_2 = \left\| \begin{matrix} 0 & 0 \\ G_m & G_a \end{matrix} \right\| + \left\| \begin{matrix} j\omega C_g & 0 \\ 0 & j\omega C_a \end{matrix} \right\|$$

$$= \left\| \begin{matrix} j\omega C_g & 0 \\ G_m & G_a + j\omega C_a \end{matrix} \right\|.$$

Substituting these values in (73) we get the complete transfer matrix

$$\mathbf{K} = \left\| \begin{matrix} 1-\frac{1}{2}\omega^2 L_g C_g & 0 & j\omega L_g(1-\frac{1}{4}\omega^2 L_g C_g) & 0 \\ \frac{1}{2}j\omega L_a G_m & 1-\frac{1}{2}\omega^2 L_a C_a+\frac{1}{2}j\omega L_a G_a & -\frac{1}{4}\omega^2 L_a L_g G_m & j\omega L_a(1-\frac{1}{4}\omega^2 L_a C_a)-\frac{1}{4}\omega^2 L_a^2 G_a \\ j\omega C_g & 0 & 1-\frac{1}{2}\omega^2 L_g C_g & 0 \\ G_m & G_a+j\omega C_a & \frac{1}{2}j\omega L_g G_m & 1-\frac{1}{2}\omega^2 L_a C_a+\frac{1}{2}j\omega L_a G_a \end{matrix} \right\|.$$

This example illustrates some of the general results which have been derived:

1. The transfer matrix is indefinite. The sum of each row and column of **C** is zero, the sum of each column of **D** is 1, and the sum of each row of **A** is 1.

2. The network is symmetrical, and obeys the reciprocity relation. **B** and **C** are symmetrical matrices, and **A** is the tranpose of **D**, as in (59). All the other relations, viz. (43) to (50) for reciprocity, and (55) to (58) for symmetry, may also be checked by a simple computation.

Distributed Amplifer

As a concluding illustrative example, we will derive the transfer matrix for the section of a distributed amplifier shown in Fig. 2. As there is an obvious ground line running through the section, the definite representation will be used.

The section may be treated as a cascade of three sections, of the series-shunt-series type. (This is a generalization of the T-section in four-pole theory.) The complete transfer matrix, in the general case, is

$$\mathbf{K} = \mathbf{K}_1\mathbf{K}_2\mathbf{K}_3 = \left\| \begin{matrix} \mathbf{I} & \mathbf{Z}_1 \\ 0 & \mathbf{I} \end{matrix} \right\| \times \left\| \begin{matrix} \mathbf{I} & 0 \\ \mathbf{Y}_2 & \mathbf{I} \end{matrix} \right\| \times \left\| \begin{matrix} \mathbf{I} & \mathbf{Z}_3 \\ 0 & \mathbf{I} \end{matrix} \right\|$$

$$= \left\| \begin{matrix} \mathbf{I}+\mathbf{Z}_1\mathbf{Z}_2 & \mathbf{Z}_1+\mathbf{Z}_1\mathbf{Y}_2\mathbf{Z}_3+\mathbf{Z}_3 \\ \mathbf{Y}_2 & \mathbf{Y}_2\mathbf{Z}_3+\mathbf{I} \end{matrix} \right\|. \quad (73)$$

This network is symmetrical, but does not obey reciprocity relation, due to presence of the vacuum tube.

Conclusion

The multi-terminal transducer is a type of network that includes the four-pole as a special case. All the results derived for the general m.t.t. apply to four-poles, although some of them appear trivial when applied to this special case. In four-poles, the various submatrices are scalars, and then they obey the commutative law in multiplication; scalars are also equal to their own transpose, and these two properties deprive relations like (47), (57), or (59) of any special meaning.

In our treatment, we did or did not specify the voltage reference points, as was more convenient in the case at hand. The use of the indefinite transfer matrix thus gives more freedom in representation or computation. This method is also applicable to four-poles, where it enables to distinguish between "T"-sections (with the series impedance in one lead only) and "H"-sections (with the series impedance distributed, equally or unequally, between both leads), or similar cases.

This paper treated linear m.t.t.'s in general, reciprocal or nonreciprocal, active or passive. The special properties of networks obeying the reciprocity relation were discussed. Energy considerations will be the subject matter of a forthcoming paper.

Prolonged Space-Wave Fadeouts at 1,046 MC Observed in Cheyenne Mountain Propagation Program*

BRADFORD R. BEAN†

Summary—During the first year of continuous operation of the Cheyenne Mountain propagation program, recordings of 1,046-mc fields at receiving locations within the radio horizon have exhibited "fadeouts" or prolonged periods of attenuation often in excess of 20 db below the monthly median level and lasting from a minute up to several hours. The occurrence of these fadeouts has been found to coincide with widespread super-refraction as evidenced by enhanced signals beyond the radio horizon and ground modification of the refractive index profile.

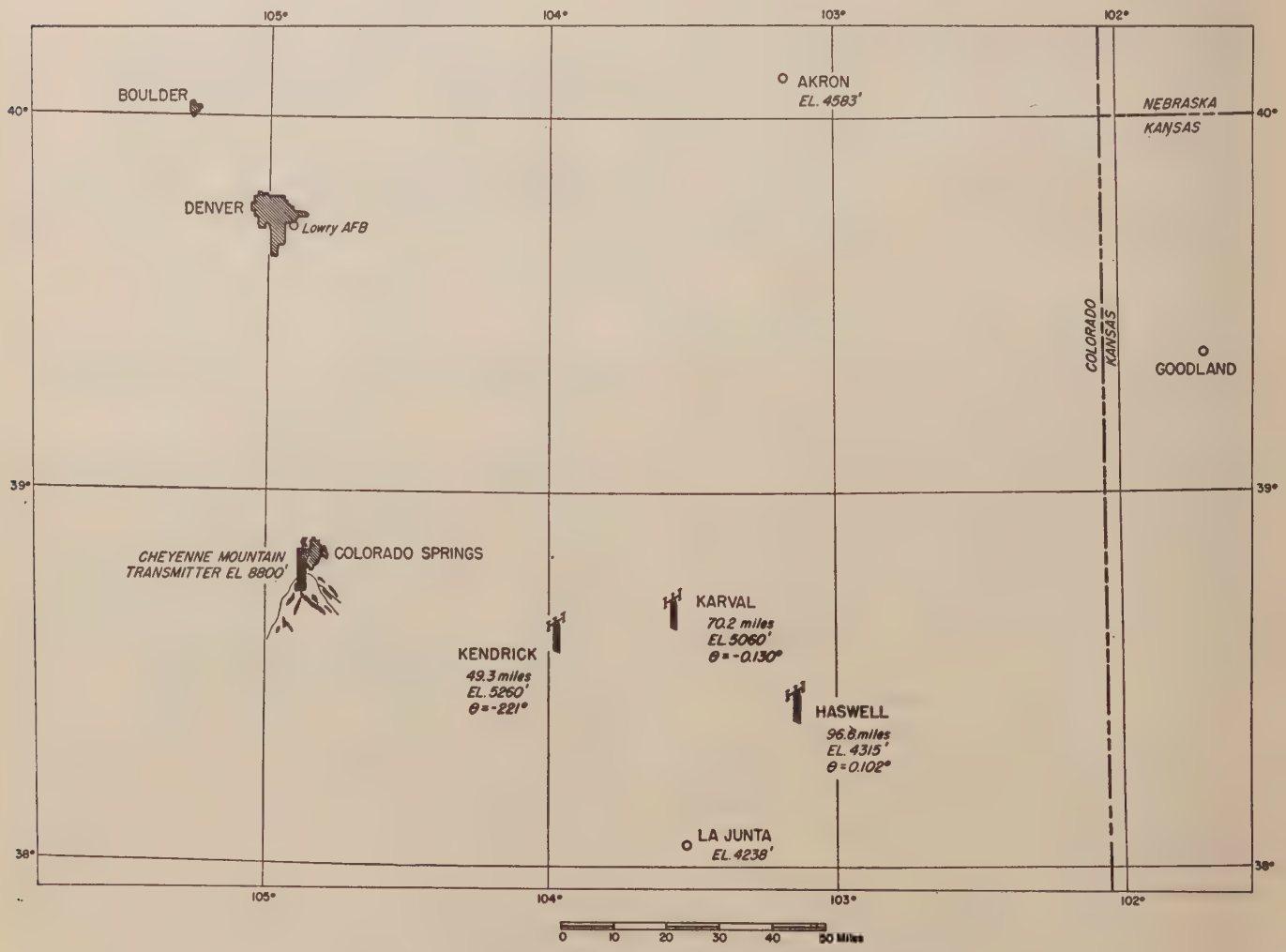

Fig. 1—Transmission paths of the Cheyenne Mountain propagation program.

THE TRANSMISSION PATHS involved in the 1,046-mc Cheyenne Mountain propagation program are shown pictorially in Fig. 1. The transmitter is located on the sheer face of Cheyenne Mountain, approximately 2,800 feet above the surrounding plains, and simulates an air-borne transmitter. The receiving sites used in this study have standard dipole antennas 43 feet above the surrounding plains and are located as follows: (*a*) well within the radio horizon at Kendrick, Colorado (49.3 miles, $\theta = -0.221$ degrees), (*b*) near the radio horizon at Karval, Colorado (70.2 miles, $\theta = -0.130$ degrees), and (*c*) beyond the radio horizon at Haswell, Colorado (96.6 miles, $\theta = 0.102$ degrees). The angle θ is the angle between the lines from the transmitting and receiving antennas to their respective radio horizons, in the great-circle plane containing the antennas for the radio standard atmosphere. This angle θ provides, in a single parameter, an indication of path length, terrain profile, and antenna heights.

Fig. 2 presents simultaneous recordings of the 100-, 192.8-, and 1,046-mc fields at Kendrick and Karval during fadeouts on the morning of September 4, 1952, and

* Decimal classification: R113.1. Original manuscript received by the IRE, September 19, 1953.

† National Bureau of Standards, Boulder, Colo.

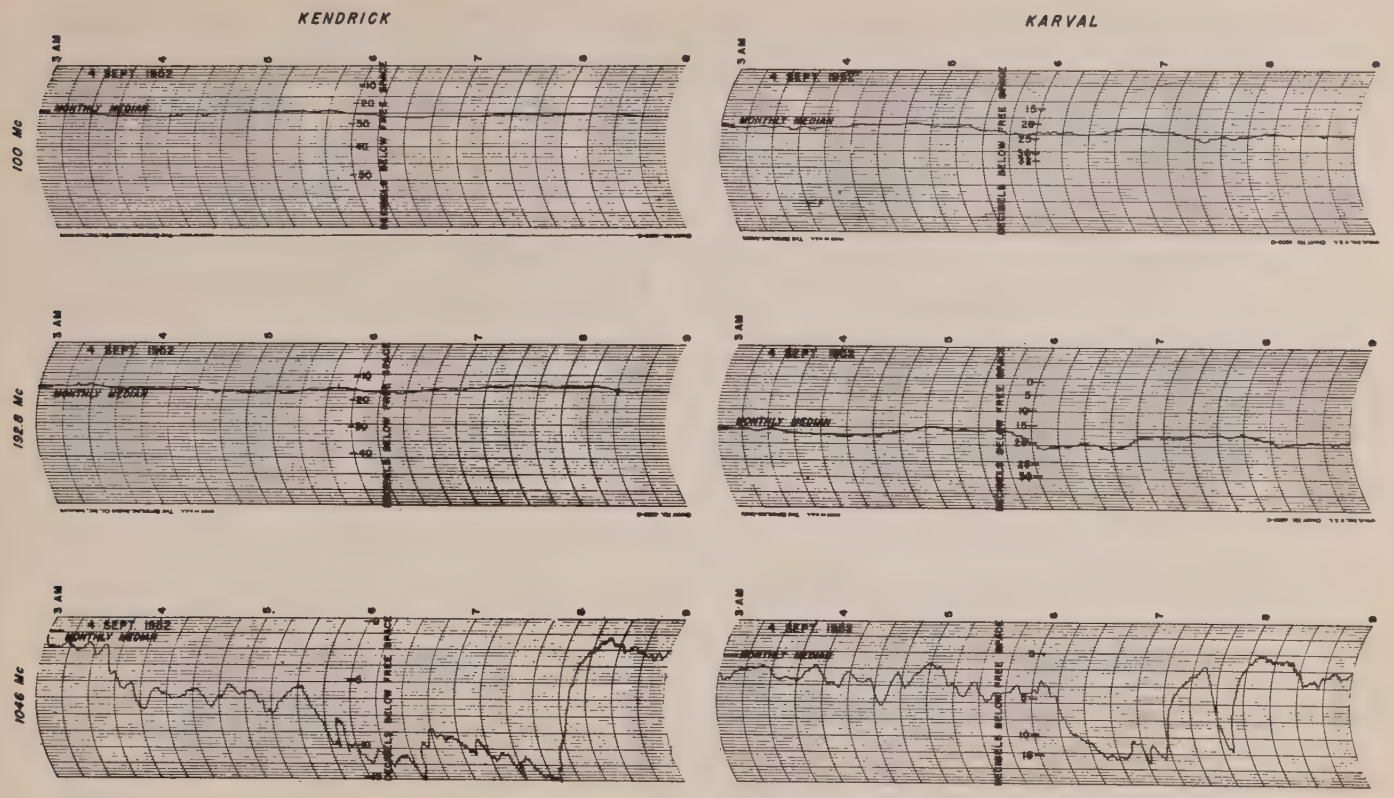

Fig. 2—Recorded 100-, 192.8-, and 1,046-mc fields for the morning of September 4, 1952, at Kendrick and Karval, Colo.

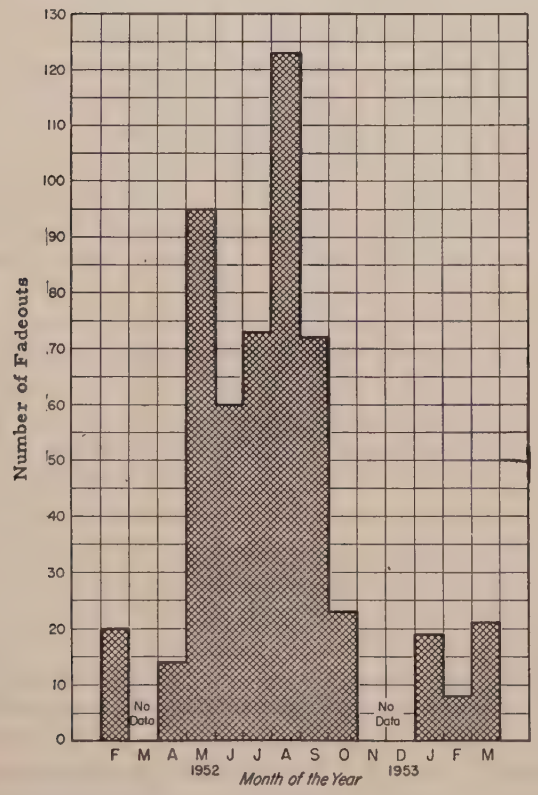

Fig. 3—Annual distribution of *normalized* occurrences of prolonged space-wave fadeouts observed at 1,046 mc at Karval, Colo.

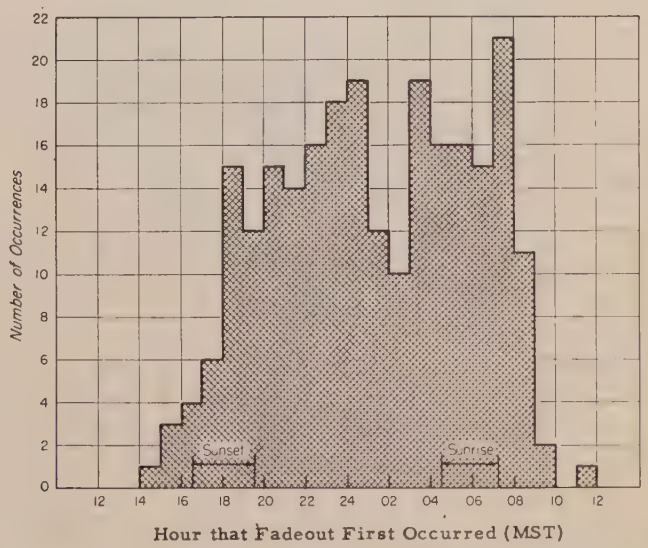

Fig. 4—Diurnal distribution of the 246 prolonged space-wave fadeouts observed at 1,046 mc at Karval, Colo., from February, 1952, to January, 1953.

illustrates the significant fact that these prolonged space-wave fadeouts are of importance only at the higher frequencies. For the purpose of this study, a prolonged space-wave fadeout is said to occur whenever the field at Karval on 1,046 mc drops 5 db or more below the monthly median level for one minute or more. The 5-db level was chosen to make the phenomenon operationally significant while the minimum of one minute rules out the rapid phase interference types of fading such as those due to passing aircraft. The normalized monthly frequency of fadeouts is shown in Fig. 3, and the diurnal distributions of the time of onset of these fadeouts is shown in Fig. 4. The data of Fig. 3 have been normal-

ized to allow for periods when the transmitter was off the air. The normalized occurrence of fadeouts was arrived at by the following relationship:

$$F_N = F_O H_M / H_O$$

where

F_N = normalized number of fadeouts,
F_O = number of fadeouts observed during the month,
H_O = number of hours of data during the month,
H_M = maximum possible hours of data during the month.

H_M/H_O had a maximum value of 5.71 in October and a minimum of 1.35 in September. Figs. 3 and 4 show a tendency for maximum occurrence during the summer months and during the nighttime hours. The fadeouts have minimum occurrence in the early afternoon, maximum occurrences a few hours after sunset and sunrise, and a secondary minimum in the early morning hours. Fig. 5 shows the percentage of total recording time that the fields were at least 5, 7, 10, and 15 db below the monthly median. May has the maximum occurrence of 5-db fadeouts with 10.5 per cent of the time, while the occurrence for any summer month is in excess of 4.9 per cent of the time; however, since the lowest occurrence, October, is 2.8 per cent, all months have occurrences which are significant for systems designed to operate out to the radio horizon.

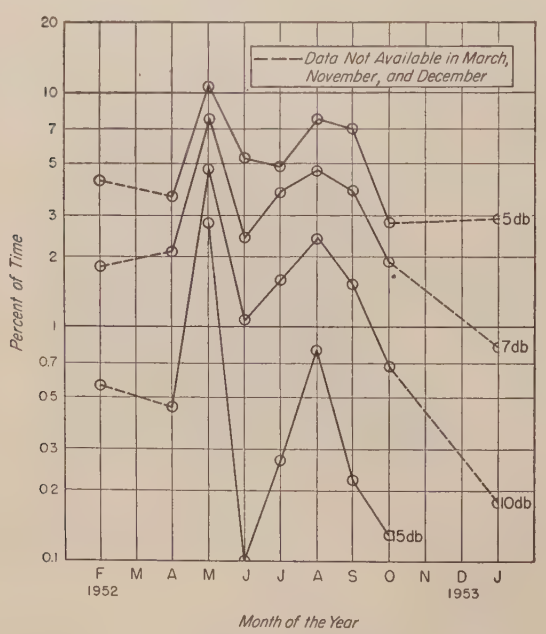

Fig. 5—Percentage of total time for each month that the 1,046-mc field at Karval, Colo., is at least 5, 7, 10, and 15 db below the monthly median field due to prolonged space-wave fadeouts.

Fig. 6 shows cumulative distributions of the durations of individual fadeouts in excess of 5, 10, and 15 db. The operational significance of the fadeouts is again emphasized by noting that 90 per cent of the fadeouts of at least 5 db are at least 6 minutes long, while 50 per cent are of at least 24 minutes duration.

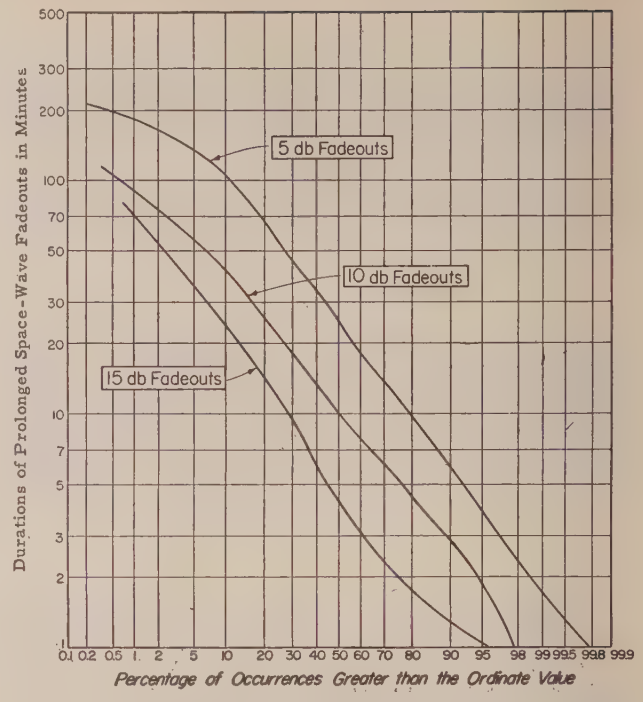

Fig. 6—Cumulative distributions of the durations of prolonged space-wave fadeouts at Karval, Colo.

The dependence of these fadeouts on meteorological conditions was presented during field trials of the NBS mobile low-level sounding system[1] at Haswell, Colorado on the night of June 21–22, 1952. The observed refractivity,[1] N, profiles are shown in Fig. 7, page 851, and show nature of the diurnal variation of the refractivity profile. There was a linear afternoon gradient of -0.0125 N units/meter between the surface and 200 meters (compared with -0.0392 N units/meter in a standard atmosphere). The surface gradient increased as the night went on, the 2320 and 0250 profiles having ducting gradients, i.e., a gradient equal to or greater than -0.157 N units/meter. After sunrise the steep surface gradients were modified by convection to produce profiles similar to those at 0554 or 0945. The 2320 and 0250 profiles are interesting in that they have ducting gradients to heights of 91 and 102 meters, respectively, while the ground-based layers display positive refractivity gradients. Fig. 7 contains reference gradients for both the radio standard atmosphere and ducting, along with their respective equivalent-earth-radius factor, k.

The recorded 1,046-mc fields at Kendrick, Karval, and Haswell during night of June 21–22, 1952, are in Fig. 8, page 852. Starting at 10 P.M., June 21, fadeouts of 5 to 15 db were observed at Kendrick and Karval, while the field beyond the radio horizon at Haswell showed a spectacular rise at 11:30 P.M. An examination of the 100- and 192.8-mc recording charts for Haswell revealed a gradual increase of field, going off scale at midnight, giving evidence of an atmospheric duct. A further con-

[1] P. D. Lowell, W. Hakkarinen, and D. L. Randall, "National Bureau of Standards mobile low-level sounding system," *Jour. Res. Nat. Bur. Stand.*, vol. 50, no. 1; 1953.

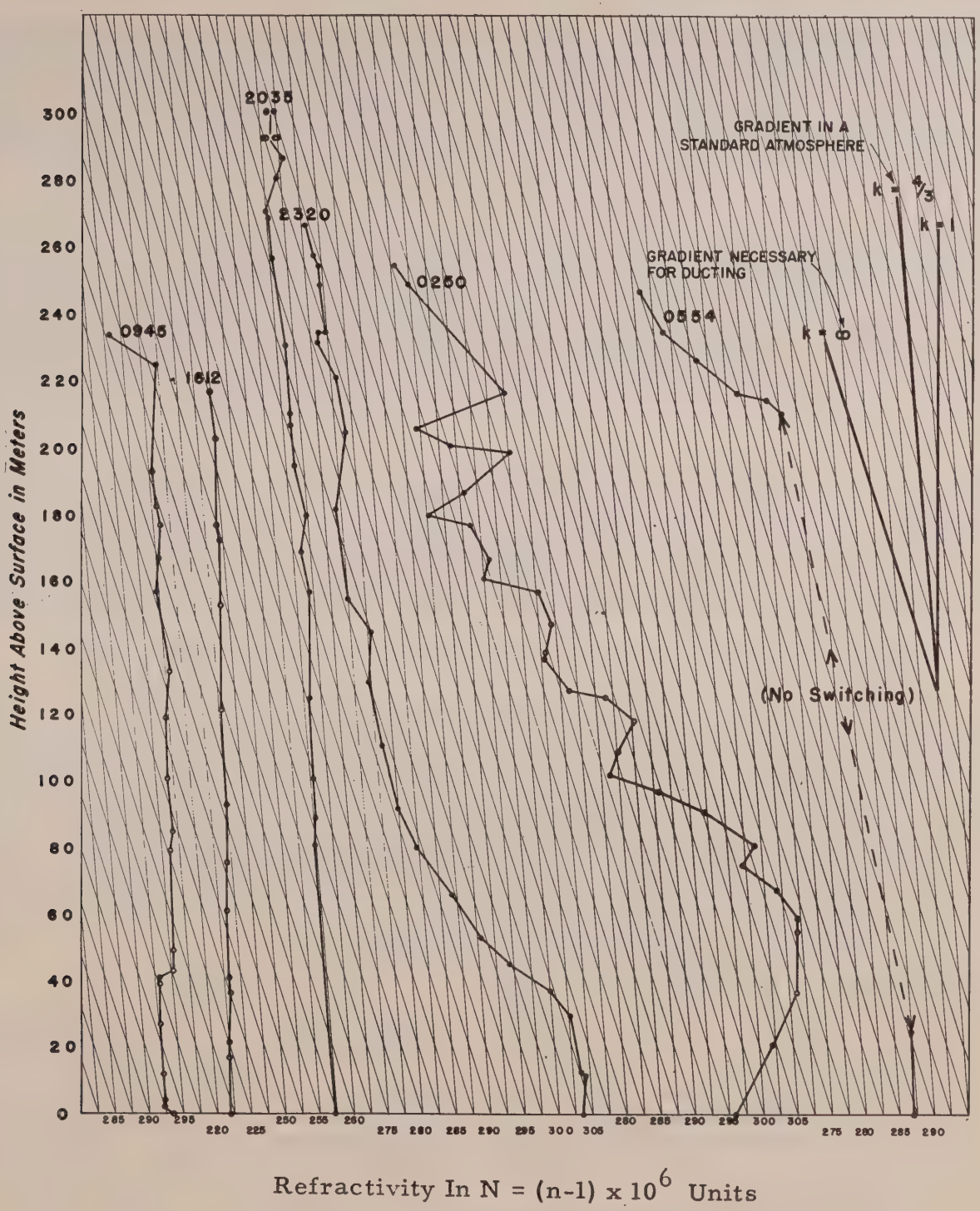

Fig. 7—Vertical distributions of refractivity at Haswell, Colo., for June 21–22, 1952.

dition to be satisfied in ducting theory[2] is that the ducting refractivity gradient extends over a sufficient height, which for 1,046 mc is about 50 meters. This latter condition is verified by the meteorological measurements described above.

Consideration of the above data sheds some light upon the possible mechanism producing the fadeouts. The most obvious feature is the presence of a duct at Haswell as borne out by both the radio and meteorological data. This conclusion is, in general, borne out by the fact that Haswell fields are above normal during 76 per cent of the time of fadeouts at Karval. For ducts to occur there must be widespread horizontal homogeneity. Therefore, it is reasonable to assume that the atmosphere near the ground in the Karval area is experiencing a degree of modification similar to that measured at Haswell and that the fadeouts at Karval occur simultaneously with increased fields at Haswell. An

[2] D. E. Kerr, "Propagation of Short Radio Waves," McGraw-Hill Book Co., Inc., New York, N. Y., p. 21; 1951.

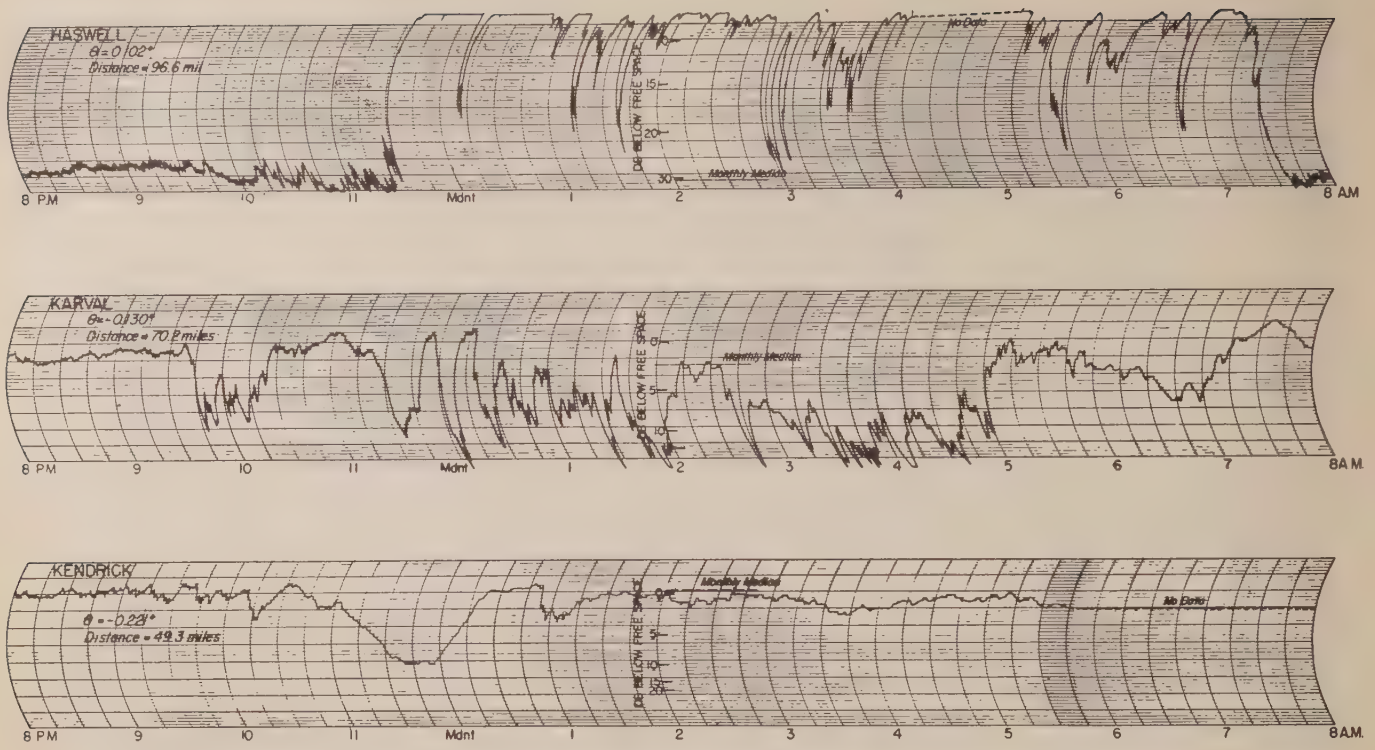

Fig. 8—1,046-mc fields for the night of June 21–22, 1952, at Kendrick, Karval, and Haswell, Colo.

analysis was then made of the radiosonde data of the U. S. Air Force Tornado Project for the period of the project, July 19 through September 26, 1952. The radiosonde data from two of these stations (see Fig. 1) La Junta (47 miles south of the path) and Akron (97 miles north of the path), were analyzed for ground modification and their refractivity profiles classified. The refractivity profiles were classified as (*a*) linear, provided the refractivity gradients were not in excess of $-0.050\ N$ units/meter (i.e., twice the average value for the period of record) and possessed no departures from this degree of linearity within 500 meters above the surface, or (*b*) ground modified, provided the refractivity gradients were at least $-0.050\ N$ units/meter or there existed a layer with a gradient of at least $-0.050\ N$ units/meter within 500 meters of the surface. Examples of this classification are found in Fig. 7 where the profiles taken at 1612 and 2035 are linear and those taken at 0945, 2320, 0250, and 0554 are ground modified.

To insure widespread horizontal homogeneity, the La Junta and Akron profiles were used only when both stations had similar profiles. The question then arises as to how long a period of time the profiles are representative of the area. Two analyses were made: (*a*) for a one-hour time interval with the radiosonde at the midpoint of the time interval, and (*b*) for a two-hour time interval with the radiosonde observations at the mid-point of the time interval. The nominal observation times are 0800 and 2000 MST.

Table I presents the results of these analyses and shows that, with ground modification present, fadeouts occurred within one hour either side of the radiosonde observation in 100 per cent of the cases; whereas, in the absence of ground modification (i.e., linear refractivity profile), fadeouts occurred in only 24 per cent of the observation periods. If, however, the time interval is shortened to one-half hour either side of the radiosonde observations fadeouts occurred in 75 per cent of the ground modified observation periods while fadeouts occurred in only 16 per cent of the observation periods with linear profiles.

TABLE I

DEPENDENCE OF PROLONGED SPACE-WAVE FADEOUTS AT KARVAL, COLO., ON THE REFRACTIVITY PROFILES AT LA JUNTA AND AKRON, COLO., FOR THE PERIOD JULY 19, 1952 TO SEPTEMBER 26, 1952.

Time Interval (Radiosonde at the mid-point)	Profile	No. of Observations	Per Cent Occurrence of Fadeouts	Per Cent Occurrence of No Fadeouts
2 Hour	Ground modified	20	100	0
	Linear	25	24	76
1 Hour	Ground modified	20	75	25
	Linear	25	16	84

At present the physical cause of these fadeouts is not clearly understood. The possibility that fadeouts are caused by simple direct- and ground-reflected-wave interference does not seem likely because calculations have proven unsatisfactory in explaining the depth of fadeouts as well as their radio frequency and θ dependence.

In the light of the above meteorological analysis, an extension of Doherty's[3] so-called "radio-hole" work would appear to be the most promising approach.

Conclusion

In conclusion, then, it would seem that the above analysis indicates that high-transmitter, low-receiver uhf links with θ near zero may be expected to display prolonged space-wave fadeouts providing that there are low-level modifications of the refractivity profile. Consequently, when planning uhf communication and navigational systems intended to operate throughout the twenty-four hours of the day in all seasons of the year, one should consider the refractivity climatology of the particular area and the transmitter power should be increased enough to insure the degree of reliability required. Detailed meteorological measurements are now being initiated in the Cheyenne Mountain propagation program with the objective of determining more precisely the mechanism causing these fadeouts and the feasibility of extending these conclusions to other areas.

[3] L. H. Doherty, "Geometrical Optics and the Field at a Caustic with Applications to Radio Wave Propagation Between Aircraft," School of Electrical Engineering, Cornell University, Ithaca, N. Y., Research Report EE 138; September 10, 1952.

Acknowledgment

The author would like to acknowledge the many helpful suggestions and criticisms of K. A. Norton during the preparation of this paper, and the work of G. J. Haines in the processing of the basic radio data.

CORRECTION

Daniel Levine, author of the paper, "Information Cells on Intensity-Modulated CRT Screens," which appeared on pages 1766–1768 of the December, 1953 issue of the Proceedings of the I.R.E., has brought the following correction to the attention of the editors.

The equations below should read

$$k\lambda > 4. \tag{5}$$

$$Q(x, y) = \frac{\tau\sqrt{\pi}}{k\lambda} I_M e^{-k^2 y}. \tag{6}$$

$$Q(x, 0) = \frac{\tau\sqrt{\pi}}{k\lambda} I_M.$$

$$i(x, y) = I_M e^{-3.67(x^2+y^2)/w^2}. \tag{8}$$

For the sake of completeness it is added that the total beam current, obtained by integration of (8), is

$$i_b = \frac{\pi}{3.67} I_M w^2 = 0.857 I_M w^2,$$

so that the equations may all be expressed in terms of the cathode current; thus, (6) becomes:

$$Q(x, y) = \frac{\sqrt{\pi}}{kv} I_M e^{-k^2 y^2}$$

$$= 1.08 \frac{i_b}{wv} e^{-3.67(y/w)^2}.$$

This form of the equation emphasizes the decrease of excitation for a long pulse when the line width (i.e. spot size) is increased for fixed values of cathode current and sweep speed.

Antenna-Scattering Measurements by Modulation of the Scatterer*

H. SCHARFMAN, ASSOCIATE, IRE,† AND D. D. KING, MEMBER, IRE

Summary—**An indoor method for measuring the radar cross section of objects by modulating the position of the object is described. The use of synchronous detection and employment of a substitution technique wherein the scattering from an object under test and a known object are measured and compared simultaneously reduce measurement error to ±2 per cent. The characteristics of the system are discussed, and its advantages and limitations are compared with other systems. Measured curves for the back-scattering cross-section of dipoles for lengths from 0.4λ to 1.6λ show good agreement with other published theoretical and experimental curves. Methods for measuring back-scattering and off-angle scattering from irregular objects by extension of this technique are indicated.**

I. Introduction

AS A CONSEQUENCE of its importance in the radar equation, the absolute value of the back-scattering cross section

$$\left(\sigma = 4\pi \frac{\text{Power scattered back to Source/unit solid angle}}{\text{Incident Power density}}\right)$$

of many objects has been studied both theoretically and experimentally for many years. The inherent mathematical difficulties in the determination of σ for all but the simplest objects place a premium on precise methods of measurement.

Several techniques have been used for the measurement of σ and related parameters. In recent years attention has turned to the use of an image plane for such measurement. Image plane methods have the advantage that errors caused by reflections and lense effects of supports are circumvented. King[1] has employed this device for determination of absorption gain and back-scattering cross section of dipoles. Dike and King[2] used a similar method for the study of the dipole, and Sevick[3] and Aden[4] made measurements on coupled dipoles and water spheres by means of this technique.

For many purposes these and simpler measurement schemes are adequate. However, use of the image-plane technique limits the investigation to objects having at least one axis of symmetry. Thus scattering investigations of the objects are confined to study of the back-scattering cross section.

In addition, it has been difficult to achieve high accuracy with the image-plane technique for several reasons. Errors arise due to drift of the source power and frequency, detector sensitivity variations, and noise. Reflections from surrounding objects and walls as well as the finite size and imperfections in the ground plane also cause errors.

As a result, errors of 5 per cent or more are quoted for previous measurements of back-scattering cross section (1, 3). The system described below avoids many of the sources of error mentioned, and also has the capability of measuring both σ and off-angle scattering from irregular objects.

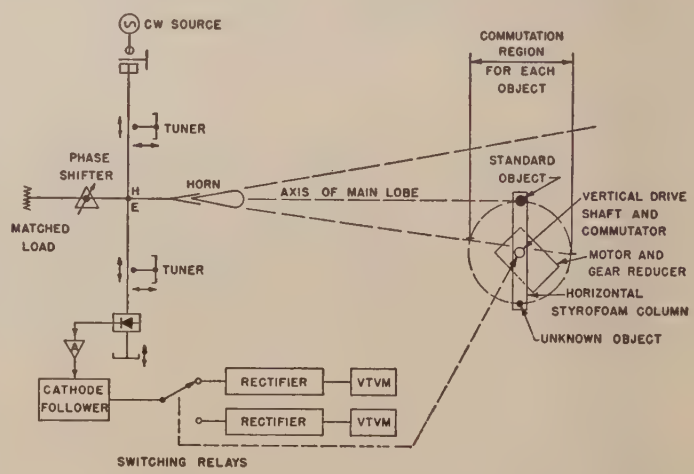

Fig. 1—System block diagram.

II. The Measuring Apparatus

The measurement system is shown in block diagram form in Fig. 1; photographs of the actual equipment set up are shown in Figs. 2, 3, 4 and 5. Referring to Fig. 1, a cw source feeds a slightly unbalanced hybrid tee; the power is divided about equally between the load and horn arms with a small amount of oscillator power entering the detector arm. The power emitted by the horn impinges on a rotating column of styrofoam (dielectric constant ≈1.03) in which are embedded a standard such as a metal sphere and opposite it, the object whose σ is to be determined. A synchronous motor and gear reducer drives a vertical wooden shaft which supports the styrofoam column. The horn is located so that its main lobe falls on only one object at a time while the rotation center is kept near the first null of the horn pattern.

As the styrofoam column rotates, each object in turn passes into the main lobe of the horn and scatters back to the horn a quantity of power proportional to its

* Decimal classification: R537.4. Original manuscript received by the I.R.E., August 11, 1953; revised manuscript received December 30, 1953.

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† Radiation Lab., Johns Hopkins University, Baltimore, Md.

[1] D. D. King, "The measurement and interpretation of antenna scattering," PROC. I.R.E., vol. 37, pp. 770–777; July, 1949.

[2] S. H. Dike, and D. D. King, "The absorption gain and back-scattering cross-section of the cylindrical antenna," PROC. I.R.E., vol. 40, pp. 853–860; July, 1952.

[3] J. Sevick, "An experimental method of measuring back-scattering cross-sections of coupled antennas," Cruft Lab., Harvard University, Cambridge, Mass, Report 151, May 28, 1952.

[4] A. Aden, "Electromagnetic scattering from metal and water spheres," Cruft Lab., Harvard University, Cambridge, Mass., Report 106, 1950.

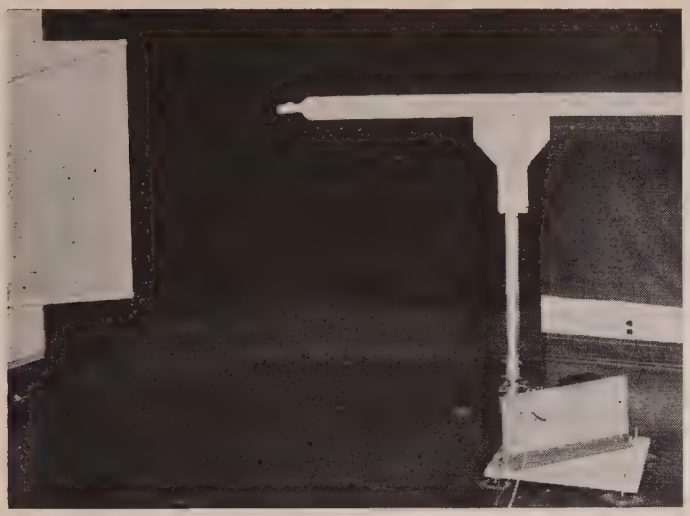

Fig. 2—View of styrofoam column and drive.

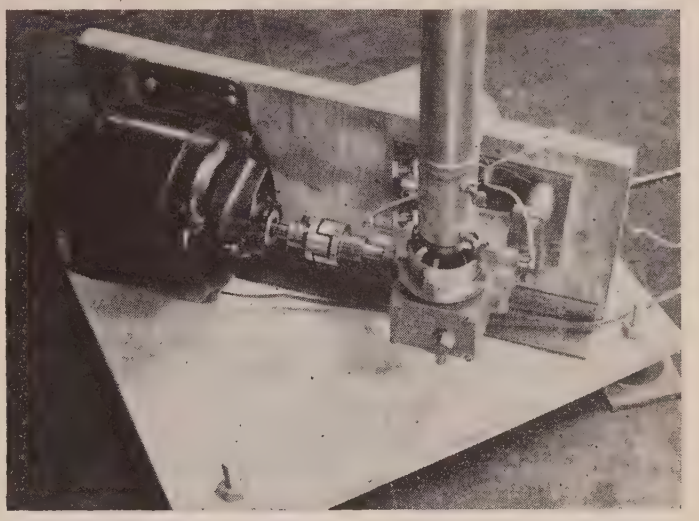

Fig. 3—Close-up of motor, gear reducer, and commutator.

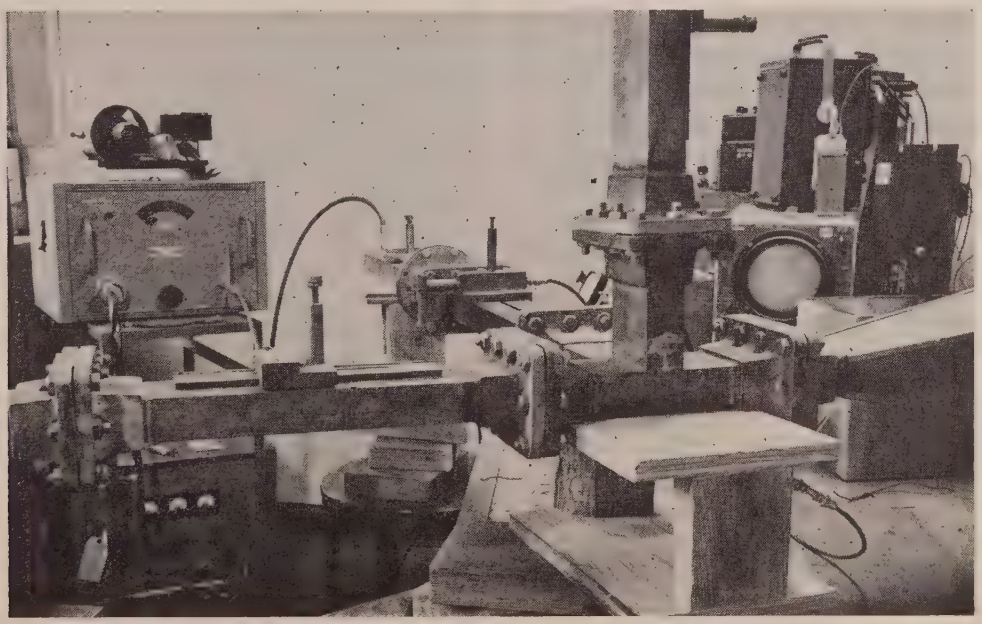

Fig. 4—View of oscillator and waveguide assembly.

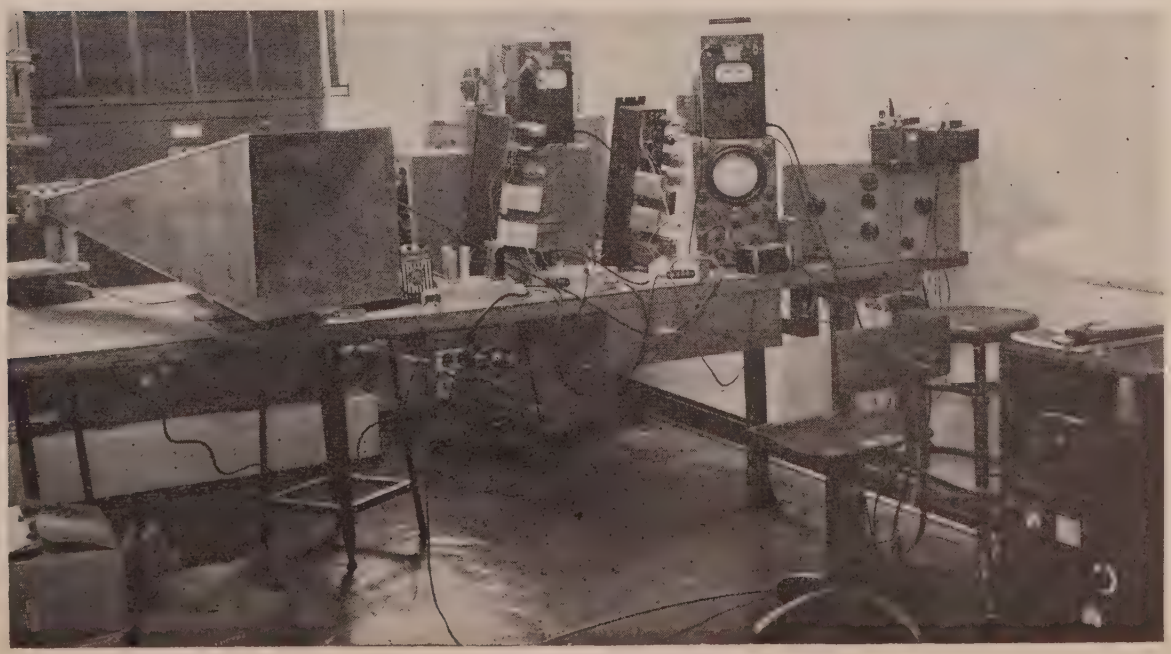

Fig. 5—View of horn and metering equipment.

back-scattering cross section. The portion of this signal reaching the detector arm mixes with the oscillator signal and is detected by a crystal. As an object moves into, through, and out of the main lobe, the power scattered back to the horn will vary in amplitude because of the lobe shape and the changing distance, but more important, the phase of the scattered field with respect to the oscillator field will change by π radians for each quarter wavelength of motion parallel to the axis of the horn. For a styrofoam column that is long compared to a wavelength, many phase reversals will occur over the semi-circular swing of the object through the main lobe. Thus, the phase of the back-scattered signal in the detector arm will vary rapidly with respect to the oscillator signal in that arm. The two signals are mixed and detected in the crystal, producing an audio signal corresponding in frequency to the doppler shift caused by the relative motion of the horn and object. In this case the detection process is that of a superheterodyne receiver with an IF equal to the doppler frequency and has been termed Synchronous Detection.[5,6] Under these conditions, the detector response is linear in that the amplitude of the audio voltage is proportional to the amplitude of the back-scattered field, and the sensitivity is about 30 db better than for direct video detection.

The output of the crystal over a full revolution of the styrofoam column consists of two signals each corresponding to the movement of one of the objects passing through the main lobe of the horn. These signals are independent of each other as only one object is in the main lobe at a time. Double reflection effects and the contribution from objects passing through the side lobes of the horn were observed to be negligible.

The combined signal from the crystal is amplified and the parts corresponding to the standard object and the object to be measured separated by means of relays and a commutator mounted on the vertical drive shaft of the styrofoam column. Commutation is arranged to occur at the time that objects are entering and leaving the main lobe. The separated doppler signals are rectified and averaged over many revolutions of the support column by a long time-constant filter. The two resulting dc voltages appear on a pair of vacuum-tube voltmeters. Each voltage is proportional to the back-scattered field of its respective object. Taking the square of the ratio of the voltages removes the proportionality constants and yields the ratio of the back-scattering cross section of the two objects. The standard scatterer (usually a metal sphere) has a σ that is known exactly[7,8,9] and the σ for the object under test is determined directly from the measured ratio and known value for the standard.

[5] M. E. Brodwin, C. M. Johnson, and W. M. Waters, Technical Report 18, Radiation Laboratory, The Johns Hopkins University, Baltimore, Md.; March 31, 1952.
[6] M. E. Brodwin, C. M. Johnson, and W. M. Waters, "Low level synchronous mixing," 1953 I.R.E. Convention Record, part 10, p. 52.
[7] J. A. Stratton, "Electromagnetic Theory," p. 563 ff, McGraw-Hill Book Co., New York, N. Y.
[8] L. N. Ridenour, "Radar System Engineering," p. 64, McGraw-Hill Book Co., New York, N. Y.
[9] D. Kerr, "Propagation of Short Radio Waves," p. 453, McGraw-Hill Book Co., New York, N. Y.

III. Characteristics of the System

The system as outlined above has several important advantages over others previously employed. As only the ac output of the crystal corresponding to moving objects is amplified and metered, only objects moving through the beam are detected. Reflections from stationary objects including the walls affect only the initial adjustment of the Tee unless the object approaches so close to a highly reflecting object that multiple reflections cannot be neglected. The latter effect is made small by keeping the styrofoam column many wavelengths from the nearest wall and lining the wall opposite the main lobe with absorbing material. The reflected power from this wall is at least 17 db below the incident power at normal incidence. In addition, the specular reflection from the wall opposite the horn was reduced by suitable orientation of the horn with respect to the wall and the styrofoam column.

Potential sources of error arising from variations in oscillator power, detector sensitivity, line voltage and aging effects are avoided by using the previously described time-sharing technique. The scattered field from each object is subject to the same drift effects, and these are eliminated when the ratio of the voltage readings is taken in the process of calculating $\sigma_{test}/\sigma_{standard}$. Short-time perturbations are removed by the long-time constant filters at the rectifier outputs.

As described above, synchronous detection is used in the measurements, and this inherently leads to several advantages. It has been shown[5,6] that about 1 dbm of reference power in the detector arm is needed for good synchronous detection, and that the detection sensitivity is essentially constant over a wide range of reference power. The leniency of these conditions allows the Tee to be properly adjusted by tuning a phase shifter in the matched load arm. This adjustment is insensitive and indicates that the system can be used over wide frequency bands with minor modifications. The crystal detector mount is the only relatively narrow-band element in the system, and this is readily modified for use in different ranges.

By virtue of its linear response, the synchronous detection process indirectly reduces other errors in the system. The crystal-output voltage is proportional to the back-scattered field strength rather than power (as in low-level video detection) and, consequently, the required linear dynamic range of the amplifiers, rectifiers, and voltmeters is reduced. Errors arising from nonlinearities in the latter devices are thus reduced.

IV. Limitations and Sources of Error

It has been stated above that this system is inherently wide band; the frequency limits will now be considered. The lowest frequency is limited by the physical size of the available range, for the object must always be kept in the Fraunhofer zone (unless one wishes to specifically measure Fresnel-zone scattering). Lower frequencies generally mean physically larger objects and more cumbersome means of getting the standard and unknown ob-

jects in and out of the beam. For outdoor ranges, provision must be made to keep moving objects out of the main lobe, and consideration must be given to the effect of moving objects close to the transmitter in its side lobes. These considerations would appear to limit the minimum frequency to roughly 1,000 mc.

The limitation at high frequencies lies in the support column which inevitably reflects energy and may act as a lens on the standard and unknown objects. At 9,000 mc styrofoam exhibits both of these faults, but measurements at 3,000 mc indicate negligible effects due to the styrofoam. The high frequency limit for high accuracy measurements, therefore, might be about 5,000 mc although usable results could perhaps be obtained at higher frequencies with a specially made support column.

Errors caused by curvature in the phase front of the incident wave over the objects are partly removed when the ratio of the meter readings is taken. This arises from the fact that slight curvature of the phase front will reduce the measured back-scattered signal for both objects.[10] As the error in each signal is in the same direction, the error in their ratio will be smaller than the error in either signal. For objects of approximately the same scattering cross section, the errors tend to cancel completely with high resultant accuracy in the ratio of the signals. By varying the range, it was found that for this system negligible error was incurred if

$$R \geqq \frac{(D+d)^2}{\lambda},$$

D = maximum dimension of horn in wavelengths
d = maximum dimension of scatterer in wavelengths
λ = wavelength
R = distance from horn to object on axis of main lobe.

The effects of nonuniform incident amplitude over the path traced by the scatterer are also partially cancelled by taking the ratio $\sigma_{\text{test}}/\sigma_{\text{standard}}$.

Errors which cannot be excluded are due to nonlinearity of the rectifiers and meters and the meter-reading error which combined should be less than 2 per cent. From intercomparison tests and the excellent repeatability, the over-all probable error in the ratio of the readings for two objects is estimated to be less than ± 1 per cent. As this ratio must be squared to obtain the ratio of the back-scattering cross section, the probable error in the latter ratio is less than ± 2 per cent. The residual errors become more significant for values of σ/λ^2 less than 0.1.

V. Experimental Results

The system of Fig. 1 was set up at 3,000 mc, and the back-scattering cross section of copper dipoles of lengths varying up to $3\lambda/2$ was measured as compared to a silver-painted sphere of radius 1.50 inches. The value of the back-scattering cross section for a sphere as given in the literature[8,9] was used as the standard. Back-scattering from various sizes of silver-painted spheres was also measured as a check.

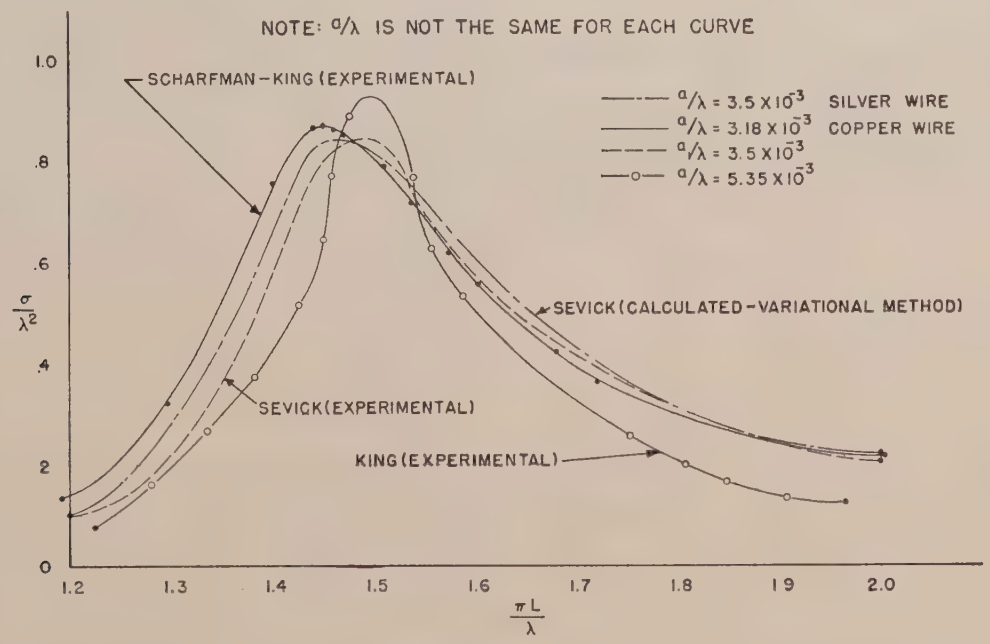

Fig. 6—Back-scattering cross section versus dipole length.

In Fig. 6 the measured curve of back-scattering cross section versus dipole length in the vicinity of the first resonance is plotted together with the measured curves of Sevick[11] and King[1] as well as the theoretical curve obtained by Sevick[11] and Tai[11] using the variational method.

In Fig. 7 a measured curve covering the first two dipole resonances is presented along with the measured curves of Dike and King,[2] Sevick,[11] and a theoretical curve by Tai.[12] The absence of a dip before the second

[10] E. H. Braun, "Gain of electromagnetic horns," Proc. I.R.E., vol. 41, pp. 109–115; January, 1953.

[11] J. Sevick, "Experimental and theoretical results on the back-scattering cross-section of coupled antennas," Cruft Lab., Harvard University, Cambridge, Mass. Report 150; 1952.

[12] C. T. Tai, "Electromagnetic back-scattering from cylindrical waves," *Jour. of Appl. Phys.*, vol. 23, pp. 909–916; August, 1952.

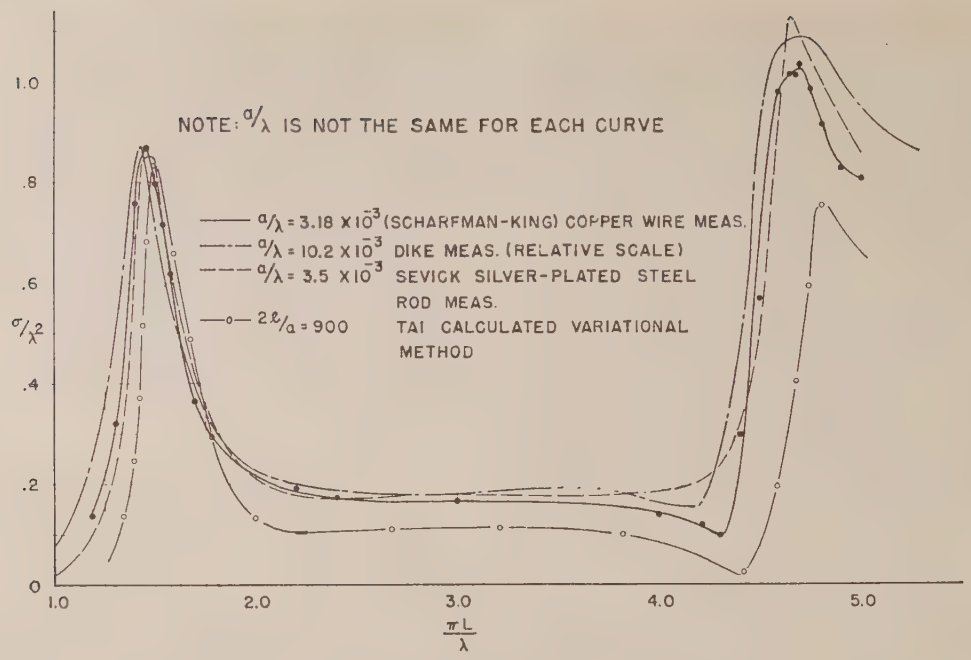

Fig. 7—Back-scattering cross section versus dipole length.

resonance in Sevick's data is ascribed to the high probable error at low values of σ reported by Sevick.

VI. Extensions of the Measurement Method

The technique described above for the measurement of back-scattering cross section has the limitation in common with image-plane techniques that only objects having at least one axis of symmetry can be measured.

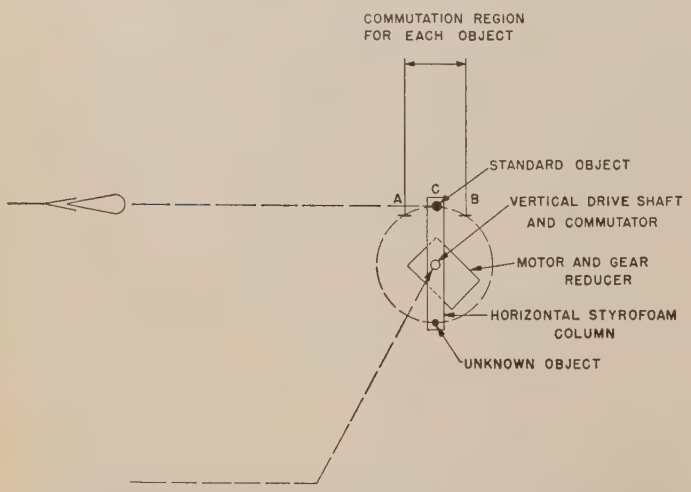

Fig. 8—System for measuring back-scattering from irregular objects.

This restriction is lifted when the commutation period is reduced so that only the signal generated near the center of the main lobe is applied to the rectifier. Scattering from irregular objects may be measured, as indicated in Fig. 8. The orientation of an object with respect to the horn is almost constant from A to B, and the error introduced by slight change in orientation from A to C to B can be varied by changing the commutation period within limits or by increasing the distance from horn to object. Off-angle scattering for arbitrary angles with respect to the horn may be measured by separating the location of receiving and transmitting horns. Reference power for synchronous detection is then supplied from a directional coupler to the remote receiver, and no hybrid junction is required.

The technique of modulating the scatterer is also applicable to image-plane systems. For absorption and loading measurements, as well as for freedom from support problems, the image system is preferable. The standard and test objects would then be mounted in a rotating disk in the image surface, or passed through the beam in some other fashion.

VII. Conclusions

It has been shown that the system of modulating the scatterer removes many sources of error in back-scattering measurements. Using the described technique, the room-reflection problem is reduced, much of the usual adjustment and tuning difficulties are avoided, and drift caused by electrical and atmospheric changes is eliminated. In addition, the method can be applied over a wide band of frequencies, and has high sensitivity. Data taken on dipole scattering confirm the accuracy of the method, and is in substantial agreement with the results of previous investigations.

With modification the system could be used to measure off-angle scattering, and, by combining some of the features of an image plane, the absorption cross section and the back-scattering cross section of terminated antennas could be measured.

VIII. Acknowledgment

C. F. Miller, M. E. Brodwin, and C. H. Grauling made valuable suggestions to the authors in the course of the work.

Correspondence

Improvement of Power Output from Pulsed Klystrons*

A recent paper[1] described the design and performance of high-power pulsed klystron amplifiers which are used on the linear electron accelerator at Stanford University. Recent tests on a modified tube of this type have given considerably higher outputs and improved efficiencies.

Early measurements showed that only 70 per cent of the beam was reaching the collector in this tube. Tests made on a beam tester showed that part of the beam was lost at the anode surface, and that by increasing the cathode-anode distance slightly and enlarging the taper at the entrance to the drift tube, the transmission would be greatly improved. These changes, made on a tube which had been provided with an insulated collector, resulted in a transmission of 90 per cent. The actual transmission through the active portion of the tube is somewhat higher, since the insulated collector is located some distance from the last cavity in the tube.

However, with these changes the tube oscillated violently at a frequency of about 5,800 mc (which is not a harmonic frequency). This oscillation made it impossible to measure the performance accurately, since directional couplers were used for power measurement. Cold tests of the cavity showed that a higher mode of the correct frequency was responsible for the oscillation. The feedback mechanism has not been determined; it probably takes place either by fast secondaries or by rf coupling through the drift tube. The latter seems improbable, since the attenuation between cavities is greater than 25 db at this frequency.

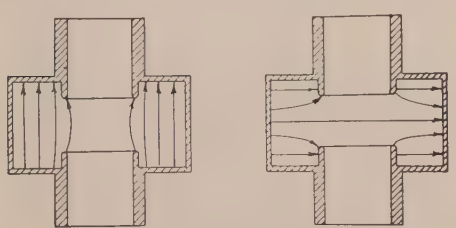

Fig. 1—Cavity modes at 10.5-cm (left) and 5.2-cm wavelengths.

The two cavity modes are shown in Fig. 1. By inserting two plates as shown in Fig. 2, it was possible to suppress the oscillation mode without affecting the fundamental resonance by a significant amount. This modification (made in the middle cavity, which has the highest Q) has effectively stopped the oscillation.

The modifications have been made on two tubes, one of which has been tested at a beam voltage of 360 kv. At this level, the tube gave an output of 30 megw with an efficiency of 41 per cent. The rf pulse length was 1.4 μsec and the repetition frequency was 60 cps. (The beam pulse length was 2 μsec, so that a 2-μsec rf pulse could have been used, had a driver been available.) The second tube was tested to 25 megw and then placed in operation on an accelerator, where it has been in use for the past six months.

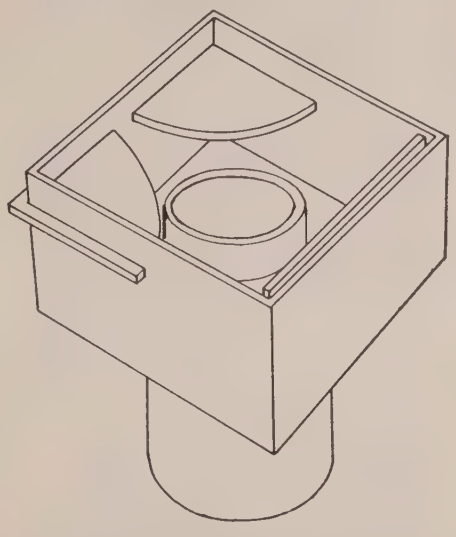

Fig. 2—Cut-away view of cavity, showing oscillation suppressors.

It is believed that the power limit on this tube has not been reached, since unmodified tubes have been operated as diodes up to 400 kv at 1.5 μsec. The present modulator does not allow variation of the pulse rate, but it should be possible to increase the rate somewhat, since the cooling is adequate. Increasing the pulse length, on the other hand, would almost certainly require redesign of the collector, to prevent melting; possibly the cathode would also need modification. The efficiency should remain about the same, since it was increasing slightly at the highest voltage used.

The above investigations were carried out at Stanford University under the sponsorship of the Office of Naval Research, Contract N6onr 25123. The cold testing and modification of cavities was carried out by K. L. Brown.

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* Received by the Institute, February 1, 1954.

[1] M. Chodorow, E. L. Ginzton, I. R. Neilsen, and S. Sonkin, "Design and performance of a high-power pulsed klystron," PROC. I.R.E., vol. 41, pp. 1584–1602; November, 1953.

A Simple Graphical Analysis of a Two-Port Waveguide Junction*

The results presented in the subject paper;[1] evaluation of scattering coefficients and measurement of a load through a junction, are justified[2] by means of elementary geometry. The authors seem to believe that justification at a higher level could make the constructions for the second problem more complicated. It is doubtful, however, whether the average engineer is more familiar with the properties of inversion than with those of so-called higher geometry. He could probably become acquainted just as easily with the latter, and thus improve his ability to solve other problems on waveguide junctions. The main acquisition would be the concept of hyperbolic distance.[3] This special distance does not have to be introduced, any more than ordinary distance, by way of group theory or other relatively obscure mathematical techniques; it is merely the result of taking a reading with the hyperbolic protractor.[4]

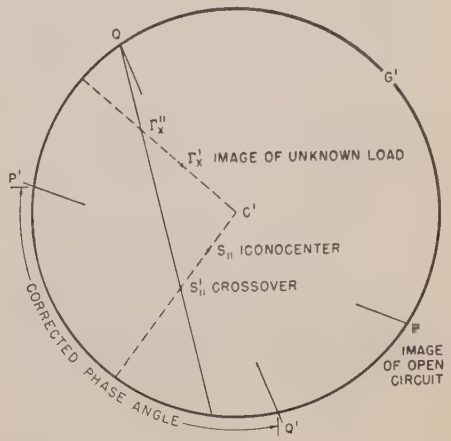

Fig. 1

With these tools available, the procedure for interpreting the measurements of an obstacle through a junction, lossy or not, is greatly simplified.[5] Using, to make comparison easier, the notation of Ref. 1, Fig. 4 of that reference is reduced to the accompanying Fig. 1. The point Γ_x'' is constructed by the transformation $\mathcal{B}$ relative to G' (Ref. 3, section 1). The corrected VSWR in decibels is then the hyperbolic distance between S_{11}' and Γ_x'', and can be obtained by a single reading with the protractor. (Compare to formula for $|\Gamma_x|$ on Fig. 4, Ref. 1, which involves 3 measurements of length.) The corrected phase angle is $P'C'Q'$, obtained by projecting P and Q through S_{11} (no necessity to determine first the phase angle of S_{22}).

It should be noted that since this construction always remains inside the circle G', there is no need for an alternative construction such as that of Fig. 5, Ref. 1.

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* Received by the Institute, November 23, 1953.

[1] J. E. Storer, L. S. Sheingold, and S. Stein, PROC. I.R.E., vol. 41, pp. 1004–1013; Aug., 1953.

[2] An analytical justification, also elementary, of the evaluation of the scattering coefficients can be found in, G. A. Deschamps, "Determination of the reflection coefficients and insertion loss of a waveguide junction," *Jour. of App. Phys.*, vol. 24, pp. 1046–1050; Aug., 1953. It was also available in the reports quoted in Footnote 1 of Ref. 1.

[3] G. A. Deschamps, "New chart for the solution of transmission line and polarization problems," *Trans. I.R.E. Professional Group on Microwave Theory and Techniques*, vol. 1, pp. 5–13; March, 1953. Also *Elect. Commun.*, vol. 30, pp. 247–254; Sept., 1953.

[4] G. A. Deschamps, "A hyperbolic protractor for microwave impedance measurements and other purposes," International Tel. and Tel. Corp., 67 Broad Street, New York, N. Y.; 1953.

[5] Ref. 3, sect. 6 and 7.

Correspondence

One Way Transmission Devices*

It might seem possible to use the nonreciprocal properties of ferrites to make lossless one-way transmission devices at microwave frequencies. For example, in the two-terminal-pair circuit shown below, energy fed in at terminals 1 would be transmitted to terminals 2 with minimum loss, whereas energy impressed on terminals 2 would be totally reflected. However, this characteristic is impossible because of the principle of conservation of energy. This can be proven in Fig. 1 above:

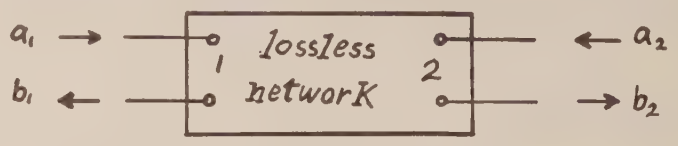

Fig. 1

If our network is lossless, all the energy entering the network minus the energy leaving must total up to zero. In equation form this can be written as:

$$\sum_{n}^{2} (a_n a_n^* - b_n b_n^*) = 0. \qquad (1)$$

The quantity, a_n, represents the amplitude of the incident wave entering the nth terminal, whereas the quantity, b_n, represents the amplitude leaving the nth terminal. The relation between the entering and the leaving amplitudes can be expressed as follows:

$$\begin{aligned} b_1 &= s_{11}a_1 + s_{12}a_2 \\ b_2 &= s_{21}a_1 + s_{22}a_2. \end{aligned} \qquad (2)$$

Since the terms $a_1a_1^*$, $a_2a_2^*$, etc., are independent, their coefficients obtained in the substitution of (2) into (1) must be equal to zero or:

$$\begin{aligned} 1 - s_{11}s_{11}^* - s_{21}s_{21}^* &= 0 \\ 1 - s_{22}s_{22}^* - s_{12}s_{12}^* &= 0 \\ s_{11}s_{12}^* - s_{21}s_{22}^* &= 0 \\ s_{12}s_{11}^* - s_{22}s_{21}^* &= 0. \end{aligned} \qquad (3)$$

These are the unitary relations shown in the reference.[1]

Let it be assumed that terminal pair 1 is matched for a matched load on terminal pair 2, i.e., $s_{11}=0$. Then it follows from the first unitary equation that:

$$|s_{21}| = 1.$$

The third equation shows that

$$|s_{22}| = 0$$

or that no reflections appear at terminal 2 if a matched load is placed on terminal one. Further, it is seen that

$$|s_{21}| = |s_{12}| = 1$$

which means that if the device transmits in one direction, it must transmit for the other direction as well.

If the network is a perfect reflector at one terminal, i.e., $|s_{11}|=1$, it can be shown that a perfect reflection occurs at the other terminal.

As a converse to the above, a passive one-way transmission device must have dissipation. Nonreciprocal devices having minimum forward loss and high absorption in the reverse direction have been demonstrated and do not contradict the above explanation.

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* Received by the Institute, January 25, 1954.

[1] C. G. Montgomery, et al., "Principles of Microwave Circuits," vol. 8, Radiation Lab. Series, p. 149; 1948.

Stretched-Log Frequency Axis*

In wide-band amplifiers with six, or more, frequency decades, a logarithmic scale frequently prevents adequate plotting of hf wiggles and cut-off conditions, as is born out by the exaggerated distributed amplifier response curve in Fig. 1(a). As a remedy, the response curve is conventionally broken off at some point and continued linear, as is shown in Fig. 1(b). Occasionally, a wiggle in the curve gets lost by this procedure, and if the decade Δf_{log} fits the plotting paper, the decade Δf_{lin} may not fit it.

If the cut-off frequencies are f_1 and f_2 and the frequencies at the discontinuity f' and f'', the equation from which a curve of this type may be plotted is of the general form

$$x = \log f \left|\frac{f'}{f_1} + Kf\right|\frac{f_2}{f''},$$

where K is a constant and x the linear scale co-ordinate. If the discontinuity here could be approximated by a continuous function, a practical formula for scale division may result from which curves similar to the one in Fig. 1b may be plotted, *free from discontinuity.*

The following suggestion for a solution dodges the discontinuity by replacing the linear scale with a stretched logarithmic scale, see Fig. 1(c), the stretching being accomplished by doubling (or quadrupling) the length of the decade from a chosen mid-frequency, here mc. As is seen, the resulting continuous curve is a good approximation of the discontinuous curve in Fig. 1(b). If the decade Δf_{log} fits the plotting paper, the decade $\Delta f_{str.log}$ also fits the plotting paper. If it is desirable also to show the lower cut-off region in detail, two-way stretching from the assumed mid-frequency may be employed, leaving the least interesting, flat portion of the response curve compressed to a small part of the graph. Modern 1,000-mc distributed amplifiers have such a region, which is made insignificant by this method, while the important wiggles and drops are enhanced in continuous presentation. The same technique applies to associated phase and time delay characteristics.

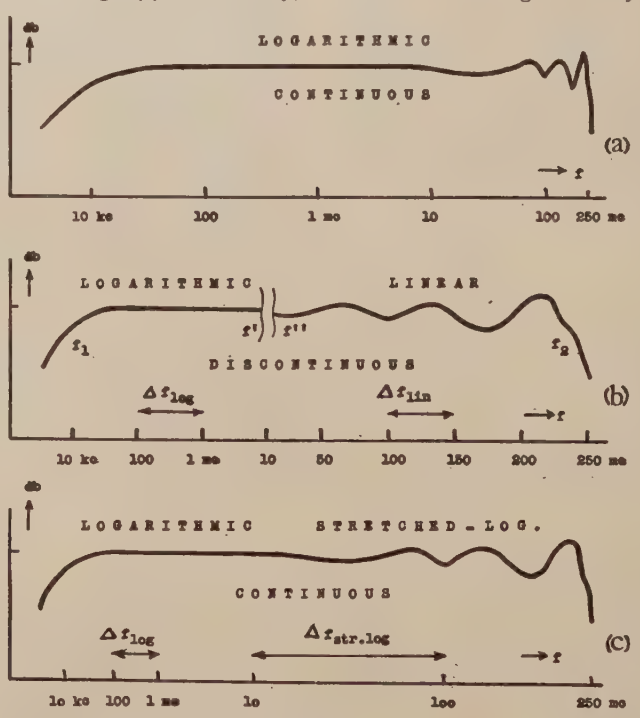

Fig. 1

To the extent printed stretched logarithmic paper is available, the above procedure simplifies to straighforward logarithmic plotting.

* Received by the Institute, July 17, 1953.

Correspondence

One may argue the point that the curve in Fig. 1(c) is not continuous, as it appears to the eye, due to the fact that in the transformation from decade to decade, the coefficients for the higher order derivatives in the appropriate Taylor's series obtain new numerical values. This is not true, for we are still dealing with the original analytical function, but the eye must get used to reading scale-implied slope changes without hesitation.

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A New Method of Driving-Point Impedance Synthesis*

R. Bott and R. J. Duffin[1] have indicated a method of driving point impedance synthesis that eliminates the coupled inductors required in the Brune realization of a minimum immitance function. Their technique is to solve for $Z(p)$ in the function

$$R(p) = \frac{kZ(k) - pZ(p)}{kZ(k) - Z(k)p}, \qquad (1)$$

which Richards[2] had showed to be positive-real if $Z(p)$ is positive-real. $Z(p)$ is thereby written as a sum of two positive-real functions, each of which can be reduced to an order two less than the order of $Z(p)$.

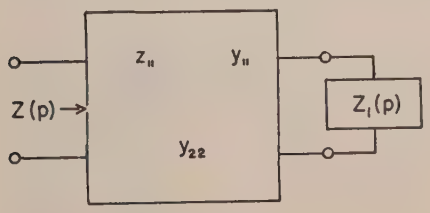

Fig. 1

Richards's function can also be used quite differently in a two-terminal-pair approach to driving-point impedance synthesis. In Fig. 1

$$Z_1(p) = \frac{z_{11}}{y_{22}} \frac{[1 - Z(p)y_{11}]}{[Z(p) - z_{11}]}. \qquad (2)$$

Given that $Z(p)$ is minimum, and $ReZ(j\omega_0) = 0$, let $Z(j\omega_0) = j\omega_0 L_0$, and assume $L_0 > 0$. Arbitrarily, let

$$z_{11} = pL_0 + \frac{k^2 L_0}{\dfrac{k^2 L_0 - pZ}{Z - pL_0} + \dfrac{\omega_0^2 p}{p^2 + \omega_0^2}} \qquad (3)$$

in which

$$k = \frac{Z(k)}{L_0}.$$

Since

$$\frac{k^2 L_0 - pZ}{Z - pL_0}$$

is a Richards function, z_{11} is positive-real. Now z_{11} can be synthesized as in Fig. 2. Let

$$\frac{k^2 L_0}{\dfrac{k^2 L_0 - pZ}{Z - pL_0} - \dfrac{Ap}{p^2 + \omega_0^2}} = Z_2(p). \qquad (4)$$

Here A is a constant which, if the elements of z_{11} are to be positive, must conform to

$$A \leqq 2 \text{ Residue} \left[\frac{k^2 L_0 - pZ(p)}{Z(p) - pL_0}\right]_{p=j\omega_0} \quad \text{for } A > 0, \qquad (5)$$

or

$$A > (-\omega_0^2) \quad \text{for} \quad A < 0.$$

$Z_2(p)$ has in general a zero at $p = j\omega_0$ and thus can be reduced to an order two less than the order of $Z(p)$.

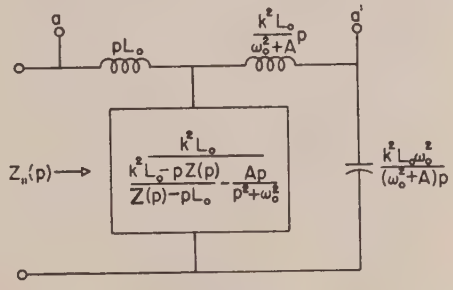

Fig. 2

$Z_1(p)$ can now be calculated with some labor from (2), this terminating network being placed across $a-a'$:

$$Z_1(p) = \frac{L_0[L_0 k^2 - pZ(p)]}{Z(p) - pL_0} + \frac{k^2 L_0 Ap}{(\omega_0^2 + A)(p^2 + \omega_0^2)}. \qquad (6)$$

If

$$\frac{k^2 A}{\omega_0^2 + A} \geqq -2 \text{ Residue} \left[\frac{L_0 k^2 - pZ(p)}{Z(p) - pL_0}\right]_{p=j\omega} \qquad (7)$$

then $Z_1(p)$ is positive real. $Z_1(p)$ has in general a pole on the $j\omega-$axis at $p = j\omega_0$ and so can be reduced to an order two less than the order of $Z(p)$. There are two important values for A:

I. $$\frac{k^2 A}{\omega_0^2 + A} = -2 \text{ Residue} \left[\frac{L_0 k^2 - pZ(p)}{Z(p) - pL_0}\right]_{p=j\omega_0}. \qquad (8)$$

Although A is negative, it is obvious from (8) that $A > (-\omega_0^2)$ and so the elements of Fig. 2 are positive.

II. $$A = 2 \text{ Residue} \left[\frac{k^2 L_0 - pZ(p)}{Z(p) - pL_0}\right]_{p=j\omega_0}. \qquad (9)$$

Under condition I, the pole at $j\omega_0$ in $Z_1(p)$ is removed and so $Z_1(p)$ is of order two less than that of $Z(p)$; $Z_2(p)$ can be reduced as previously indicated. Under condition II, zero at $p = j\omega_0$ is removed and so $Z_2(p)$ is of order two less than that of $Z(p)$; $Z_1(p)$ can be reduced as indicated before.

In each case the reduction is accomplished with one less element per reduction cycle than in the Bott-Duffin method, and

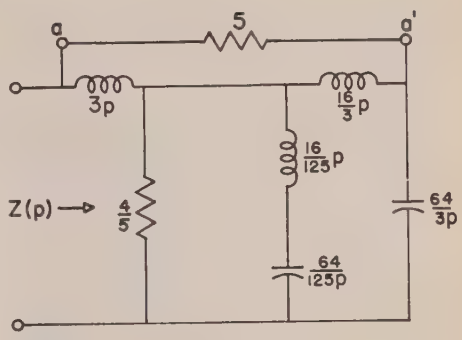

Fig. 3

no mutual inductance is used. If $L_0 < 0$, duality leads to a corresponding realization.

As an example, let

$$Z(p) = \frac{5p^2 + 18p + 8}{p^2 + p + 10}.$$

From which

$$\omega_0 = 2, \quad L_0 = 3, \quad \text{and} \quad k = 2/3,$$

all units being in the usual practical MKS system for impedance. The realization under condition I is shown in Fig. 3; that under condition II in Fig. 4; the Bott-Duffin synthesis leads to Fig. 5.

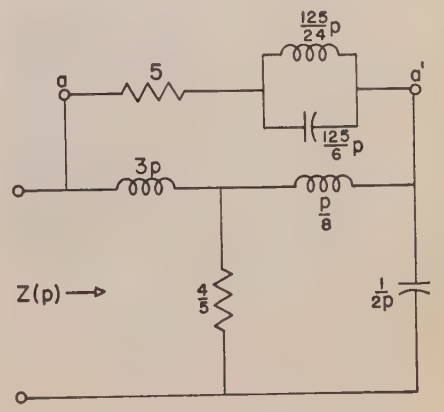

Fig. 4

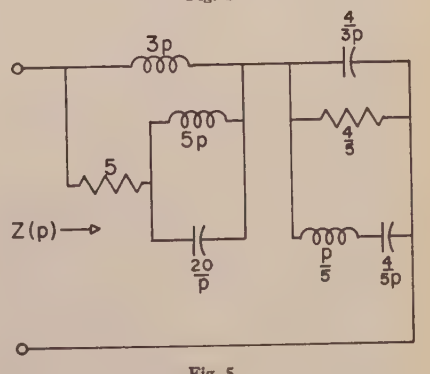

Fig. 5

This procedure is a result of work done with the support of the U. S. Army, U. S. Air Force, and U. S. Navy. A more detailed discussion of this and other new synthesis techniques is in preparation.

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* Received by the Institute, January 19, 1954.

[1] R. Bott and R. J. Duffin, "Impedance synthesis without use of transformers," *Jour. Appl. Phys.*, vol. 20, p. 816; Aug., 1949.

[2] P. I. Richards, "A special class of functions," *Duke Math. Jour.*, vol. 14, pp. 777–786; 1947.

Correspondence

The Effect of Space Charge on Beam Loading in Klystrons*

In the usual derivation of klystron beam loading, the effect of space charge is neglected.[1] It is desirable to know whether this assumption is justified and to determine under what beam conditions space charge can no longer be ignored. An interesting by-product of this investigation is a new explanation for the excessive beam loading caused by secondary electrons.

A very general solution of the diode loading problem has been obtained by Llewellyn and Peterson[2] which is readily applicable to the problem at hand. Their equation (16) gives the impedance of a diode with any degree of space charge between the input and output planes. The last term of (16) may be written as $-i2\zeta(1-\zeta)(\sqrt{V_{Da}}+\sqrt{V_{Db}})^2/I_D\theta$ in which case the general diode impedance is given by

$$Z = r_0\left[\frac{r}{r_0}+i\frac{x}{r_0}-i\frac{3(1-\zeta)}{\zeta}\cdot\frac{1}{\theta}\right], \quad (1)$$

where

$$r_0 = 2\zeta^2(\sqrt{V_{Da}}+\sqrt{V_{Db}})^2/3I_D$$

$$\frac{r}{r_0} = 12\left[2(1-\cos\theta)-\theta\sin\theta\right]/\theta^4$$

$$\frac{x}{r_0} = -12\left[(\theta^3/6)+\theta(1+\cos\theta)-2\sin\theta\right]/\theta^4$$

V_{Da} = dc voltage at entrance plane a

V_{Db} = dc voltage at exit plane b

I_D = dc current density, amperes/cm²

θ = electron transit angle between a and b planes.

The factor ζ specifies the degree of space charge between the a and b planes. It is related to the ratio of the direct-current density I_D to the limiting-current density that can pass between the planes without electron reversal I_m. Llewellyn and Peterson[2] give a cubic equation relating ζ and I_D/I_m. This equation may be solved algebraically for ζ. Discarding extraneous roots, we find

$$\zeta = \begin{cases} 2\left[1-\cos\left(\frac{1}{3}\cos^{-1}\left|\,1-2I_D/I_m\right|\right)\right] \\ \qquad 0\leqq I_D/I_m\leqq\tfrac{1}{2} \\ 2\left[1+\cos\left(\frac{1}{3}\cos^{-1}\left|\,1-2I_D/I_m\right|\right.\right. \\ \qquad \left.\left.+120°\right)\right] \quad \tfrac{1}{2}\leqq I_D/I_m\leqq 1. \end{cases} \quad (2)$$

The limiting-current density is given by the expression

$$I_m = \frac{2.33}{10^6}\,\frac{(\sqrt{V_{Da}}+\sqrt{V_{Db}})^3}{X^2} \quad (3)$$

where I_m is in amperes/cm², V_D is in volts, and gap spacing X is in cm.

In a klystron gap it is customary to calculate admittance rather than impedance, hence it is useful to invert and rationalize (1). If we let $V_{Da}=V_{Db}=V_D$, the conductive component of the admittance may be written as

$$g = \frac{I_D}{V_D}\,\frac{3}{8}\,\frac{1}{\zeta^2}\cdot\frac{r}{r_0}\Big/\left[\left(\frac{r}{r_0}\right)^2+\left[\frac{x}{r_0}-\frac{3(1-\zeta)}{\zeta}\,\frac{1}{\theta}\right]^2\right]. \quad (4)$$

For small space charge the last term in the denominator predominates and (4) becomes simply

$$g = \frac{I_D}{V_D}\,\frac{1}{2\theta^2}\left[2(1-\cos\theta)-\theta\sin\theta\right]\frac{1}{(1-\zeta)^2}.$$

Except for the last factor, this is identical with the beam loading result derived by the nonspace-charge method. Thus the presence of a small amount of space charge merely scales up the customary loading by the factor $1/(1-\zeta)^2$.

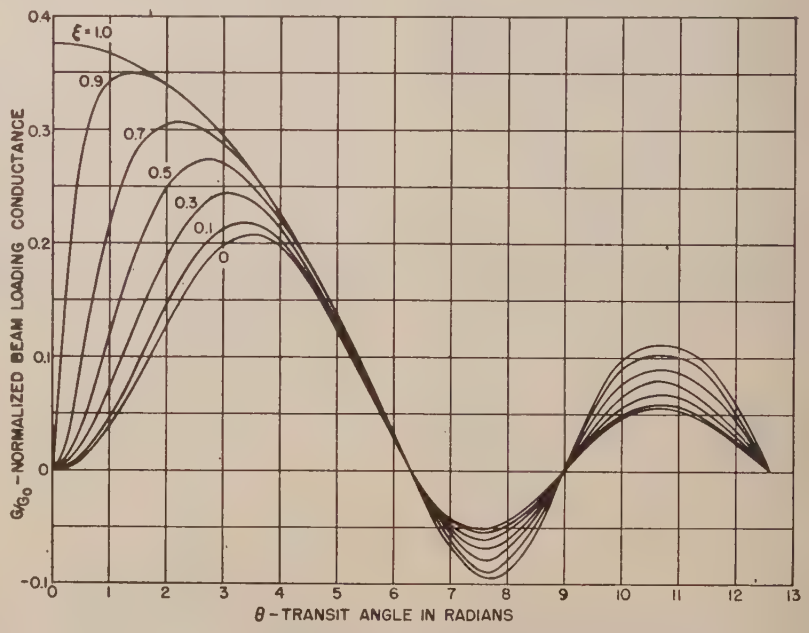

Fig. 1—Normalized beam loading conductance for various degrees of space charge.

Equation (4) is plotted in Fig. 1 for several values of the parameter ζ. The curve for $\zeta=0$ is the familiar primary beam-loading curve for klystrons. The most interesting characteristic of the other curves is the order of magnitude increase of beam loading at small transit angles and large degrees of space charge. It is also significant to note that when complete space charge prevails a virtual cathode is established and loading takes on a new characteristic. No longer is transit time necessary for loading to take place. Indeed, loading is most severe at zero transit angle. This, of course, is typical of the loading found in diodes and grounded-grid triodes where complete space-charge conditions usually exist.

To see whether space charge alters primary beam loading appreciably in a klystron gap, it is necessary only to calculate ζ for typical conditions. Let $V_D=1{,}000$ volts, and $X=0.100$ inch; then from (3) $I_m=9.13$ amps/cm². If a typical value for I_D is 0.1 amp/cm², $I_D/I_m=0.011$, and from (2) $\zeta=0.0049$. According to (5), the space-charge correction is only 1 per cent under these conditions, so space charge is truly negligible in the calculation of primary electron beam loading.

It is interesting to consider the effect of space charge on secondary electron beam loading. The majority of secondary electrons in a small-signal gap travel with extremely small velocities. As a first approximation let us consider that all secondaries are emitted with 5 volts of energy into a 0.100 inch gap. From (3) the limiting current density for these electrons is 3.24 ma/cm². This current density would be attained if a 100 ma/cm² primary beam suffered only 3.24 per cent interception. From this calculation it seems likely that the space-charge factor for secondaries is quite large. It is possible that a virtual cathode of secondary electrons could be set up in the center of the gap. Electrons coming to rest there would not be reflected immediately but would leave in response to an applied rf voltage and thus cause heavy loading. Secondaries of different energies would either pass through the virtual cathode or be reflected before they reach it. In either case they would be traveling under heavy space-charge conditions and their loading would be correspondingly increased. Loading would be especially severe in the case of small transit angles and large degrees of space charge, as Fig. 1 shows. In view of these results it seems likely that a successful theory of secondary electron beam loading definitely must include the effect of space charge.

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* Received by the Institute, September 23, 1953.

[1] D. R. Hamilton, J. K. Knipp, and J. B. H. Kuper, "Klystrons and Microwave Triodes," McGraw-Hill Book Co., Inc., New York, N. Y., p. 54; 1948.

[2] F. B. Llewellyn and L. C. Peterson, "Vacuum-tube networks," Proc. I.R.E., vol. 32, pp. 144–166; March, 1944.

Contributors

J. Aarons (S'44–A'51) was born on October 3, 1921 in New York, N. Y. He received a B.S. degree in education from the City College of New York in 1942. In 1949 he received an M.A. in physics from Boston University.

J. Aarons

During the Second World War he was a member of the U. S. Air Force, he took a radar course at Harvard and M.I.T. and served as radar officer.

Mr. Aarons was employed from 1946–1953 as a physicist by the Geophysics Research Directorate, Air Force Cambridge Research Center, where he did research in solar noise and low frequency magnetic fluctuations of the earth's field.

At present Mr. Aarons is on leave from the Cambridge Research Center to do research as a Fulbright grantee at the Sorbonne and the Laboratory of the Physics of the Atmosphere in Paris.

He is a member of RESA.

❖

B. R. Bean was born September 14, 1927 in Boston, Mass. He earned a B.S. in Meteorology from the University of New Hampshire in 1949. He worked for the U. S. Weather Bureau Station at Durham, N. H. from 1945–1949.

B. R. Bean

In March, 1950 Mr. Bean transferred to the Central Radio Propagation Laboratory of the National Bureau of Standards. Since then he has been doing research on radio meteorology. He is now in charge of radio meteorological operations and analysis group.

Mr. Bean is a professional member of The American Meteorological Society.

❖

T. J. Bridges was born at Gillingham, Kent, England in December, 1923. He was educated at Maidstone Grammar School and London University where he received the B.Sc. (Eng.) degree from London University in 1944.

T. J. Bridges

He joined the Admiralty in 1945 and is at present a member of the Royal Naval Scientific Service at the Services Electronics Research Laboratory, Baldock, England.

I. Gumowski was born in Poland on March 28, 1928. He received the B.Sc. degree in electrical engineering from Laval University, Quebec, P. Q., in 1951. For a period in 1951 he was employed as an engineer in the Transmission and Development Laboratory of the Canadian Broadcasting Corporation in Montreal, but returned to Laval University to continue his studies. In 1952 he received the M.Sc. degree, and in 1953 he completed the research required for the D.Sc. degree, while working on a project sponsored by the Defense Research Board of Canada. At present he is a member of the teaching staff at Laval University. Mr. Gumowski is also a registered professional engineer in the province of Quebec.

I. Gumowski

❖

K. W. Harrison (SM'52) was born in London, England, in 1911. He studied mechanical and electrical engineering at the Borough and Regent Street Polytechnics, and in 1933 passed the graduateship examination of the Institution of Electrical Engineers. He became an associate member in 1938.

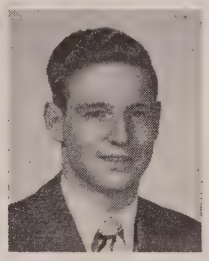
K. W. Harrison

At the end of 1930, after two and a half years with Standard Telephones and Cables, Mr. Harrison joined the Telephone Manufacturing Company. In 1940 he was sent to India to supervise installation of carrier systems and later, served as a Major in the Royal Signals.

In 1946 Major Harrison took up an appointment with the East African Posts and Telegraphs, returning to the Telephone Manufacturing Company in 1949, where he is now head of the Line-Transmission Department and responsible for the design, development and planning of transmission systems.

❖

R. H. Kingston (S'48–A'52) was born in Somerville, Massachusetts, on February 13, 1928. He received the B.S. and M.S. degrees in Electrical Engineering from the Massachusetts Institute of Technology in 1948. Under the co-operative plan in electrical engineering during the years 1945 to 1948, he worked at the Philco Corporation on component testing, radar system design, and (research on) frequency modulation and circuit theory. In June, 1951 he received the Ph.D. degree in Physics from M.I.T., where his major field was the physics of solids. He joined the staff of the Lincoln Laboratory at M.I.T., at this time, where he has been engaged in transistor research and development. From September, 1951 to September, 1952, on leave of absence from the Lincoln Laboratory, he was a member of the technical staff at Bell Telephone Laboratories in New Jersey.



R. H. Kingston

❖

J. T. Mendel (S'49–A'51) was born in Palo Alto, Calif., on March 10, 1928. He attended Stanford University from 1945–1952, receiving the B.S. degree in electrical engineering in 1949, the M.S. degree in 1950, and the Ph.D. degree in 1952. In 1951–52 he held an Atomic Energy Commission Fellowship at Stanford University. Since 1952 he has been a member of the technical staff, Bell Telephone Laboratories, Murray Hill, N. J.

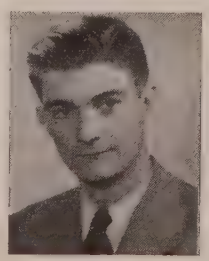
J. T. Mendel

Dr. Mendel is a member of Tau Beta Pi, Sigma Xi, and Phi Beta Kappa.

❖

R. L. Pritchard (S'45–A'51) was born in Irvington, N. J., on Sept. 8, 1924.

He attented M.I.T. on the V-12 program. He graduated from Brown University in 1946 with a B.S. degree. In 1947, he received the M.S. degree from Harvard. He received the Ph.D. degree in acoustics at Harvard in 1950.



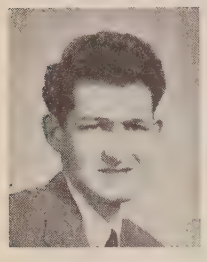
R. L. Pritchard

Since then, he has been with General Electric Research Lab. in Schenectady, working on communication and acoustic projects. For the past three years he has been doing research on the transistor from an electric-circuit point of view.

Dr. Pritchard is a member of the Acoustical Society of America and Sigma Xi.

Contributors

C. F. Quate (S'43–A'50) was born on December 7, 1923, in Baker, Nev. He received the B.S. degree from the University of Utah in 1944, and the Ph.D. in electrical engineering from Stanford University in 1949.



C. F. Quate

In 1949 he joined the Bell Telephone Laboratories, Murray Hill, N. J., as a member of the technical staff engaged in research on microwave tubes.

Dr. Quate is a member of Tau Beta Pi, and Sigma Xi.

❖

J. Shekel (A'52) was born in Bialystok, Poland on January 6, 1926. He received his engineering education at the Hebrew Institute of Technology in Haifa, Israel, and graduated with the degree Ingénieur (E.E.) in 1951.

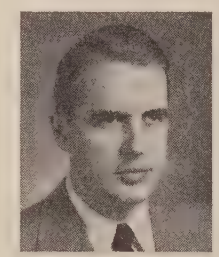
Jacob Shekel

After his graduation, Mr. Shekel was employed by the Scientific Department of the Ministry of Defense of Israel, where he does research and development work on electronic projects. Mr. Shekel works in the fields of network analysis and synthesis, and ultra-high-frequency techniques. Since February of 1953, Mr. Shekel has been a visiting lecturer in engineering at the Israel Institute of Technology, in Haifa, in addition to his regular Ministry of Defense duties.

During the summer of 1949 Mr. Shekel came to the United States and was a guest student on the Foreign Students Summer Project at the Massachusetts Institute of Technology. While he was there he took a course on microwaves.

❖

W. H. Yocom was born in Oberlin, Ohio, on May 15, 1919. He attended Oberlin College receiving the B.A. degree in Physics in 1940. He received the B.S. degree in electrical engineering from the Massachusetts Institute of Technology in 1942. In 1950 he received the M.S. degree in electrical engineering from Stevens Institute of Technology.



W H. Yocom

Since 1942 Mr. Yocom has been employed by Bell Telephone Laboratories, Murray, Hill, N. J. Mr. Yocom is concerned with the various aspects of carrier and microwave system development, including the design of testing equipment. Since 1952 he has been engaged in research in the field of microwave electronics. He is currently investigating backward-wave oscillators for millimeter waves.

Mr. Yocom holds membership in Sigma Xi.

❖

A. H. Zemanian (A'51) was born on April 16, 1925, in Bridgewater, Massachusetts. He received the B.E. degree from the College of the City of New York in 1947. He received the M.E.E. and the Eng. Sc. D. degrees in 1949 and 1953, respectively, from New York University.

A. H. Zemanian

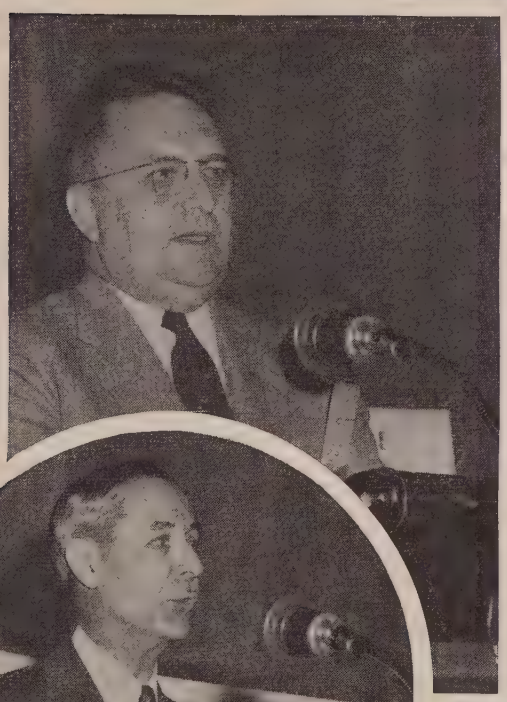
He was an instructor in the electrical engineering department of the College of the City of New York for the academic year 1947–1948. He also was an inspector of electrical circuit installation for the Board of Transportation in New York. He then joined the Maintenance Co., New York, N. Y., as an engineer, where his duties included the design and supervision of the electrical phases of contracting jobs on electric elevators. Since 1952, he has been an instructor in the electrical engineering department of New York University.

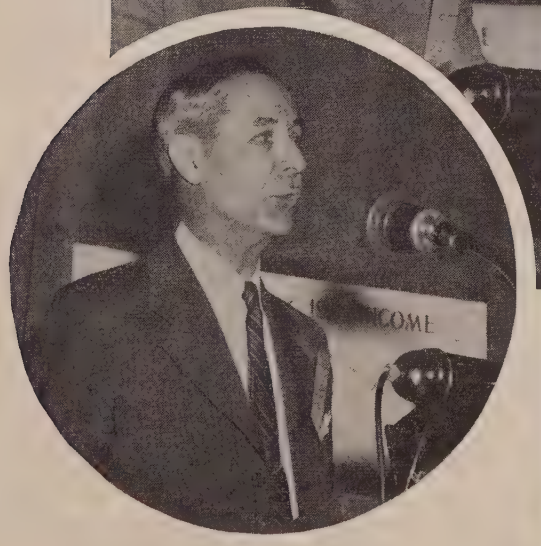
Dr. Zemanian is an associate member of the AIEE. He also holds membership in Tau Beta Pi, and Eta Kappa Nu. He holds a professional engineering licence from the state of New York.

1954 ENGINEERING SHOW
LARGEST IN IRE'S HISTORY

Panoramic view of the 604 electronics exhibits that covered the New York Kingsbridge Armory's four acres of floor space, March 22nd through March 25th.

1954 ANNUAL MEETING





(*Above, left to right*) Dr. James W. McRae, IRE 1953 President, makes the symbolic gesture of handing over his office by presenting the gavel to incoming 1954 President William R. Hewlett.

(*Left*) John D. Ryder, principal speaker, addresses the 1954 Annual Meeting on "Electronic Horseless Carriages."

(*Inset, right*) W. R. G. Baker, IRE Treasurer and (*inset, left*) Haraden Pratt, IRE Secretary, give their 1954 reports at the meeting.

COLOR TV SYMPOSIUM

(*Right*) Tremendous interest in the subject of color TV is shown by the packed hall at the Waldorf-Astoria Color TV Symposium at the recent IRE Convention. There were more than 100 standees during most of the three-hour session.

(*Below*) Speakers for this Symposium were (*left to right*) J. T. Mullin (Bing Crosby Enterprises); W. L. Hughes (Iowa State College); W. L. Brewer (Eastman Kodak); Harry F Olson (RCA); Lewis Winner (Chairman, IRE Professional Group on Broadcast Transmission Systems); E. K. Jett (WMAR-TV, Session Chairman); J. H. Haines (Allen B. DuMon Labs); and J. F. Fischer (Philco).

ove, left) IRE's President William R. Hewlett (*left*) presents 1954 Medal of Honor to William L. Everitt (*right*) during the nual Banquet at the Waldorf-Astoria, New York City, March 24

low) At the same function, Alfred N. Goldsmith (*right*) one of three IRE Founders and Editor Emeritus of the Institute, re-es the Founders Award from Past President James W. McRae.

ANNUAL CONVENTION HAS RECORD 40,108

New York City became the engineering capital of the world on March 22 to 25 when 40,108 members, exhibitors, and guests gathered from all parts of the world for the 1954 IRE National Convention. The attraction which caused this record-breaking attendance was the wealth of technical information presented in 243 technical papers and 604 exhibits at the Kingsbridge Armory, Waldorf-Astoria Hotel, and Hotel Shelton.

The convention opened on Monday morning, March 22, with the Annual Meeting of the Institute at which Prof. John D. Ryder, University of Illinois, gave the principal address entitled, "Electronic Horseless Carriages." The meeting also heard reports from IRE officers, including vice-president Maurice J. H. Ponte, who had travelled from France to attend the convention.

Technical Sessions

The 51-session program of technical papers got under way on Monday afternoon in three meeting halls at the Kingsbridge Armory, three at the Waldorf-Astoria, and one at the Shelton, and continued morning and afternoon of the following three days. The entire program was organized with the assistance of all 21 Professional Groups of the IRE, resulting in coverage of virtually every phase of activity in the radio-electronics field, as described in detail in the March issue of Proceedings.

Highpoints of the program were two special symposia held on Tuesday evening: "Engineering Based on Biological Design" organized by the Medical Electronics Group, and "High Fidelity in Audio Engineering" sponsored by the Audio Group. Each symposium featured talks by a panel of distinguished experts followed by discussions with members of the audience.

Exhibits

Perhaps of greatest general interest to those attending the convention was the new exhibit site—the Kingsbridge Armory—and the exhibits themselves. The huge, four-acre drill floor of the Armory was completely filled with 604 exhibits, providing members and visititors with the unparalleled opportunity to view the latest products of a major portion of an entire industry and to speak first-hand with company representatives. Color television picture tubes, transistors, electronic computers, project Tinkertoy assemblies, ultrasonic equipment, and medical electronic apparatus were a few of the many interesting and important developments on display. The busy visitor was greatly aided

SPEAKERS' TABLE AT THE

k Row (*left to right*)—Mr. Stuart L. Bailey, Vice-Chairman, 1954 IRE National Convention; Dr. John R. Pierce, Editor IRE; Mr. Haraden Pratt, Secretary IRE; Dr. Donald B. clair, Senior Past President IRE; Dr. William L. Everitt, 1954 Medal of Honor winner; Dr. Alfred N. Goldsmith, Founders Award winner and Guest Speaker. Front Row in Special vards Table. (*Left*) Mr. Alda V. Bedford, 1954 Vladimir K. Zworykin Television Prize; (*right*) Dr. Harold A. Zahl, 1954 Harry Diamond Memorial Award.

AND ENGINEERING SHOW ATTENDANCE REGISTER

by the grouping of exhibits according to general subjects and the labelling of exhibit aisles, such as, Television Avenue, Transistor Way, Radar Road, etc.

Visitors found the Kingsbridge Armory, located in the Bronx, readily accessible, both by bus and by subway, from the midtown area. Specially chartered busses provided free service between the Waldorf and the Armory at 10-minute intervals.

Social Events

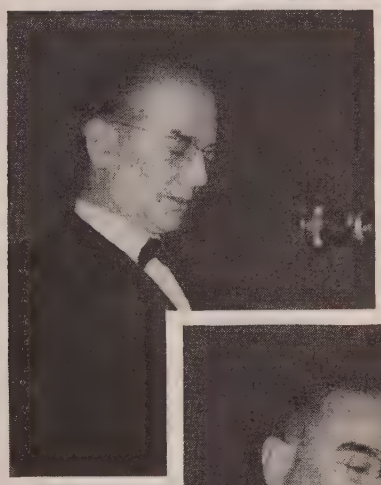
The social activities of the convention received a send-off on the first evening of the convention when a "get-together" cocktail party was held in the spacious Grand Ballroom of the Waldorf-Astoria, providing an excellent opportunity for visitors from all parts of the country to renew old acquaintances and make new ones.

A capacity crowd attended the IRE Annual Banquet, held in the Grand Ballroom on Wednesday evening, to hear Dr. Alfred N. Goldsmith, co-founder and Editor Emeritus of the IRE, speak on "IRE—Past and Future," the text of which appeared in the April issue of Proceedings, and to see him receive the IRE Founders Award from Past President James W. McRae for outstanding leadership in the radio engineering profession. William H. Doherty, Bell Telephone Laboratories, ably performed the duties of toastmaster.

The occasion was highlighted by the presentation of the Medal of Honor, IRE's highest technical award, to William L. Everitt, a past president of IRE and Dean of the School of Engineering at the University of Illinois. President William R. Hewlett made the presentation. Vice President M. J. H. Ponte responded on behalf of the 76 newly elected Fellows, whose photographs appear on the following pages.

Other annual awards presented by President Hewlett were: the Morris Liebmann Memorial Prize to Robert R. Warnecke, Campagnie Generale de Telegraphie Sans Fil, Paris, France; the Browder J. Thompson Memorial Prize Award to R. L. Petritz, U. S. Naval Ordnance Laboratory and The Catholic University of America; the Harry Diamond Memorial Award to Harold A. Zahl, Signal Corps Engineering Laboratories; the Editor's Award to P. W. Howells, General Electric Company; and the Vladimir K. Zworykin Television Prize Award to Alda V. Bedford, Radio Corporation of America.

The social program was rounded out with a tea at IRE headquarters for wives of members attending the convention.



(*Above*) Dr. William M. Doherty performs the duties of Toastmaster at the recent IRE Annual Banquet.

(*Above, center*) Dr. Maurice J. H. Ponte, IRE Vice-President, livers acceptance speech on behalf of the 1954 Fellow Award winn at the Banquet.

(*Above*) Dr. Alfred N. Goldsmith, Guest Speaker of the Banqu and winner of the 1954 Founders Award, holds a slender volume 124 pages containing all IRE publications for the year 1913. It m a startling comparison to today's Proceedings with no less th 128 text pages per monthly issue. The title of Dr. Goldsmith's t was "IRE—Past and Future."

1954 IRE ANNUAL BANQUET

Back Row (*left to right*)—Dr. William H. Doherty, Toastmaster; Mr. William R. Hewlett, President IRE; Dr. J. W. McRae, Junior Past President IRE; Dr. Maurice J. H. Ponte, V President IRE; Dr. W. R. G. Baker, Treasurer IRE; Mr. George W. Bailey, Chairman, 1954 IRE National Convention. Front Row in Special Awards Table. (*Left to right*) Mr. Paul Howells, 1954 Editor's Award; Dr. Richard L. Petritz, 1954 Browder J. Thompson Memorial Prize Award; Dr. Robert R. Warneke, 1954 Morris Liebmann Memorial Prize Awar

IRE Awards, 1954

Medal of Honor—1954

WILLIAM L. EVERITT

"For his distinguished career as an educator and scientist; for his contributions in establishing electronics and communications as a major branch of electrical engineering; for his unselfish contributions to the defense of our country; for his leadership in the affairs of The Institute of Radio Engineers."

Founders Award

ALFRED N. GOLDSMITH

"For outstanding contributions to the radio engineering profession through wise and courageous leadership in the planning and administration of technical developments which have greatly increased the impact of electronics on the public welfare."

Morris Liebmann Memorial Prize—1954

ROBERT R. WARNECKE

"For his many valuable contributions and scientific advancements in the field of electron tubes, and in particular, the magnetron class of traveling-wave tubes."

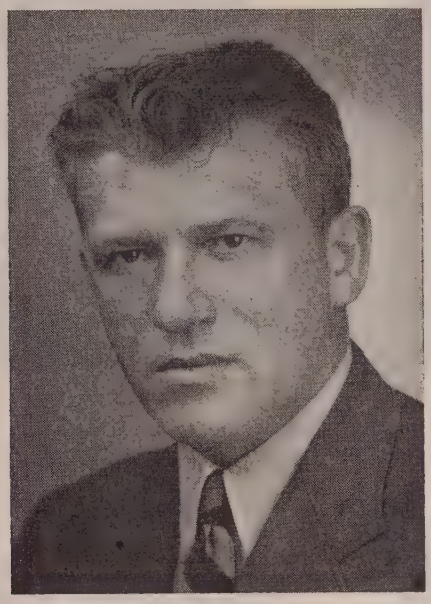
Browder J. Thompson Memorial Prize Award—1954

R. L. PETRITZ

"For his paper entitled, 'On the Theory of Noise in P-N Junctions and Related Devices,' which appeared in PROCEEDINGS OF THE I.R.E.," November, 1952.

IRE Awards, 1954

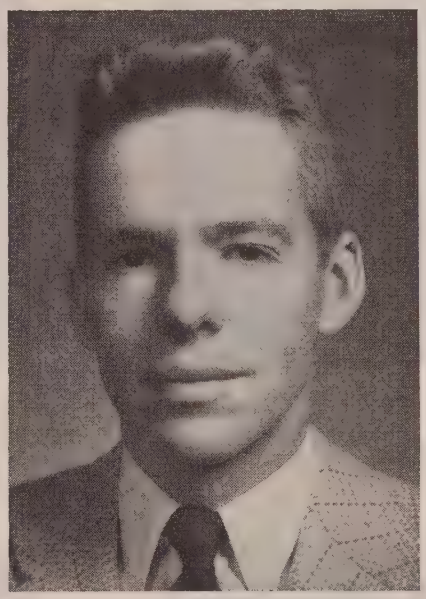
Harry Diamond Memorial Award—1954



HAROLD A. ZAHL

"For his technical contributions, his long service, and his leadership in the U. S. Army Signal Corps research program."

Editor's Award—1954



P. W. HOWELLS

"For his paper entitled, 'Transients in Color Television,' which appeared in the 1953 CONVENTION RECORD OF THE I.R.E., Part 4, Session 2."

Vladimir K. Zworykin Television Prize Award—1954

ALDA V. BEDFORD

"For his contributions to the principle of mixed highs and its application to color television."

New Fellows

E. I. ANDERSON

"For his contributions to the improvement and simplification of radio and television circuitry."

A. V. ASTIN

"In recognition of his distinguished leadership and administration of science and engineering."

W. F. BAILEY

"For his contributions to the theory, practice, and standardization of television."

D. S. BOND

"For his contributions to the development of communication and navigation apparatus and systems."

M. R. BRIGGS

"For his contributions to the design of radio transmitters, and his leadership in the establishment of industry standards."

I. F. BYRNES

"For his contributions to development and design of equipment in the field of marine communication and navigation."

MADISON CAWEIN

"For his contributions to television theory and circuitry."

J. G. CHAFFEE

"For his contributions to transmission and relay systems and frequency modulation."

BRITTON CHANCE

"For his contributions to radar development, and to the application of electronics to biophysics."

New Fellows

MARVIN CHODOROW

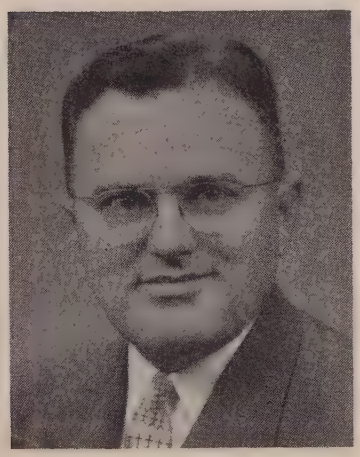
"For his contributions to the theory and design of klystron tubes."

D. H. CLEWELL

"For his research in the application of instrumentation to petroleum exploration and production."

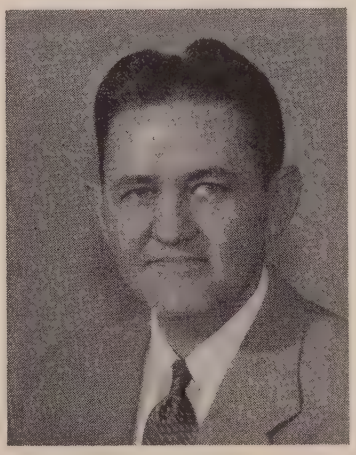
R. B. COLTON

"For his pioneering contributions to the development and application of radar in the U. S. Army."



M. S. CORRINGTON

"For mathematical analysis of frequency modulation, and of transients in networks and loudspeakers."



P. M. CRAIG

"In recognition of his leadership in the design of radio and television receivers."

C. A. CULVER

"For his contributions to the art and literature of electronic and allied fields."

F. M. DOOLITTLE

"For his pioneering work in radio communication and radio broadcasting."

W. L. DUNN

"For his contribution in the development of radio receivers."

N. E. EDLEFSEN

"For his contribution to radar, and his leadership in the development of military electronic equipment."

New Fellows

M. A. EDWARDS

"For his creative contributions to the development of the amplidyne and other control systems."

C. L. ENGLEMAN

"For his contributions to the administration of electronic programs of the U. S. Navy."

D. C. ESPLEY

"For his creative contributions to microwave and television techniques in England."

D. H. EWING

"For his contributions to the development of electronic aids to air navigation and traffic control systems."

R. M. FANO

"For his contributions in the fields of information theory and microwave filters."

E. P. FELCH

"For his contributions in the field of precision measurement and instrumentation of communication circuits."

J. L. FINCH

"For his contributions to radio communications, and particularly those associated with VLF antenna design."

R. M. FOSTER

"For basic mathematical contributions to modern network theory."

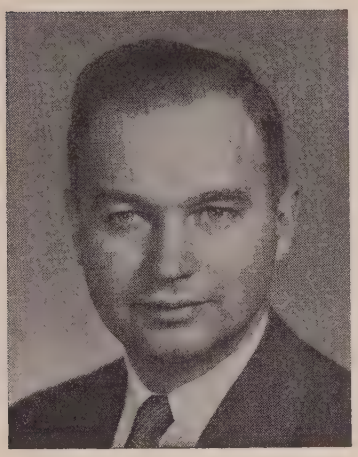
A. W. FRIEND

"For his contributions to tropospheric echo research, and to the development of magnetic materials and circuits."

New Fellows

E. G. FUBINI

"For his many contributions to the operational analysis of electronic counter measures."



I. A. GETTING

"For his contributions to the development of automatic tracking radar systems."

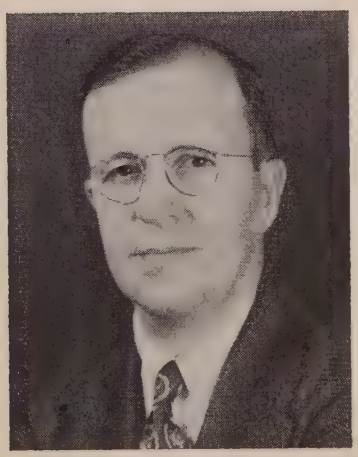
J. E. GORHAM

Posthumously awarded.

"For his leadership and contributions in the field of high-power electron tubes for military applications."

C. A. GUNTHER

"For his contributions to the design and his leadership in the development of military electronic equipment."



J. P. HAGEN

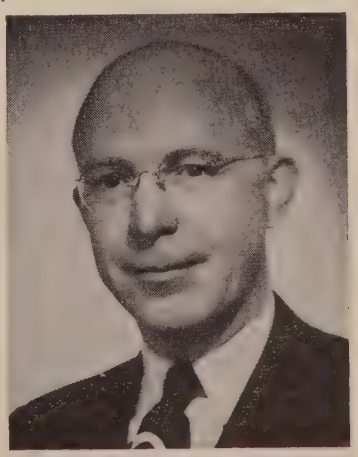
"For contributions to the development of military electronic equipment, and for measurements of solar microwave radiation."

N. L. HARVEY

"For his contributions in the fields of electronic navigation and communication."

WALTER HAUSZ

"For his contributions in the fields of radar and telemetry."



J. E. HAYES

"For contributions in Canada to the development of outstanding short wave antennas and switching designs."

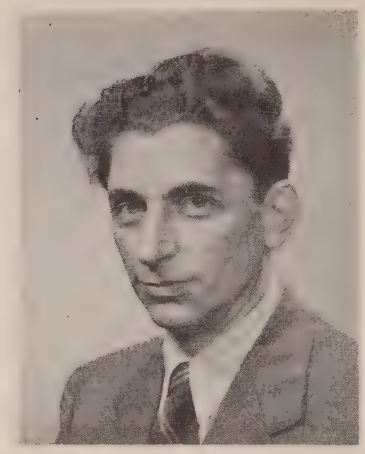
W. D. HERSHBERGER

"For his early contributions to the development of radar, and to frequency stabilization using microwave absorption lines."

New Fellows

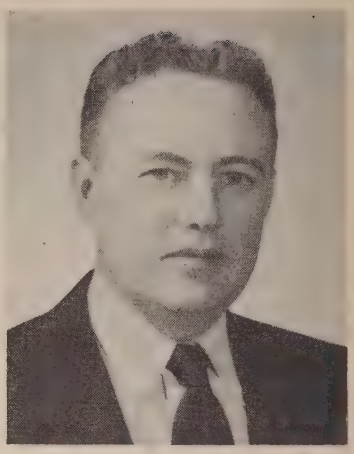
R. D. HUNTOON

"For his contributions in the fields of electronic ordnance and electron physics."



T. G. E. HUTTER

"For his contributions in the field of electron ballistics and electron optics."

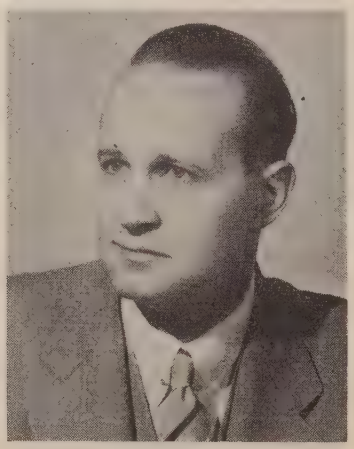


J. F. JORDAN

"For his contributions in the application of electronics to musical instruments."

W. R. KOCH

"For his creative inventions and developments in radio and television circuitry."

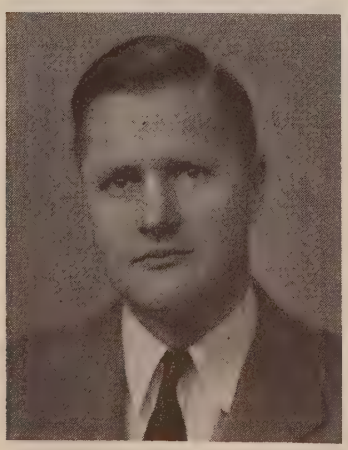


J. D. KRAUS

"For his leadership in antenna development, and his research contributions in radio astronomy."

J. B. H. KUPER

"For his contributions to nucleonic instrumentation and health physics."

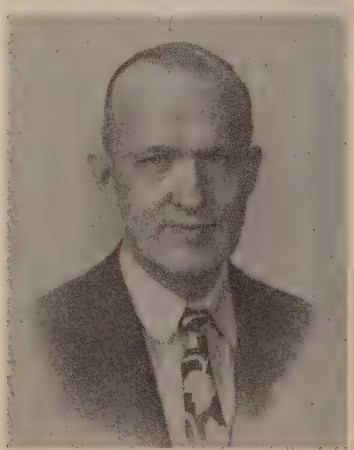


J. M. LAFFERTY

"For his research contributions to microwave tubes and high-current density cathodes."

J. J. LAMB

"For his technical contributions to amateur radio activities, and for direction of radio and electronic circuitry development."



REUBEN LEE

"For his contributions to the design and development of inductive components."

New Fellows

W. D. Lewis

"For his contributions to research, particularly in the fields of microwave filters and switching systems."

W. W. Mieher

"For his engineering contributions and technical leadership in the development of precision radar systems."

A. W Montgomery

"For his leadership in radio and telecommunication research in England, and his services in the international liaison in these fields."

Yasujiro Niwa

"For his leadership in radio engineering in Japan, and his contributions to vocational education."

B. M. Oliver

"For his contributions to communications, particularly in the field of information theory and coding systems."

J. M. Pettit

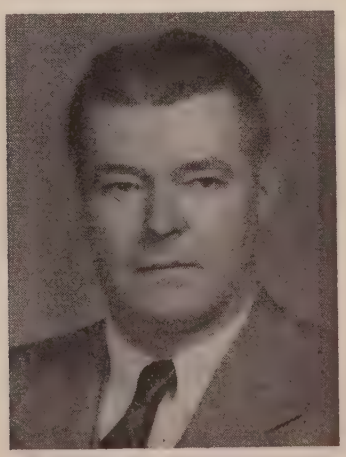
"For his outstanding work as engineer and educator in the field of high-frequency and microwave communications."

M. J. H. Ponte

"For his contributions to high-power electron tubes, and his unswerving leadership in electronic research in France."

Donald A. Quarles

"For his distinguished service in the administration of major technical programs in the fields of communications, military electronics and atomic ordnance."



W. H. Radford

"In recognition of his contributions through teaching and research in the field of radio communications."

New Fellows

EUGEN REISZ

"In recognition of his pioneering contribution to the development of grid-controlled electronic tubes."

HERRE RINIA

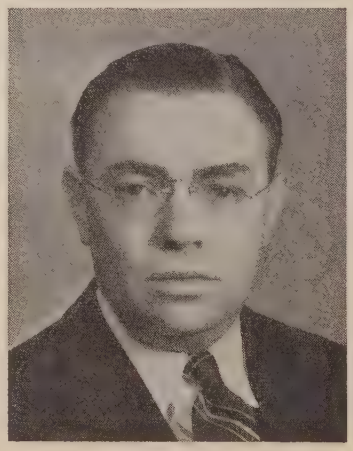
"For his creative contributions to radio engineering of Holland, and his leadership in the field of television."

T. C. RIVES

"For his leadership in military electronic research and development."

G. M. ROSE, JR.

"For his contributions to vacuum tube research, design, and manufacture, and his transistor developments."



P. C. SANDRETTO

"For his contributions to aeronautical communication and navigation."

KURT SCHLESINGER

"For many contributions to cathode-ray tubes and television circuitry."

A. H. SCHOOLEY

"For his pioneering development of fire control radar, and his contributions to electronic measurements."

H. J. SCHRADER

"For his application of information theory to navigation and television systems."

A. C. SCHROEDER

"For his contributions to television receiver circuitry, and his pioneering work in the development of color television."

New Fellows

G. R. Shaw

"For his technical contributions to the manufacture of radio tubes and their standardization."

R. F. Shea

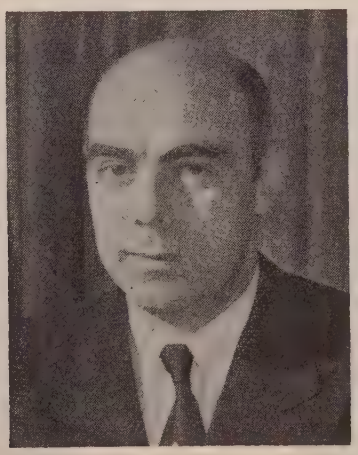
"For his contributions to FM receiver design, and his pioneering work in transistor applications."

Samuel Silver

"For research in the field of electromagnetic radiation."

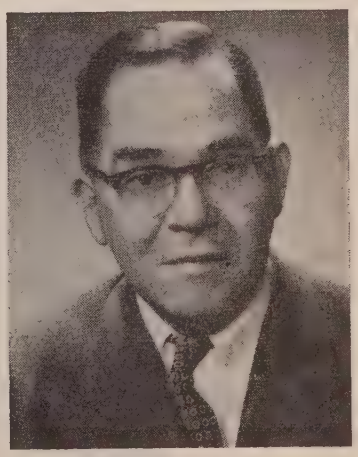
George Sinclair

"For his contributions to the development of radiating systems and model techniques in antenna measurements."



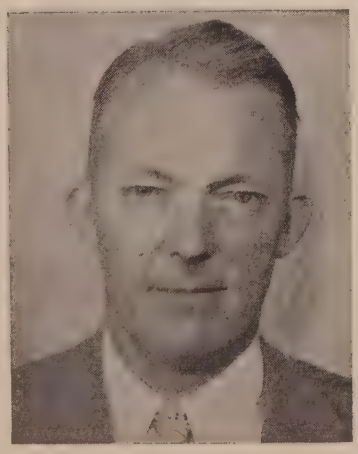
E. A. Speakman

"In recognition of his leadership in administration of electronic research and development."



B. R. Tupper

"In recognition of his application of radio techniques to the extension of toll-telephone services in Canada."



W. L. Webb

"For early contributions and engineering leadership in radio direction finding, radar, and radio communication."

P. T. Weeks

"For his contributions to electron tube research, engineering, and manufacture."

J. O. Weldon

"For his contributions to the design of high power transmitters and their use in international broadcasting."

New Fellows

K. R. WENDT

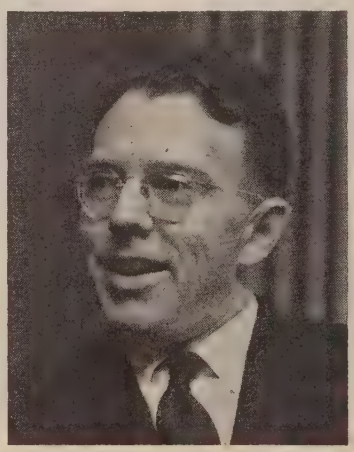
"For his contributions to the development of television equipment and circuits."

E. M. WILLIAMS

"For his contributions to the development of military electronic equipment."

J. R. WILSON

"For his stimulating leadership in research on electron tubes and devices, and in their development and manufacture."



D. E. WOOLDRIDGE

"For his contributions to physics and electronics research, and his leadership in development efforts for national defense."

1954 Convention Record of the I.R.E.

All available papers presented at the 1954 IRE National Convention will appear in the CONVENTION RECORD OF THE I.R.E., to be published in June. The CONVENTION RECORD will be issued in eleven Parts, with each Part devoted to one general subject.

Instructions on Ordering

1. If you were a member of an IRE Professional Group and had paid the Group assessment as of April 30, 1954, you will automatically receive, free of charge, that Part of the CONVENTION RECORD pertaining to the field of interest of your Group, as indicated in the chart below.

2. If you are not a member of an IRE Professional Group, CONVENTION RECORD Parts may be purchased at the prices listed in the chart below. Orders must be accompanied by remittance, and to assure prompt delivery, should be sent immediately to The Institute of Radio Engineers, 1 East 79 Street, New York 21, N. Y.

Convention Record of the I.R.E.

PART	TITLE	FREE TO PAID MEMBERS OF FOLLOWING IRE PROFESSIONAL GROUPS	PRICES FOR MEMBERS (M) LIBRARIES (L) NON-MEMBERS (NM)		
			M	L	NM
1	Antennas & Propagation SESSIONS: 30, 37, 42, 49	Antennas & Propagation	$1.25	$3.00	$3.75
2	Circuit Theory SESSIONS: 28, 35, 39, 46	Circuit Theory	1.25	3.00	3.75
3	Electronic Devices & Component Parts SESSIONS: 6, 14, 25, 32, 40, 47	Electron Devices Component Parts	1.50	3.60	4.50
4	Electronic Computers & Information Theory SESSIONS: 2, 12, 19, 27, 34	Electronic Computers Information Theory	1.50	3.60	4.50
5	Aeronautical Electronics & Telemetry SESSIONS: 3, 5, 8, 10, 15, 45	Aeronautical & Navigational Electronics Radio Telemetry & Remote Control	1.50	3.60	4.50
6	Audio & Ultrasonics SESSIONS: 11, 18, 23, 41, 48	Audio Ultrasonics Engineering	1.50	3.60	4.50
7	Broadcasting & Television SESSIONS: 13, 20, 26, 33	Broadcast Transmission Systems Broadcast & Television Receivers	1.50	3.60	4.50
8	Communications & Microwave SESSIONS: 1, 7, 21, 43, 50	Communications Systems Microwave Theory & Techniques Vehicular Communications	1.50	3.60	4.50
9	Medical & Nuclear Electronics SESSIONS: 17, 22, 24, 31	Medical Electronics Nuclear Science	1.50	3.60	4.50
10	Instrumentation & Industrial Electronics SESSIONS: 29, 36, 38, 44	Instrumentation Industrial Electronics	1.25	3.00	3.75
11	Engineering Management & Quality Control SESSIONS: 4, 9, 16	Engineering Management Quality Control	1.00	2.40	3.00
	Complete Convention Record (All Eleven Parts)		15.25	36.60	45.75

IRE News and Radio Notes

International Convention to Be Held This Spring

Fifty years ago the first paper dealing with Oxide-Coated Cathodes was published. To celebrate the anniversary, the Société Française des Ingénieurs Techniciens du Vide is organizing an International Convention to be held in Paris on the 24th and 25th of June.

The program will cover the following subjects: Basic Metal, Carbonates and Coating Process, Emission Theory of Oxide-Coated Cathodes, Definition of the Cathode Properties, Special Cathodes, and Stability of Emissive Properties.

Papers accepted by the Organizing Committee will be presented by the authors or a summary will be read by the Secretary at the sessions of the Convention. The texts will be published in a special issue of the periodical *Le Vide*, the official bulletin of the S.F.I.T.V. in their original language with, if possible, a French translation.

For further information, please write to the Société Française des Ingénieurs Techniciens du Vide, 44 rue de Rennes, Paris VI°, France.

Papers for Nuclear Science Conference Invited

The IRE Professional Group on Nuclear Science will hold its First Annual National Conference in Chicago, Ill., October 6–7 at the Sherman Hotel. The theme of the Conference will be Nuclear Reactor Technology, Particle Accelerators, and Nucleonic Instrumentation. A review of European Reactor Technology by noted scientists will be presented. The first day's session will coincide with the last day of the meeting of the 1954 National Electronics Conference.

Papers are invited on the topics mentioned above, to be about 3,000 words in length. Titles and abstracts should be submitted by mail before June 1, 1954 to: T. Brill, Argonne National Laboratory, Box 299, Lemont, Ill.

Registration, program information, housing and social details will be announced in the near future.

International Conference on Semiconductors Planned

The Netherlands Physical Society, with the support of the International Union of Pure and Applied Physics and U.N.E.S.C.O., will organize an International Conference on Semiconductors, to be held at Amsterdam, Holland from June 29–July 3.

The following scientists will deliver lectures: J. Bardeen, W. H. Brattain, H. B. G. Casimir, F. A. Kröger, D. Polder, M. Schön, W. A. Shockley, R. A. Smith, H. J. Vink, and H. Welker on subjects such as bulk recombination; surface conductivity; surface trapping; surface recombination; intermetallic compounds; the band picture in polar and nonpolar semiconductors; photoconductivity in semiconductors like PbS, PbTe, PbSe, ZnS, and CdS; and the application of general physical and chemical laws for the preparation of semiconductors with specific properties.

Discussions will be held in connection with these main lectures and there will be opportunity to discuss problems in more detail in sectional meetings. In these sectional meetings short communications of about 15 minutes each can be given.

Scientists who would like to participate in the conference or want to give a scientific contribution should communicate with the Secretary of the Society, Dr. H. J. Vink, Floralaan 142, Eindhoven, Holland, as soon as possible.

Symposium on Instrumentation Scheduled

The Institute of Industrial Health and the School of Public Health of the University of Michigan are co-sponsors of a Symposium on Instrumentation, to be held from May 24–27 at Ann Arbor, Michigan. Both manufacturers and "users" will be meeting at the University to exchange ideas and information about what is available and what is needed in the field of instrumentation for air velocity, air pollution, ionizing radiation, sound, air sampling and analysis. The Symposium is designed to be of interest to the manufacturer, safety engineer, industrial hygienist, physicist, chemist, engineer, meteorologist, and noise investigator.

Eight major speakers will give comprehensive review addresses in the eight areas in instrumentation of Industrial Hygiene. In addition, nineteen technical papers will be presented by authorities from specific fields of instrumentation. Manufacturers of instruments for industrial hygiene have been invited to exhibit and also to act as special faculty members for the Symposium.

An immediate result of the meeting will be an "encyclopedia" type of book. Papers presented at the Symposium will be published in an illustrated volume, which will also include the comprehensive review and data supplied by the manufacturers of the instruments.

Those interested in obtaining an information booklet about the Symposium should write to the following address: Director, Continued Education, 109 South Observatory Street, Ann Arbor, Michigan.

Cumulative Index Available

The Cumulative Index for the Proceedings of the I.R.E., Transactions of the I.R.E. Professional Groups, and the Convention Record of the I.R.E., for the period 1948–1953, will soon be available. The Index will be sold to members at a price of $1.00, and to nonmembers at a price of $3.00. To receive a copy of the Index, send orders with payments to The Institute of Radio Engineers, 1 East 79 Street, New York 21, N. Y.

Calendar of COMING EVENTS

IRE National Airborne Electronic Conference, Biltmore Hotel, Dayton, Ohio, May 10–12

IRE-AIEE-IAS-ISA National Telemetering Conference, Morrison Hotel, Chicago, Ill., May 24–26

Eighth NARTB Broadcast Engineering Conference, Palmer House, Chicago, Ill., May 25–27

IAS Annual Summer Meeting, IAS Building, Los Angeles, Calif., June 21–24

IRE Symposium on Global Communications, Washington, D. C., June 23–25

British Institute of Radio Engineers, 1954 Convention, Christ Church, Oxford, England, July 8–12

Sixth Annual Oak Ridge Summer Symposium, "Modern Analytical Chemistry," Oak Ridge, Tenn., August 23–27

IRE-WCEMA Western Electronic Show & Convention, Pan Pacific Auditorium, Los Angeles, Calif., August 25–27

1954 International Congress of Mathematicians, Royal Tropical Institute, Amsterdam, Holland, September 2–9

National Electronics Conference, Hotel Sherman, Chicago, Ill., October 4–6

Society of Motion Picture & TV Engineers Seventy-sixth Semi-Annual Convention, Ambassador Hotel, Los Angeles, Calif., October 17–22

IRE-RETMA Radio Fall Meeting, Hotel Syracuse, Syracuse, N. Y., October 18–20

IRE Quality Control Symposium, Statler Hotel, New York, N. Y., November 12–13

1955 Southwestern IRE Conference and Electronics Show, Baker Hotel, Dallas, Tex., February 10–12

Penn State to Offer Seminars

Pennsylvania State University, in cooperation with industry, will offer summer seminars on the following subjects: Transistors, June 9–18; Color Television, June 21–July 2; Analog Computers, June 21–July 2; Electric Contacts, June 28–July 2; and Electrostatic Precipitation, June 21–25.

The courses are primarily designed to disseminate information on the recent technological advances in these various fields. Newcomers in these fields, as well as experienced men, will benefit by the seminars.

Further information may be obtained by writing R. L. Riddle, Chairman, Short Course Committee, Pennsylvania State College, State College, Pa.

IRE News and Radio Notes

Technical Committee Notes

The **Feedback Control Systems** Committee convened on February 5th under the chairmanship of J. E. Ward. Chairman Ward introduced the following guests who have been invited to become members: Professor D. P. Lindorff, University of Connecticut; D. L. Lippitt, General Electric Co.; and T. K. Maples, Doelcam Corp. Chairman Ward discussed the use of such names as decilog (dg) and logit (lg) for this logarithmic ratio, but the Committee did not reach a definite conclusion. The Committee discussed the new phrasing of the Scope of Activities for Committee 26 as suggested by L. G. Cumming, Technical Secretary. The Committee recommended accepting the wording as is, with the deletion of the second paragraph. Chairman Ward reported to the Committee on his conversation with J. C. Lozier regarding the action taken by Subcommittee 26.1 at their last meeting. The Subcommittee after further discussion could not agree that a summing point should be used in the Basic Feedback Control Loop, since some of the members of the subcommittee still preferred the use of the more general mixing point. For the benefit of the guests of the Committee, some of the discussions regarding the choice of a summing point in the Basic Feedback Control Loop were reviewed. There was a discussion of Measurements (Methods of Testing) and Standards of Performance for Feedback Control Systems. A list of terms was compiled which are in common use for Feedback Control Systems. The list pertained to Measurements and Standards of Performance. These terms were grouped and assigned to members of the Committee to investigate existing definitions and suggested definitions thereby eliminating duplication of effort.

Under the chairmanship of P. C. Sandretto the **Navigation Aids** Committee convened on February 9th. Final consideration was given to the list of Navigation terms submitted by Mr. Gray. General Sandretto notified the members of his resignation as Chairman of the Navigation Aids Committee and announced that Professor Mimno would be the new Chairman for the coming year. He expressed his thanks to the Committee members for their co-operation and for the excellent work performed by them during his tenure as Chairman.

On February 18th the **Radio Transmitters** Committee met under the chairmanship of P. J. Herbst. The scope of Committee 15 as revised at the February 11 meeting of the Standards Committee was read to the group. The committee discussed the unsatisfactory aspects of the revised scope. It was decided by the committee to write a letter to Chairman Jensen of the Standards Committee, suggesting the formation of an ad hoc committee to study the problem of either assigning systems work to existing committees, or creating new committees to handle certain systems aspects. There was a report of the Ad Hoc Committee on Spurious Responses. Suggestions were made by the main committee. The report will be referred to the Ad Hoc Committee with the request that they continue to clarify the Report. A vote of confidence was given to the committee members for their work. There was a review of the Proposed Standard on Television Transmitters. A number of suggestions were made and this standard will be referred to the subcommittee. The proposed Standard on Pulses: Methods of Measurement of Pulse Quantities with changes incorporated according to suggestions received from the Grand Tour, was referred to the Standards Committee for review and possible approval, at its April meeting. The Committee thanked Chairman H. Goldberg and members of the Subcommittee 15.4 for their work on this Standard. Chairman Herbst read a report from Mr. Brown, Chairman of Subcommittee 15.5. It stated that their work on definitions has in general been completed and a list of tentative definitions agreed upon. Work is well under way on methods of measurement of the various characteristics of single sideband transmitters. A tentative standard should be completed within the next four months.

On February 17th the **Video Techniques** Committee convened under the chairmanship of W. J. Poch. A letter from the West Coast representative, Cameron Pierce, was read in which the names of R. DeBaun and E. Benham were suggested for the Committee roster. Approval for these additions to the membership was unanimous. E. Weppler agreed to become Chairman of a new Definitions subcommittee, thereby satisfying a long-felt need in the Video Techniques Committee. J. L. Jones reported that his subcommittee was making good progress toward the preparation of the proposed standard on the Methods of Measurement of Differential Gain and Differential Phase. It is hoped that the material will be ready for submission to the Committee within several months. J. R. Hefele reported in Mr. Baracket's absence that the Subcommittee on Video Systems and Components had not met since November but that several projects were being actively pursued by the membership. Dr. Athey reported on the general status of the work in the video recording field. A list of television signal measurement terms circulated to the Committee membership was reviewed at some length. Some corrections were made and final approval was given. It was agreed that these definitions should be given to the Definitions Co-ordinator.

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Abstracts and References

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NOTE: The Institute of Radio Engineers does not have available copies of the publications mentioned in these pages, nor does it have reprints of the articles abstracted. Correspondence regarding these articles and requests for their procurement should be addressed to the individual publications, not to the I.R.E.

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The number in heavy type at the upper left of each Abstract is its Universal Decimal Classification number and is not to be confused with the Decimal Classification used by the United States National Bureau of Standards. The number in heavy type at the top right is the serial number of the Abstract. DC numbers marked with a dagger (†) must be regarded as provisional.

A new section has been introduced covering the general technique of electromagnetic waves, oscillations and pulses (i.e. transmission lines and circuits), as distinct from specific applications to telecommunications. The new section is numbered 621.37, with subdivisions. Full details of the new classification, and of the numbers which become obsolete as a result of its introduction, are given in PE Note 535, obtainable from The International Federation for Documentation, Willem Witsenplein 6, The Hague, Netherlands, or from The British Standards Institution, 2 Park Street, London, W.1., England.

Section 621.396.67, dealing with Antennas, has been modified and expanded; details of the new classifications are given in PE Note 519.

Section 621.396.96, with subdivisions, has been introduced to cover Radar; details of the new classifications are given in PE Note 518.

New Subject Section

A section headed Automatic Computers has been introduced.

The Index to the Abstracts and References published in the PROC. I.R.E. from February 1953 through January 1954 is published by *Wireless Engineer* and included in the March 1954 issue of that journal. Copies of this issue may be purchased for $1 (including postage) from the Institute of Radio Engineers, 1 East 79th Street, New York 21, N.Y. As supplies are limited, the publishers ask us to stress the need for early application for copies. Included with the Index is a selected list of journals canned for abstracting with publishers' addresses.

ACOUSTICS AND AUDIO FREQUENCIES

534:061.3 919
International Congress on Acoustics, Delft (Netherlands), 15th–25th June 1953—P. Chavasse and R. Lehmann. (*Ann. Télécommun.*, vol. 8, pp. 335–338; Oct., 1953.) A brief report of the proceedings, with titles of the papers presented. See also 2855 of 1953.

534.1.001.362 920
Transformer Couplings for Equivalent-Network Synthesis—B. B. Bauer. (*Jour. Acous. Soc. Amer.*, vol. 25, pp. 837–840; Sept., 1953.) Equivalent circuits for mechanical arrays, based on the impedance analogy, are obtained in a simple way by introducing ideal transformer couplings.

534.121.2 921
On the Vibration of Mass-Loaded Membranes—E. T. Kornhauser and D. Mintzer. (*Jour. Acous. Soc. Amer.*, vol. 25, pp. 903–906; Sept., 1953.) The conventional treatment originated by Rayleigh is discussed, and some of the results are shown to be invalidated by an unfounded assumption.

534.13:534.414 922
Theory of the Free Vibrations of Isotropic Elastic Media—E. Ledinegg and P. Urban. (*Acta Phys. Austriaca*, vol. 7, pp. 420–435; Sept., 1953.) First-order perturbation theory is used to calculate the natural frequencies of an inhomogeneous system of arbitrary form. The method is used to determine the variation of the natural frequency of a system when its elastic constants are subjected to small spatially distributed variations. Calculations are made of the variation of resonance frequency of a cylindrical acoustic resonator on introduction of (*a*) an axial rod, and (*b*) a transverse plate. Application to the determination of Lamé's constants is discussed.

534.13:621.395.61/.62 923
The Coupling of Mechanical and Acoustic Vibration Systems—W. Güttner. (*Acustica*, vol. 3, pp. 201–206; 1953. In German.) A wideband transducer is obtained by coupling a Helmholtz resonator to the mechanical system. The optimum bandwidth conditions are attained when the reception, or the radiation, does not involve the resonator. As an example, the frequency response of a crystal microphone is determined.

534.141.4 924
On Edge Tones and Associated Phenomena—A. Powell. (*Acustica*, vol. 3, pp. 233–243; 1953. In English.) Edge-tone phenomena occurring at high jet speeds have been photographed and analyzed.

534.2 925
The Concept of Radiation Scattering and Its Application to Reinforced Cylindrical Shells—M. C. Junger. (*Jour. Acoust. Soc. Amer.*, vol. 25, pp. 899–903; Sept., 1953.) A method of analysis is developed permitting investigation of complex structures which cannot be treated by the methods previously described (3299 of 1952). The radiation pressure is expressed as the sum of two terms, one corresponding to the contribution from a rigid scatterer and the other to the contribution due to the vibration of the actual scatterer in response to the incident wave. Special cases are examined.

534.2 926
Acoustic Streaming at Low Reynolds Numbers—J. M. Andres and U. Ingard. (*Jour. Acoust. Soc. Amer.*, vol. 25, pp. 932–938; Sept., 1953.)

534.2 927
Acoustic Streaming at High Reynolds Numbers—J. M. Andres and U. Ingard. (*Jour. Acoust. Soc. Amer.*, vol. 25, pp. 928–932; Sept., 1953.)

534.2 928
Acoustic Streaming Equations: Laws of Rotational Motion for Fluid Elements—W. L. Nyborg. (*Jour. Acoust. Soc. Amer.*, vol. 25, pp. 938–944; Sept., 1953.)

534.2 929
The Propagation of Sound in a Medium with Random Fluctuations of the Refractive Index—V. Ya. Kharanen. (*Compt. Rend. Sci. U.S.S.R.*, vol. 88, pp. 253–256; Jan. 11, 1953. In Russian.) The effect of small fluctuations in a medium whose refractive index is nearly unity is considered theoretically. An expression for the deflection of the sound wave is derived. The results of the particular cases considered can be applied to the problem of the propagation of sound in the surface layer of the sea. See also 847 of 1947 (Bergmann).

534.2–13 930
The Attenuation of Sound Propagated over the Ground—J. D. Hayhurst. (*Acustica*, vol. 3, pp. 227–232; 1953.) The effect of wind on sound propagation over a concrete runway was investigated experimentally for distances up to 800 m. The only significant parameter is the component of wind along the direction of propagation, but the attenuations are statistically independent of the absolute wind speed. The change of attenuation with frequency and with height above ground is shown graphically.

534.2–14 931
Wave Propagation in a Randomly Inhomogeneous Medium: Part 1—D. Mintzer. (*Jour. Acous. Soc. Amer.*, vol. 25, pp. 922–927; Sept., 1953.) "The propagation of sound pulses from a point source in a medium where the index of refraction varies randomly is studied by means of a Born approximation to the wave equation. The coefficient of variation (standard deviation of the amplitude of a series of pulses, expressed as a percentage of the mean amplitude of the series) is evaluated for pulse lengths short compared with the time in which the refractive index varies significantly, and for ranges large compared with the wavelength of the sound. The results are in agreement with the experiments of Sheehy [1055 of 1950]."

534.21–16 932
Magnetic Damping of Acoustic Waves in Conducting Media—E. T. Kornhauser. (*Jour. Acous. Soc. Amer.*, vol. 25, pp. 1011–1012; Sept., 1953.) An elementary analysis is presented and the results are compared with those of Anderson (2635 of 1953).

534.213–14 933
The Acoustic Properties of Gas-Bubble

Screens in Water—E. Meyer and E. Skudrzyk. (*Akust. Beihefte*, no. 3, pp. 434–440; 1953.)

534.213.4–14:534.64 934
Acoustic Impedance of Rectangular Tubes—J. K. Wood and G. B. Thurston. (*Jour. Acous. Soc. Amer.*, vol. 25, pp. 858–860; Sept., 1953.) Theoretically derived values of the impedance of a section of a system of infinite parallel planes are in good agreement with measured values for long narrow tubes, when the latter are corrected for end effects.

534.213.4–14:534.64 935
End Corrections for a Concentric Circular Orifice in a Circular Tube—G. B. Thurston and J. K. Wood. (*Jour. Acous. Soc. Amer.*, vol. 25, pp. 861–863; Sept., 1953.) The values of the end correction found experimentally for the acoustic impedance of a tube with an orifice are in good agreement with values derived theoretically, treating the orifice as a plane-piston source.

534.23 936
Optimum Directivity Patterns for Linear Point Arrays—R. L. Pritchard. (*Jour. Acous. Soc. Amer.*, vol. 25, pp. 879–891; Sept., 1953.) Extended version of 3292 of 1948 (Pritchard and Rosenberg).

534.23 937
Approximate Calculation of the Directivity Factor of Linear Point Arrays—R. L. Pritchard. (*Jour. Acous. Soc. Amer.*, vol. 25, pp. 1010–1011; Sept., 1953.)

534.232:546.431.824–31 938
Barium Titanate Admittance-Temperature Characteristics—A. L. Lane. (*Jour. Acous. Soc. Amer.*, vol. 25, pp. 873–878; Sept., 1953.) Admittance variations over the temperature range 34 degrees to 67 degrees F. are studied for underwater transducers made of $BaTiO_3$ with and without addition of lead as a temperature stabilizer. Though the dielectric and piezoelectric constants of the material vary widely over this temperature range, it is possible for the transducer admittance to vary very little at certain frequencies near resonance.

534.232–14:621.395.61 939
Notes on Measurements of the Efficiency of Transducers for Water-Borne Sounds—H. Thiede. (*Akust Beihefte*, no. 3, pp. 449–451; 1953.) A method particularly applicable in the sonic and lower ultrasonic frequency ranges is described. Measurements of power absorbed for constant excitation at frequencies well below, well above and at the mechanical resonance frequency are made, the last-mentioned measurement being repeated with the transducer radiating in air. Efficiency can then be calculated.

534.321.9 940
Intensities Produced by Jet-Type Ultrasonic Vibrators—H. O. Monson and R. C. Binder. (*Jour. Acous. Soc. Amer.*, vol. 25, pp. 1007–1009; Sept., 1953.) An investigation of the influence of inlet pressure, cup position and physical proportions on the intensity and frequency of the vibrations.

534.414:534.845 941
The Impedance of a Resistance-Loaded Helmholtz Resonator—U. Ingard and R. H. Lyon. (*Jour. Acous. Soc. Amer.*, vol. 25, pp. 854–857; Sept., 1953.) Resonators with a damping screen are considered, and the dependence of the impedance on the distance between screen and aperture is investigated theoretically; the resistance and reactance components are differently affected. The results are discussed in relation to the design of perforated facings.

534.62+621.317.3.029.63/.64 942
A New Anechoic Chamber for Sound Waves and Short Electromagnetic Waves—E. Meyer, G. Kurtze, H. Severin, and K. Tamm. (*Akust. Beihefte*, no. 3, pp. 409–420; 1953.) Description of the design and construction of a chamber of dimensions 5.5m×10m×14m, with a reflection factor <10 per cent for sound waves of frequency >70 cps and for em waves of frequency >1 kmc. The treatment of the walls is basically similar to that in the pre-war anechoic chamber at the Technische Hochschule, Charlottenburg [6 of 1948 (Meyer et al.)], but behind the fibreglass wedges a cavity 12 cm deep has been provided which together with slots between the wedges, forms a resonator system. Graphite powder drawn drawn into the wedges by suction provides the attenuation of the em waves.

534.62 943
On the Measurement of Power Radiated from an Acoustic Source—W. G. Cady and C. E. Gittings. (*Jour. Acous. Soc. Amer.*, vol. 25, pp. 892–896; Sept., 1953.) An account is given of experiments which verified the theory developed by Borgnis (2536 of 1953). A slab of ρc rubber, a 90 degree wedge and a cavity radiometer were used as absorbers.

534.78 944
General Analysis of the Acoustic Structure of Speech Sounds by the Superposition of Periodic and Random Components—R. Husson and R. Saumont. (*Compt Rend. Acad. Sci., (Paris)*, vol. 237, pp. 1555–1556; Dec. 9, 1953.) A discussion of the physiological origin of the different types of component and of the ability of the ear to discriminate between them. A new classification of phonemes is given.

534.78 945
A Study of the Building Blocks in Speech—C. M. Harris. (*Jour. Acoust. Soc. Amer.*, vol. 25, pp. 962–969; Sept., 1953.) A technique is used in which speech elements are recorded on magnetic tape and joined together as required.

534.78 946
A Speech Synthesizer—C. M. Harris. (*Jour. Acoust. Soc. Amer.*, vol. 25, pp. 970–975; Sept., 1953.) The device uses a keyboard arrangement to select for reproduction speech elements magnetically recorded on a drum.

534.78 947
Some Experiments on the Recognition of Speech, with One and with Two Ears—E. C. Cherry. (*Jour. Acous. Soc. Amer.*, vol. 25, pp. 975–979; Sept., 1953.)

534.832 948
Vibration of Plates Covered with a Damping Layer—A. van Itterbeek and H. Myncke. (*Acustica*, vol. 3, pp. 207–212; 1953.) The time taken for a halving of the initial amplitude was determined experimentally for steel plates 3–26 mm thick. The variation with the mass of paint used and with temperature is shown graphically.

534.832:534.121.1 949
Sound Radiation from Plates Excited in Flexural Vibration—K. Gösele. (*Acustica*, vol. 3, pp. 243–248; 1953. In German.) The sound energy radiated per unit area of the plate is a function of the ratio f/f_{cr} where f is the frequency of the sound and f_{cr} the critical frequency of the plate. For values of this ratio >1 the radiation is independent of the total area of the plate, for unity value the radiation increases with the area and for values <1 it decreases to a limiting value. The application of the theoretical results to the construction of single- and double-shell walls is mentioned.

534.84 950
Room Acoustics—H. J. Sabine. (*Trans. I.R.E.*, vol. AU-1, pp. 4–12; July/Aug., 1953.) A brief survey of the transmission of sound through the air in a room, and of the ways in which a sound signal may be modified by the acoustic characteristics of the room.

534.84 951
Measurements of the Acoustical Properties of Two Roman Basilicas—A. C. Raes and G. G. Sacerdote. (*Jour. Acous. Soc. Amer.*, vol. 25, pp. 954–961; Sept., 1953.)

534.845–14 952
Pulse Method of Measurement of the Reflection from Absorbers of Water-Borne Sound in Tubes—W. Kuhl, H. Oberst, and E. Skudrzyk. (*Akust. Beihefte*, no. 3, pp. 421–433; 1953.) Measurements are made on a cro of the amplitude of pulses reflected from resonance-type absorbers at the end of a water-filled tube. Pulse durations between 0.5 and 1.5 ms are used, with frequencies in the range 9–22 kc. Sources of error, and their elimination, are considered in detail.

621.395.61 953
Derivation of the Laws of Force for All Magnetic and All Electric Transducers from a Single Law in Each Case—F. A. Fischer. (*Akust. Beihefte*, no. 3, pp. 441–448; 1953.) If the magnetic energy of a system of two coupled coils and the electric energy of two coupled electric conductors respectively are calculated, the mechanical force is given by the derivative of the energy with respect to x, assuming that the x co-ordinate is the only one along which the system varies. The different types of transducers result from a proper choice of this variable co-ordinate.

621.395.61 954
The Theoretical Sensitivity of Three Types of Rectangular Bimorph Transducers—E. G. Thurston. (*Jour. Acous. Soc. Amer.*, vol. 25, pp. 870–872; Sept., 1953.) Application of the method discussed previously (926 of 1953) to the cases of (*a*) the corner-loaded square bimorph, (*b*) the rectangular bimorph supported along two sides and uniformly loaded, and (*c*) the rectangular bimorph loaded only along a strip.

621.395.61 955
A Miniature Piezoelectric Microphone—J. Medill. (*Jour. Acous. Soc. Amer.*, vol. 25, pp. 864–866; Sept., 1953.) A Rochelle-salt microphone developed for use as a secondary standard in acoustic testing has a diameter of $1\frac{1}{8}$ inch. Details are given of its construction and performance.

621.395.61 956
A Miniature Microphone for Transistorized Amplifiers—B. B. Bauer. (*Jour. Acous. Soc. Amer.*, vol. 25, pp. 867–869; Sept., 1953.) The construction and performance of a moving-iron microphone suitable for use in hearing aids are described.

621.395.612.451 957
The Uniaxial Microphone—H. F. Olson, J. Preston, and J. C. Bleazey. (*Trans. I.R.E.*, vol. AU-1, pp. 12–19; July 1, Aug., 1953.) Reprint. See 2199 of 1953.

621.395.623.7:534.321.9 958
Interaction of Two Ultrasonic Waves in Air—J. Maulois. (*Toute la Radio*, vol. 20, pp. 352–354; Nov., 1953.) Report of experiments made by S. Klein using two ionophone loudspeakers to produce either (*a*) unmodulated waves of slightly different frequencies, or (*b*) an unmodulated and a modulated wave of the same carrier frequency. Audible beats were produced. The possibility of designing a loudspeaker requiring neither horn nor baffle is discussed.

621.395.623.7.011.21 959
Loudspeaker Impedance—V. Salmon. (*Trans. I.R.E.*, vol. AU-1, pp. 1–3; July/Aug., 1953.) The operation and testing of loudspeakers, and the calculation of available input power, require knowledge of five impedances, which are here discussed and defined. In estimating loudspeaker performance with different output stages or at extreme frequencies, the source regulation is an important factor. Its measurement is described.

ANTENNAS AND TRANSMISSION LINES

621.315.212 960

Modern Coaxial-Cable Technique in Great Britain—E. Baguley. (*Elec. Commun.*, vol. 30, pp. 186–216; Sept., 1953.) A discussion of the design, manufacture, and testing of long-distance telephone and television wide-band coaxial cables. Comparative performance figures and curves for cables made between 1935 and 1950 are given. The tests described include a 5-kv flashover test, pulse-echo test for detecting irregularities of impedance, HF attenuation measurement, impedance measurement between 50 kc and 10 mc and cross-talk measurements.

621.315.212:621.372.2 961

The Transmission Characteristics of Coaxial Cables—E. Adam. (*Öst. Z. Telegr. Teleph. Funk Fernsehtech.*, vol. 7, pp. 98–107 and 122–134; July/Aug. and Sept./Oct., 1953.) A simple treatment of the subject based on work published by various authors, from Maxwell up to the present time.

621.372.2+538.566 962

The Effect of the Radiation Condition in the case of Complex Wave Number, and its Significance in the Problem of Surface Waves—Haug. (*See* 1162.)

621.372.2 963

Normalized Impedance and Reflection Coefficient—P. A. Lindsay. (*Wireless Eng.*, vol. 31, pp. 43–47; Feb., 1954.) Starting from the formula giving the impedance Z at the end of a transmission line of length l in terms of the characteristic impedance Z_o, the propagation constant and l, perspective drawings are constructed of surfaces representing the relations between the normalized load impedance Z/Z_o and the voltage coefficient of reflection.

621.372.2:537.226 964

Propagation in a Dielectric Transmission Line—M. Jouguet. (*Compt. Rend. Acad. Sci., (Paris)*, vol. 237, pp. 1656–1657; Dec., 1953.) In addition to the E_o and H_o modes, mixed modes with m-order symmetry round the circumference of the line can exist. Two infinite series of modes designated $(EH)'_{m,n}$ and $(EH)''_{m,n}$ correspond to any value of m. The values of the cut-off frequency, phase velocity and attenuation are determined for various cases.

621.372.2+538.566]:621.3.012 965

New Chart for the Solution of Transmission-Line and Polarization Problems—G. A. Deschamps. (*Elect. Commun.*, vol. 30, pp. 247–254; *Trans. I.R.E.*, vol. MTT-1, pp. 5–13; March, 1953.) This chart, an orthographic projection of the Poincaré sphere (see 2418 of 1951), is considered as a modification of the Smith chart. It is normally used with a hyperbolic protractor. The examples given of its applications include the evaluation of the reflection coefficient of a stratified medium and of junction-insertion loss. Polarization problems are treated by assimilating polarization ratio to reflection coefficient.

621.372.43 966

Transmission-Line Matching System—R. E. Collin and J. Brown. (*Wireless Eng.*, vol. 31, pp. 31–35; Feb., 1954.) The two-section quarter-wave transformer is analyzed. By tolerating a certain mismatch at the center of the band it is possible to obtain a perfect match at two frequencies located symmetrically about the center frequency. This arrangement gives an increase of 20–45 per cent in bandwidth over the commonly used binomial transformer, which gives perfect matching at the center frequency only. Simple design procedure and performance curves are presented.

621.372.8 967

Velocity of Energy in Waveguides—G. Sincich. (*Alta Frequenza*, vol. 22, pp. 239–243; Oct., 1953.) A calculation is made of the instantaneous energy velocity for the TM or TE mode. The velocity is always less than that for a free wave; its value is zero, maximum or equal to the group velocity according as the longitudinal component of the electric field (for the TM mode) or the magnetic field (for the TE mode) is maximum, zero, or equal to its rms value respectively.

621.372.8 968

On the Propagation Constant in Gentle Circular Bends in Rectangular Wave Guides—Matrix Theory—A. T. de Hoop. (*Jour. Appl. Phys.*, vol. 24, pp. 1325–1327; Oct., 1953.)

621.396.67 969

A Note on the Cylindrical Antenna of Non-circular Cross Section—Y. T. Lo. (*Jour. Appl. Phys.*, vol. 24, pp. 1338–1339; Oct., 1953.) The equivalent radii are derived for antennas of polygonal and elliptical cross section.

621.396.67.012.12 970

An Automatic Recorder of Aerial Radiation Diagrams—E. G. Hamer and J. B. L. Foot. (*Jour. Brit. I.R.E.*, vol. 14, pp. 33–42; Jan., 1954.) The instrument described compares a local af reference signal with an af signal of the same frequency and of amplitude proportional to the rf signal from the antenna. Arrangements are provided for mounting large aerials at various heights, and for the presentation of results in different ways.

621.396.67.029.64:621.317.3 971

Measurements on Aerials for Centimetre Waves—Bouix. (*See* 1127.)

621.396.674.3:538.566 972

The Interaction between Electromagnetic Wave and Dipole—F. Borgnis. (*Arch. elekt. Übertragung*, vol. 7, pp. 463–466; Oct., 1953.) The expression for the mean active power abstracted from the em wave by an electric or magnetic dipole is derived interms of the electric or magnetic field, the length of the dipole (d), and the conjugate of the dipole complex current, from considerations of the near field. The expressions are valid for $d<<\lambda$ and $r<<\lambda$, where r is the distance from the dipole. The maximum effective cross-sectional area of either type of dipole with respect to a plane wave is shown to be $(3/8\lambda)\lambda^2$.

621.396.677.71 973

Vertically Polarized Microwave Antenna—R. K. Thomas and M. E. Ringenbach. (*Radio & Telev. News, Radio-Elec. Eng. Sec.*, vol. 50, pp. 13, 27; Oct., 1953.) A design giving a fan-shaped beam uses a stub-supported rigid coaxial line with longitudinally spaced circumferential slots.

621.396.677.71 974

Automatically Deiced X-Band Beacon Antenna—C. E. Thomas, Jr. (*Electronics*, vol. 27, pp. 152–155; Jan., 1954.) Details are given of a slotted-waveguide array with an omnidirectional radiation pattern and narrow vertical beam width, built to withstand wind velocities >150 knots. The radome heating circuit is controlled by a bridge which includes a capacitor exposed to the weather.

621.396.677.859:621.396.96 975

Designing Radomes for Supersonic Speeds—S. S. Oleesky. (*Electronics*, vol. 27, pp. 130–135; Jan., 1954.) Designs which will not attenuate or distort the radar beam are discussed.

AUTOMATIC COMPUTERS

681.142 976

Phantastron computes Pulse-Width Ratios—L. D. Findley. (*Electronics*, vol. 27, pp. 164–167; Jan., 1954.) Description of an analog computer developed for determining the ratio of the widths of two pulses occurring simultaneously in two channels of a radar system. The output is readily convertible into digital form. Other applications are indicated.

681.142 977

Programme Control of an Electronic Computer—H. Harmuth. (*Acta Phys. austriaca*, vol. 7, pp. 390–401; Sept., 1953.)

681.142 978

Digital Computers at Manchester University—T. Kilburn, G. C. Tootill, D. B. G. Edwards, and B. W. Pollard. (*Proc. IEE*, part II, vol. 100, pp. 487–500; Oct., 1953.) An account is given of the development of the universal high-speed computer; its features include a cr tube store for 10,240 binary digits, an intermediate store for more than 280,000 digits, completely automatic transfer between the two stores, a fast multiplier, and input and output systems using 5-hole teleprinter tape.

681.142 979

The Construction and Operation of the Manchester University Computer—B. W. Pollard and K. Lonsdale. (*Proc. IEE*, part II, vol. 100, pp. 501–512; Oct., 1953. Discussion, pp. 540–543.) The account particularly stresses techniques used to achieve reliability and ease of maintenance, and includes details of the performance over a period of about 61 weeks. *See* also 978 above.

681.142 980

Universal High-Speed Digital Computers: a Decimal Storage System—T. Kilburn and G. Ord. (*Proc. IEE*, part II, vol. 100, pp. 513–522; Oct., 1953. Discussion, pp. 540–543.) Three decimal-digit storage systems using cr tubes are described. The best compromise as regards speed of operation, storage capacity and reliability is obtained with a cr-tube screen comprising ten separate elements. In a serially controlled store holding 32 words of 8 decimal digits the digit period was 40 μs and could be reduced to 30 μs. These results compare favorably with those for the binary store used in the Manchester University computer.

681.142:621–526 981

The Design and Testing of an Electronic Simulator for a Hydraulic Remote-Position-Control Servo Mechanism—F. J. U. Ritson and P. H. Hammond. (*Proc. IEE*, part II, vol. 100, pp. 568–569; Oct., 1953.) Discussion on 1054 of 1953.

681.142:621–526 982

Analogue Computers for Feedback Control Systems—R. A. Bruns. (*Trans. AIEE*, vol. 71, part II, pp. 250–254; 9152.)

681.142:621.314.222 983

Transformer-Analogue Network Analysers—M. W. H. Davies and G. R. Slemon. (*Proc. IEE*, part II, vol. 100, pp. 469–480; Oct., 1953. Discussion pp. 481–486.) The use of voltage transformers as multipliers for complex quantities is explained. A network analyzer is described comprising a measuring section and a number of identical panels each with two multiplying units. These panels can be used to represent impedances, admittances, etc. Negative resistance and unilateral mutual inductance can easily be represented. Applications discussed include the solution of linear simultaneous equations.

681.142:621.385 984

Valve Reliability in Digital Calculating Machines—L. Knight. (*Elec. Eng.*, vol. 26, pp. 9–13; Jan., 1954.) A discussion of measures which can be taken to reduce the number of tube failures and to reduce the inconvenience caused when failures do occur.

681.142:621.385.832 985

Recent Advances in Cathode-Ray-Tube Storage—F. C. Williams, T. Kilburn, C. N. W. Litting, D. B. G. Edwards, and G. R. Hoffman. (*Proc. IEE*, part II, vol. 100, pp. 523–539; Oct., 1953. Discussion, pp. 540–543.) Modifications have been made to the system described by Williams and Kilburn (2258 of 1949) and used in the Manchester University computer; the defocus-focus system and the signal ampli-

fier are discussed particularly. An account is given of experimental work undertaken to clarify the storage mechanism. The theoretical basis of the system is examined; the original theory, though inaccurate in detail, is satisfactory for practical purposes and can be used to predict storage results. The tube operation is analogous to that of a triode tube, the collector, bombarded spot, and surrounding screen surface corresponding respectively to the anode, cathode, and control grid.

CIRCUITS AND CIRCUIT ELEMENTS

621.318.572:621.385.832 986

A Multi-Decade Predetermined Counter—A. Dorn. (*Elec. Appl. Bull.*, vol. 14, pp. 91–101; June/July, 1953.) The automatic four-decade counter described is a modified version of that noted in 603 of February. The maximum counting rate is 12,500 per sec. and the minimum duration of a complete cycle of counts is $\frac{1}{3}$ ms.

621.372 987

Directional Coupling with Transmission Lines—W. L. Firestone. (*Tele-Tech.*, vol. 12, pp. 95–97, 178; Nov., 1953.) Adaptation of devices developed in connection with waveguides to forms with open-wire lines, for use at lower radio frequencies.

621.372 988

Response of an Amplifier Stage with Anti-resonant Circuit to an Input Voltage whose Instantaneous Frequency varies linearly as a Function of Time—P. Poincelot. (*Compt. Rend. Acad. Sci.*, (*Paris*) vol. 237, pp. 1314–1315; Nov. 23, 1953.) Analysis is given for an input which can be represented by a particular Fourier integral.

621.372.2:534.321.9 989

Equivalent-Network Representations for Solid and Mercury Delay Lines—S. Dairiki, T. E. Lawrence, and R. A. Mapleton. (*Jour. Acous. Soc. Amer.*, vol. 25, pp. 841–853; Sept., 1953.) An equivalent-network representation is obtained from consideration of the piezoelectric properties of the quartz crystal used to convert the electrical energy into ultrasonic energy. A formula is derived for the transfer function of the line. When mercury and fused-quartz lines are compared on the basis of the gain-bandwidth product only, the fused-quartz line is superior; but other factors (e.g. production of sparious signals) must be taken into account when selecting a line for a specific application.

621.372.413 990

Characteristics of an Elliptical Electromagnetic Resonant Cavity operating in the TE_{111} Mode—T. P. Higgins and A. W. Straiton. (*Jour. Appl. Phys.*, vol. 24, pp. 1297–1299; Oct., 1953.) "Equations and graphs are presented of the resonant wavelength and quality factor of elliptical resonant cavities operating in the TE_{111} mode. Both components of this this mode are treated and distinctive characteristics of each are shown as a function of eccentricity of the cavity."

621.372.413 991

Interaction between Classical Electron and Quantized Electromagnetic Field—I. R. Senitzky. (*Phys. Rev.*, vol. 91, pp. 1309–1311; Sept. 13, 1953.) The equations of motion are solved by a method involving the corresponding difference equations, and an expression is derived for the probability of transition between a high-energy state and neighboring states of the field inside a resonant cavity.

621.372.5 992

The Transformation Properties of Loss-Free Quadripoles between Homogeneous Lines, and a Proof of Weissfloch's Transformation Law using Circle Geometry—H. Lueg. (*Arch. elekt. Übertragung*, vol. 7, pp. 478–484; Oct., 1953.) A method for evaluating certain parameters in the application of Weissfloch's transformation law (3287 of 1943 and back references). This law has applications in measuring technique and circuit theory in the *cm*-λ and *dm*-λ regions.

621.372.5 993

On the Synthesis of Reactance 4-Poles—H. J. Carlin and R. La Rosa. (*Jour. Appl. Phys.*, vol. 24, pp. 1336–1337; Oct., 1953.) Using theory developed by Bleevitch (1547 of 1952), formulas are derived for the elements of the impedance or admittance matrix of a loss-free quadripole which, when terminated by a resistor, realizes a prescribed physical driving-point impedance.

621.372.5.012 994

A Block-Diagram Approach to Network Analysis—T. M. Stout. (*Trans. AIEE*, vol. 71, part II, pp. 255–260; 1952.) An extension of the method described by Graybeal (1534 of 1952) to electrical networks, with examples of application in the analysis of ladder, bridged-T and parallel-T networks.

621.372.5.029.64:621.318:1 995

Ferrites in Microwave Applications—J. H. Rowen. (*Bell. Sys. Tech. Jour.*, vol. 32, pp. 1333–1369; Nov., 1953.) A discussion of the microwave Faraday effect and of the mechanism of power absorption in ferrites. Plane-wave theory is given, and is extended to cover waveguides. The measurement and applications of the effects are surveyed. See also 1233 of 1952 (Hogan).

621.372.512.029.63:621.372.21 996

Coupled Lines as High-Frequency Transformer—O. Bronder. (*Fernmeldetech. Z.*, vol. 6, pp. 475–480; Oct., 1953.) Systems comprising two parallel two-conductor lines are considered Voltage and current distribution and input impedance are calculated in terms of terminating impedance, overlap length and coupling. An indication is given of the requirements to be satisfied for wide-band operation, and some particular arrangements are discussed.

621.372.54 997

Tentative Graphical Representation of the Image Transfer Coefficient of Ladder-Type Filters with or without Losses—J. E. Colin. (*Câbles & Trans.*, vol. 7, pp. 242–262; July, 1953.) The method is based on the expression of the image transfer coefficient $p=a+jb$ as a function of frequency in accordance with the classification of filters given previously (2748 of 1949 and 575 of 1950).

621.372.543.2.029.5 998

A Set of 0.4-Octave Band-Pass Filters for Frequencies between 10 and 200 kc/s—H. Pursey. (*Electronic Eng.*, vol. 26, pp. 31–33; Jan., 1954.) Filters designed for eliminating harmonics are described. In order to achieve the desired sharp cut-off on the high-frequency side the filters are operated into an abnormally low impedance, about a tenth of the center-frequency impedance.

621.372.56:621.372.8 999

Broadband Rotary Waveguide Attenuator—B. P. Hand. (*Electronics*, vol. 27, pp. 184–185; Jan., 1954.) The attenuator comprises three collinear sections of round waveguide, each with a resistive film arranged in a diametral plane; the middle section is rotated to obtain the desired attenuation. Some performance figures are given for an experimental model; the characteristics are independent of frequency.

621.372.6 1000

A General Network Theorem, with Applications—B. D. H. Tellegen. (Proc. I.R.E., (*Australia*) vol. 14, pp. 265–270; Nov., 1953.) Slightly modified reprint of paper noted in 62 of 1953.

621.373.42 1001

"Chameleon" Oscillator—T. Roddam. (*Wireless World.*, vol. 60, pp. 52–55; Feb., 1954.) An oscillator giving high frequency stability and using relatively few components is based on a cathode-coupled Hartley circuit with resistance in the cathode lead.

621.373.421.13 1002

The Measure of Activity of Oscillators and the Performance Index—W. Herzog. (*Arch. elekt. Übertragung*, vol. 7, pp. 470–472; Oct., 1953.) The measure of activity of an oscillator is defined as the ratio of grid-voltage amplitude to anode-current amplitude. In crystal-controlled oscillators it is proportional to the crystal performance index. The measure of activity is evaluated for several common types of oscillator.

621.373.424.029.4 1003

Design of Heterodyne Oscillators—H. Lenmartz. (*Funk u. Ton*, vol. 7, pp. 526–534; Oct., 1953.) The choice of frequencies for producing a given beat frequency, and the design of the variable tuning capacitor are considered. Precautions required to ensure frequency stability are discussed.

621.373.43 1004

The Operation of Generators for Steep-Fronted Waves—J. Lagasse, J. Favarel, and P. Sido. (*Compt. Rend. Acad. Sci.*, (*Paris*), vol. 237, pp. 1151–1152; Nov. 9, 1953.) When the self-inductance of the generator is taken into account in the analysis, a third-order differential equation is obtained; simplification is introduced by taking advantage of a particular solution. Graphical methods of investigation, based on the work of Bergeron, are also outlined.

621.373.52 1005

Transistor Oscillator for Use in Multifrequency Pulsing Current Supply—F. E. Blount. (*Bell Sys. Tech. Jour.*, vol. 32, pp. 1313–1331; Nov., 1953.) Discussion of an oscillator for use in the transmission of digital information, at frequencies in the band 700–1700 cps, over telephone trunk-lines. Performance details are given.

621.374:621.318.572 1006

Introduction to the Use of Electronic Tubes in Pulse Techniques—P. A. Neeteson. (*Elec. Appl. Bull.*, vol. 14, pp. 120–112; June/July, 1953.) Circuit phenomena produced by switching are briefly reviewed and the operation of tubes as nonlinear or switching elements is discussed as an introduction to the consideration of some well known pulse circuits.

621.374.33 1007

A Special Time-Discriminating Selector for Electronic Pulses—A. Alberigi, F. Lepri, and G. Stoppini. (*Nuovo Cim.*, vol. 9, pp. 365–368; April 1, 1952.) A circuit is described which is able to select the last of a series comprising an arbitrary number of pulses distributed in an arbitrary way within a fixed time interval.

621.375.1:621.396.822 1008

Determination of the Noise Factor of Various Amplifiers by a Comparison Method—W. Mansfeld. (*Funk u. Ton*, vol. 7, pp. 501–507; Oct., 1953.) The relative noise factor can be determined by using a standard-type signal generator connected to the input of the amplifier and and a band-pass filter coupled to a tube voltmeter at the output.

621.375.232 1009

Wide-Band I.F. Amplifiers—H. S. Jewitt. (*Wireless World*, vol. 60, pp. 86–90; Feb., 1954.) Bandwidths >10 mc, such as are required for dealing satisfactorily with very short pulses, are achieved without necessitating alignment of the amplifier with a particular set of tubes, by using a negative-feedback design technique. The amplifier is treated as a series of pairs of tubes, the first of each pair operating without and the second with feedback. A table gives component values in a form requiring a minimum of computation. A particular 12-stage amplifier designed by this method for a centre frequency of 60 mc had a bandwidth of 20 mc to the -1db points and an over-all gain of about

90 db; the response varied only very slightly when all the tubes were changed.

621.375.3 1010
Magnetic Amplifier uses Conventional Inductors—A. I. Bennett, Jr. (*Electronics*, vol. 27, pp. 181–183; Jan., 1954.)

621.375.3 1011
Negative Inductance cuts Magnetic-Amplifier Lag—G. M. Ettinger. (*Electronics*, vol. 27, pp. 162–163; Jan., 1954.) The speed of response of a magnetic amplifier is improved, without reducing its sensitivity, by balancing out the effect of control-circuit inductance by means of a "negative-inductance" device comprising a valve with an iron-cored mutual inductor.

621.375.3.024 1012
The Design of a Practical D. C. Amplifier Based on the Second-Harmonic Type of Magnetic Modulator—S. W. Noble and P. J. Baxandall. (*Proc. IEE*, part II, vol. 100, pp. 567–570; Oct., 1953.) Discussion on 72 of 1953.

621.375.4 1013
Transistors: Theory and Application: Part II—Cascading Transistor Amplifier Stages—A. Coblenz and H. L. Owens. (*Electronics*, vol. 27, pp. 158–161; Jan., 1954.) Formulas and values of circuit parameters are tabulated for various possible arrangements using point-contact or junction transistors.

621.375.4.029.3:621.314 1014
Power Transistors for Audio Output Circuits—Giacoletto. (*See* 1247.)

621.396.822:621.372 1015
Physical Basis of Thermal Noise—D. A. Bell. (*Wireless Eng.*, vol. 31, pp. 48–50; Feb., 1954.) An examination is made of methods developed by various workers for analyzing circuit noise due to random motion of charge carriers. It is concluded that it is not necessary to postulate thermodynamic equilibrium in order to be able to calculate the thermal noise.

621.375.4+621.373.52 1016
Principles of Transistor Circuits [Book Review]—R. F. Shea (Ed). Publishers, Chapman & Hall, London, 535 pp., 88s. (*Wireless Eng.*, vol. 31, pp. 51–52; Feb., 1954.) A comprehensive work on transistors and their circuitry, written by nine members of the staff of the electronics laboratory of the G.E.C. Company of America.

621.376+621.314.26 1017
Modulators and Frequency Changers [Book Review]—D. G. Tucker. Publishers: Macdonald & Co., London, 232 pp., 28s. (*Brit. Jour. Appl. Phys.*, vol. 4, p. 319; Oct., 1953.) For research, development, and maintenance engineers. High-power transmitting modulators are not treated.

GENERAL PHYSICS

534.014.5 1018
Graphical Presentation of Oscillator Resonance, Applicable to the Study of Nonlinear Systems—P. Liénard. (*Acustica*, vol. 3, pp. 212–223; In French.)

535.31:535.13 1019
Derivation of the Laws of Geometrical Optics from Maxwell's Field Theory—G. Mandl. (*Acta Phys. austriaca*, vol. 7, pp. 365–389; Sept., 1953.)

535.42:538.56:621.372.8 1020
The Diffraction of Electromagnetic Waves by a Semi-infinite Circular Waveguide—J. D. Pearson. (*Proc. Camb. Phil. Soc.*, vol. 49, part 4, pp. 659–667; Oct., 1953.) Formulas are obtained for the currents on the waveguide, and asymptotic expressions are derived for the longitudinal components of the electric and magnetic vectors in the waveguide at large distances from the mouth.

537.122:538.56 1021
A Collective Description of Electron Interactions: Part 2—Collective vs Individual Particle Aspects of the Interactions—D. Pines and D. Bohm. (*Phys. Rev.*, vol. 85, pp. 338–353; Jan. 15, 1952.) Part 1: 2975 of 1951.

537.221:546.47 1022
The Electromotive Force developed by a Creeping Zinc Crystal—F. D. Coffin and S. L. Simon (*Jour. Appl. Phys.*, vol. 24, pp. 1333–1334; Oct., 1953.)

537.311.1 1023
Theory of Plasma Waves in Metals—P. A. Wolff. (*Phys. Rev.*, vol. 92, pp. 18–23; Oct. 1., 1953.) The Hartree approximation is used to investigate the effect of the crystal lattice on plasma oscillations in metals. A formula is derived for the oscillation frequency; for free electrons this reduces to that given by Bohm and Gross (88 and 89 of 1950); for insulators there are no oscillations. In metals with occupied d bands there is strong coupling between the plasma oscillations and the d electrons, causing frequency broadening.

537.311.33 1024
Theory of A.C. Space-Charge Polarization Effects in Photoconductors, Semiconductors, and Electrolytes—J. R. MacDonald. (*Phys. Rev.*, vol. 92, pp. 4–17; Oct. 1, 1953.) Linear theory is developed for solid or liquid materials containing charge carriers which can move freely within the material but are blocked by the electrodes. The carriers may be electrons, holes, ions or ion vacancies. The general solution for the admittance of the material is obtained for an arbitrary ratio between the mobilities of positive and negative carriers, and is discussed for some special cases. Results are compared with those derived from other theories.

537.525.5 1025
The Field-Emission-Initiated Vacuum Arc: Part 1—Experiments on Arc Initiation—W. P. Dyke, J. K. Trolan, E. E. Martin, and J. P. Barbour. (*Phys. Rev.*, vol. 91, pp. 1043–1054; Sept. 1, 1953.) The experiments were carried out under conditions of high vacuum and clean cathode surface. Results show (*a*) that breakdown is immediately preceded by recognizable emission anomalies, (*b*) that breakdown occurs at a critical current density of the order of 10^8 A/cm^2, for applied microsecond pulses in the range 8.8–60.1 kv, and (*c*) that breakdown is independent of cathode bombardment.

537.525.5 1026
The Field-Emission-Initiated Vacuum Arc: Part 2—The Resistively Heated Emitter—W. W. Dolan, W. P. Dyke, and J. K. Trolan. (*Phys. Rev.*, vol. 91, pp. 1054–1057; Sept. 1, 1953.) Theory is developed for the heat flow in an emitter of geometry approximating to that used in the experiments reported in 1025 above. The calculations support the view that breakdown is initiated by a resistive heating process.

537.533:537.525.92 1027
Space-Charge Effects in Field Emission—J. P. Barbour, W. W. Dolan, J. K. Trolan, E. E. Martin, and W. P. Dyke. (*Phys. Rev.*, vol. 92, pp. 45–51, Oct. 1, 1953.) "A progressive reduction of the observed field current below values expected from the empirical law for increasing values of the potential is attributed to space charge. The current density expected in the presence of space charge is calculated from the Fowler-Nordheim field emission theory using values of the cathode electric field obtained from a solution of Poisson's equation for plane electrodes with boundary conditions appropriate to field emission. The result is a generalization of Child's equation, and is asymptotic to it when the applied potential is large compared with the value required for appreciable field emission."

537.533:538.691 1028
"Brillouin Flow" with Thermal Velocities—J. R. Pierce and L. R. Walker. (*Jour. Appl. Phys.*, vol. 24, pp. 1328–1330; Oct., 1953.) "A type of electron flow in a constant magnetic field is described. The beam of electrons is supposed to be everywhere in thermal equilibrium and the usual Brillouin flow is found when the equilibrium temperature tends to zero. Some considerations are put forward bearing on the choice of a suitable beam temperature in specific problems."

537.533.8 1029
A Modified Theory of Production of Secondary Electrons in Solids—A. van der Ziel. (*Phys. Rev.*, vol. 92, pp. 35–39; Oct. 1, 1953.) Difficulties in previous theories of energy loss and secondary-electron production in metals [see e.g. 3069 of 1952 (Dekker and van der Ziel)] are removed by replacing the Coulomb interaction between a primary electron and a lattice electron by a screened Coulomb interaction.

537.56:538.63 1030
Magnetic-ionic Theory of Weakly Ionized Gases in the Presence of an Oscillating Electric Field and a Constant Magnetic Field—R. Jancel and T. Kahan. (*Jour. Phys. Radium*, Oct., 1953.) The velocity distribution functions for a nonuniformly ionized gas are calculated by solving the Boltzmann differential equations. Explicit expressions are derived for the magneto-ionic conductivity, the dielectric tensor, the Hall effect, the deflection of the electron beam and the generalization of Langevin's mobility formula. Comparison is made with other methods of calculation.

537.562:538.63 1031
The Conductivity Tensor of Electros Plasmas in the presence of a Constant Magnetic Field—M. Bayet, J. L. Delcroix, and J. F. Denisse. (*Compt. Rend. Acad. Sci.*, (*Paris*), vol. 237, pp. 1503–1505; Dec. 9, 1953.) Assuming the mean collision frequency to be independent of velocity, it is shown that the form of the conductivity tensor given by Jancel and Kahan (1030 above) differs from that given by Huxley (1266 of 1952). In order to represent the magnetoresistance terms correctly by Janceal and Kahan's method, closer approximation is necessary.

538.22 1032
Some Magnetic Properties of Metals: Part 6—Surface Corrections of the Landau Diamagnetism and the de Haas-van Alphen Effect—R. B. Dingle. (*Proc. Roy. Soc. A.*, vol. 219, pp. 463–477; Oct. 7, 1953.) An expression is given for the magnetic moment of a large but finite system of electrons in terms of the volume and two mutually perpendicular surface areas together with constants independent of the shape of the system. The theory is verified by calculating the values of the constants for two different shapes. Corrections to part 5 (1967 of 1953) are noted.

538.221 1033
A Two-Electron Example of Ferromagnetism—J. C. Slater, H. Statz, and G. F. Koster. (*Phys. Rev.*, vol. 91, pp. 1323–1341; Sept. 15, 1953.) The method of dealing with problems of ferromagnetism on the basis of energy-band theory [*Rev. Mod. Phys.*, vol. 25, p. 199; 1953. (Slater)] is applied to a simple case.

538.51 1034
The Law of Induction—G. Vallauri. (*Alta Frequenza*, vol. 22, pp. 211–238; Oct., 1953.) A survey of the literature indicates a lack of clarity on this subject. The law of induction is formulated generally in terms of a geometrical entity called the "flux line," which exists in the field of a solenoidal vector.

538.52 1035
Induction Phenomena consequent on the Movement of Material in Primary Magnetic Fields, and their Experimental Applications: Part 3—Fundamental Theory for Very General

Cases—H. Hinteregger. (*Acta Phys. austriaca*, vol. 7, pp. 337–354; Sept., 1953.) Theory applicable to any type of motion is developed. Part 2: 388 of March.

538.561:537.533 1036

Motion of an Electron in a Magnetic Undulator—R. Combe and M. Feix. (*Compt. Rend. Acad. Sci.*, (*Paris*), vol. 237, pp. 1318–1320; Nov. 23, 1953.) Rigorous solutions are obtained for the equations of motion of an electron in a millimeter-wave generator of the type described by Motz et al. (3582 of 1953).

538.561:537.533 1037

Frequencies and Power of Waves Radiated by a Magnetic Undulator—R. Combe and M. Feix. (*Compt. Rend. Acad. Sci.*, (*Paris*) vol. 237, pp. 1660–1662; Dec. 21, 1953.) Continuation of investigation noted in 1036 above. The dimensions of an undulator for generating mm waves are determined from rigorous calculations of electron trajectories. The results differ greatly from values derived from classical theory. The values found for radiated power are of the same order.

538.566:537.56 1038

The Equations of Propagation of Electromagnetic Waves in an Ionized Gas—A. M. Confetta. (*R. C. Ist. Lombardo*, vol. 85, pp. 495–502; 1952.) Two different methods are used to derive the equations for the propagation of em waves in an ionized gas subjected to a time-varying em field and to a constant magnetic field; terms depending on past states are found to occur.

538.566.029.64:535.51-7 1039

Effect of a Metal Plate on Total Reflection W. Culshaw and D. S. Jones. (*Proc. Phys. Soc.*, vol. 66, pp. 859–864; Oct., 1, 1953.) The phase changes which occur on total reflection are modified when a metal plate is placed near and parallel to the reflecting boundary. By moving the plate, the phase difference between plane waves of equal amplitudes polarized in and perpendicular to the plane of incidence can be varied from $-\pi$ up to the positive value obtained without the plate. This has been confirmed experimentally using a wavelength of 1.25 cm and a 45 degrees perspex prism. When the refractive index is $\geqq 1+\sqrt{2}$, there are two positions of the plate for which the reflected wave is circularly polarized, the electric vector rotating in opposite directions for these two positions.

538.566.2 1040

Theory of a [linear] Conductor near the Boundary of Two Media—A. A. Pistol'kors. (*Compt. Rend. Acad. Sci. U.S.S.R.*, vol. 86, pp. 941–943; Oct. 11, 1952. In Russian.) The electromagnetic field due to an oscillation of angular frequency ω in the conductor can usually be split into two components: a cylindrical-wave field and a spherical-wave field. Near the axis of the conductor, which is parallel to the boundary plane, the field is of a special type. The condition for the existence of the two components is derived and applied to particular cases.

538.691 1041

Deflection of High-Energy Electrons in Magnetized Iron—S. Berko and F. L. Hereford. (*Phys. Rev.*, vol. 91, pp. 1127–1130; Sept. 1., 1953.) The effective magnetic field acting on electrons traversing a ferromagnetic medium is shown, by experiments, to be equal to the flux density.

GEOPHYSICAL AND EXTRATERRESTRIAL PHENOMENA

523.5:621.396.96 1042

Radio Echo Studies of Meteor Ionization—T. R. Kaiser. (*Advances Phys.*, vol. 2, pp. 495–544; Oct., 1953.) The physical properties of the meteors and the columns of ionization which they produce, and the properties of the atmosphere in the meteor ionization region were investigated using, mainly, pulse equipment operating in the 4-10-m waveband. The development of a satisfactory theory of scattering of radio waves from meteor trails, together with the theory of evaporation from meteors, makes possible the detailed interpretation of the radio echoes obtained. In particular, the estimated incident flux of meteors is found to be sufficient to maintain a continuous ionospheric layer extending from ~130 km down to 115 km, below which ionization will be distributed in patches becoming fewer (and larger) with decreasing height. Over 40 references.

523.72:621.396.822 1043

Observations of Solar Radio-Noise Storms—B. Vauquois, P. Coupiac, and M. Laffineur. (*Compt. Rend. Acad., Sci.*, (*Paris*) vol. 237, pp. 1630–1632; Dec. 21, 1953.) Observations indicate that simple storms, i.e. those caused by a single bipolar group of sunspots, exhibit a directional effect varying with wavelength.

523.745:621.396.822.029.62 1044

Radio Noise Bursts from Solar *M*-Regions—A. K. Das and B. N. Bhargava. (*Nature*, (*London*), vol. 172, pp. 855–856; Nov. 7, 1953.) The passage of *M*-regions is characterized by a notable increase in the short-term variability of noise intensity at 3 m λ.

523.755:537.56 1045

Ionization and Recombination Processes in a Plasma and the Ionization Formula of the Solar Corona—G. Elwert. (*Z. Naturf.*, vol. 7a, pp. 432–439; June, 1952.) For a plasma not in thermodynamic equilibrium, formulae are derived for (*a*) the photo-ionization, (*b*) the ionization by collisions, (*c*) the recombination with emission of light, and (*d*) the recombination due to triple collisions. By equating the number of collision ionizations with the number of photo-recombinations, an ionization formula is obtained for the solar corona from which the temperature of the corona can be estimated.

523.852:621.396.822.029.62 1046

Radio-Frequency Radiation from the Spiral Nebula Messier 81—R. H. Brown and C. Hazard. (*Nature*, (*London*), vol. 172, p. 853; Nov. 7, 1953.) The intensity measurement of radio source observed has been combined with previous observations to determine the ratio between the radio flux and the light flux for individual nebulae. The calculated flux from a nebula of apparent magnitude +10 (Shapley-Ames) is $\sim 4.2\times10^{-26}$ W/m^2/(cps) at 158.5 mc.

523.852.3:621.396.822 1047

An Extended Radio-Frequency Source of Extragalactic Origin—R. H. Brown and C. Hazard. (*Nature*, (*London*), vol. 172, pp. 997–998; Nov. 28, 1953.) Extension of observations noted in 2215 of 1952.

550.3:519.2 1048

Statistical Analysis of Geophysical Time Series—R. P. W. Lewis, and D. H. McIntosh. (*Met. Mag.*, (*London*), vol. 82, pp. 239–242; Aug., 1953.) Notes are provided on dealing with the effects of linear and periodic trends on standard deviation and correlation coefficient. Significant tests in coherent series are discussed.

550.384 1049

The Phased-Superposed-Epoch Method of Analysis, and an Application to Geomagnetic Activity—E. J. Chernosky. (*Trans. Amer. Geophys. Union*, vol. 34, pp. 519–528; Aug., 1953.) A method is described in which geomagnetic-activity variations are analyzed by grouping the data according to their phases (increasing, decreasing, or no-change), as determined by considering the characteristics for pairs of successive days.

550.384:523.74/.75 1050

Geomagnetic Activity and the Sunspot Cycle—R. Ananthakrishnan. (*Nature*, (*London*), vol. 172, pp. 854–855; Nov. 7, 1953.) A table is given, for the last four sunspot-maximum epochs, of the numbers of magnetically calm days and of slightly, moderately and highly disturbed days, together with sunspot numbers and a measure of the prominence activity. Geomagnetic activity shows correlation with prominence activity rather than with sunspot activity.

550.384.4 1051

Diurnal Magnetic Variations near the Magnetic Equator—S. K. Pramanik and P. S. Hariharan. (*Indian Jour. Met. Geophys.*, vol. 4, pp. 353–358; Oct., 1953.) Observations made at nine places in South India are analysed in conjunction with observations for South America. High diurnal ranges of H occur in a belt of 5 degrees to 6 degrees of latitude, with a maximum near the magnetic equator.

551.510.535 1052

Abnormal Ionospheric Absorption observed during February, 1952—R. Eyfrig. (*Ann. Géophys.*, vol. 9, pp. 325–327; Oct./Dec., 1953.) Observations from stations all over the world are analyzed. Very high absorption was observed during the period 24th–28th, February within a limited region. The eastern and western limits of the affected region are not known accurately, because of the sparseness of stations; no effect was observed in America or in the far east. The southern limit was observed in Africa, roughly along the line joining Casablanca and Djibouti. An explanation involving corpuscular radiation fits the observations better than one based directly on wave radiation from the sun. The absorption is thought to occur in low-altitude ionized layers such as that detected by Gnanalingam and Weekes (3427 of 1952).

551.510.535 1053

A Self-Consistent Calculation of the Dissociation of Oxygen in the Upper Atmosphere: Part 2—Three-Body Recombinations—H. E. Moses and Ta-You Wu. (*Phys. Rev.*, vol. 91, pp. 1408–1409; Sept. 15, 1953.) Continuation of work noted in 102 of 1953. Results based on the assumption that recombination occurs mainly as a three-body nonradiative process indicate that dissociation occurs at a level about 5 km higher than that previously calculated.

551.510.535 1054

Bifurcation of the E Region—J. R. Lien, R. J. Marcou, J. C. Ulwick, D. R. McMorrow, D. B. Linsford, and O. C. Haycock—(*Phys. Rev.*, vol. 92, pp. 508–509; Oct. 15, 1953.) The retardation time, relative to an uhf reference signal, of a radio signal of frequency about 1 mc above the critical frequency was measured by rocket-borne instruments. Electron-density /altitude graphs were derived from the retardation-time/time-of-flight records; bifurcation of the E layer is clearly indicated. The results are in close agreement with those calculated from records obtained with the N.B.S. Model C-3 ionosphere recorder. The separation of the electron-density maxima was 13–18 km.

551.510.535 1055

Direction-Finding Studies of Large-Scale Ionospheric Irregularities—E. N. Bramley. (*Proc. Roy. Soc. A*, vol. 220, pp. 39–61; Oct. 22, 1953.) Continuation of work noted in 3091 of 1951 (Bramley and Ross). A report is given of measurements on pulse-modulated transmissions in the frequency range 2–15 mc, reflected from the ionosphere. Both vertical-incidence and oblique-incidence observations were made. The directional variations of the reflected signals were studied, and are interpreted as due to tilting or wrinkling of the constant-ionic-density surfaces in the ionosphere. The results indicate large-scale horizontal movements in the ionosphere. In the E_s layer these appear to be of the nature of drifting clouds of ionization, while in the F layer the effects are consistent with horizontally travelling ripples having wavelengths of 50–400 km and speeds up to 350 m/s. The direction of motion is more often towards east or west than towards north or

south, and evidence of a diurnal variation has been observed.

551.510.535 1056

Reflexions from Irregularities in the Ionosphere—G. H. Munro. (*Proc. Roy. Soc. A*, vol. 219, pp. 447–463; Oct. 7, 1953.) Complexities in records of *h't* and *h'f* for the *F* region are shown to be due to curvature of the reflecting surface. The order of curvature associated with travelling disturbances is calculated and the variation in the type of complexity is shown to result from differing group retardations in the different paths. A 50-km-base triangular system of three 5.8-mc pulse-transmitters, and a panoramic type 1-15-mc recorder were used respectively for *h't* and *h'f* recording.

551.510.535 1057

Examination of the Formation of the Ionosphere F Region—K. Rawer and E. Argence. (*Ann. Géophys.*, vol. 9, pp. 1–25; Jan./March, 1953.) Continuation of work noted in 987 of 1952 (Rawer et al.). Two different possible mechanisms are discussed, namely (*a*) photoionization of O_2, and (*b*) ionization by soft X rays. The photoionization hypothesis is only acceptable if the dissociation of the O_2 is assumed to take place at a relatively high altitude of about 130 km. Ionization by X rays of solar origin certainly occurs, but a recent calculation of the spectral distribution of solar energy indicates too high an ionization at 95–110 km.

551.510.535 1058

Recombination Coefficient in the F-Regions: a Possible New Process of Ionization of Nitrogen Molecules—R. B. Banerji, (*Nature*, (*London*), vol. 172, pp. 953–954; Nov., 21, 1953.) Discussion of a possible dissociative recombination process according to which the daytime solar radiation at visible wavelengths may be instrumental in ionizing nitrogen molecules if positively charged particles are present.

551.510.535:523.78 1059

The Solar Eclipse of 25-2-1952 at Gao—F. Delobeau. (*Ann. Géophys.*, vol. 9, pp. 317–324; Oct./Dec., 1953.) Ionosphere observations made during the eclipse are reported. Though the occultation was only partial, a diminution of *E*- and F_2-layer ionization compared with normal days was observed; the effect was very pronounced at 0735 hours U.T., when the sunspot near the western limb of the sun was obscured. On the other hand, when this sunspot disappeared on 26th February, as a result of the rotation of the sun, there was no noticeable reduction of the *E*-layer ionization. The effect is difficult to explain in terms of classical theory. The diminution of ionization of the F_2 layer is significant.

551.510.535"1953":621.396.11 1060

Ionosphere Review: 1953—T. W. Bennington. (*Wireless World*, vol. 60, pp. 66–68; Feb., 1954.) Twelve-month running averages and monthly-mean values are plotted for the sunspot numbers and the noon and midnight F_2 critical frequencies from 1947. It is inferred that the critical frequencies are likely to decrease only very slightly before the next sunspot minimum, and that the latter is to be expected between March, 1954 and April, 1955.

551.594:621.317.72 1061

The Agrimeter for Continuous Recording of the Atmospheric Electric Field—J. A. Chalmers. (*Jour. Atmos. Terr. Phys.*, vol. 4, pp. 124–128; Oct., 1953.) Description of an instrument based on measurement of the "bound" charge on a portion of the earth's surface.

551.594.221 1062

The Pilot Streamer in Lightning and the Long Spark—B. F. J. Schonland. (*Proc. Roy. Soc. A.*, vol. 220, pp. 25–38; Oct. 22, 1953.)

551.594.221:621.396.969.029.63 1063

The Study of Lightning Streamers with 50 cm Radar—F. J. Hewitt. (*Proc. Phys. Soc.*, vol. 66, pp. 895–897; Oct. 1, 1953.) Radar echoes were observed during the whole period between the successive strokes in a lightning flash. Some components of these echoes appear to be associated with the processes leading up to the next stroke. Photographic cro-trace records of the stroke, at its commencement and 12 ms later, are reproduced and discussed.

551.594.6 1064

Methods of Synchronizing the Observations of a "Sferics" Network—A. L. Maidens. (*Met. Mag.*, (*London*), vol. 82, pp. 267–270; Sept., 1953.) Methods in use in the British network are described.

LOCATION AND AIDS TO NAVIGATION

621.396:551.594.6 1065

The Reception of Atmospherics at High Frequencies—F. Horne. (*Jour. Atmos. Terr Phys.*, vol. 4, pp. 129–140; Oct., 1953.)Report on hf atmospheric-noise observations made at a single point (*Teddington*) with a view to assessing the practical possibility of basing estimation of source distances on them. Atmospherics were recorded simultaneously at two frequencies in the upper part of the hf band, and the source identity was checked from crdf data. Observations were expected to, and generally did, exhibit the following features: (*a*) reception at both frequencies from sources within ground-wave range; (*b*) no reception when the observation station lies in the skip zone for both frequencies; (*c*) reception on either the lower frequency or both frequencies, according as the observation station lies beyond the skip zone either for the lower frequency or for both frequencies. Difficulties of the method and results achieved are discussed.

621.396.96 1066

Theoretical Performance of Airborne Moving Target Indicators—F. R. Dickey, Jr. (*Trans. I.R.E.*, no.PGAE-8, pp. 12–23; June, 1953.) Fluctuations in the strength of echoes from fixed objects, due to the movement of the antenna, are considered and simple formulae developed to facilitate ground clutter attenuation calculations. The expression for the ratio of mean-square voltage change between two successive signals received from the ground to to the mean-square signal voltage is the sum of four terms, one involving rotation, and the other three the displacement components. Each term is developed in a power series, three different beam shapes being taken into account.

621.396.96:621.396.677.859 1067

Designing Radomes for Supersonic Speeds—S. S. Oleesky. (*Electronics*, vol. 27, pp. 130–135; Jan., 1954.) Designs which will not attenuate or distort the radar beam are discussed.

621.396.963.325 1068

A High-Definition General-Purpose Radar—J. W. Jenkins, J. H. Evans, G. A. G. Wallace, and D. Chambers. (*Jour. Brit. IRE*, vol. 14, pp. 5–21; Jan., 1954. Discussion pp. 21–23.23.) A detailed description is given of equipment for operation in the 3-cm waveband; it is compact enough to be mobile, and can be used with only minor modifications for local aircraft control, airfield surface-movement control, harbor control, etc. Variable pulse width and automatically centering timebase are used. The paraboloid antenna reflector is moulded from a plastic of good dimensional stability. Ppi display is given on two 15-inch tubes.

621.396.969 1069

Some Possible Reductions in Gust Loads through Use of Radar in Transport Airplanes—H. B. Tolefson. (*Bull. Amer. Met. Soc.*, vol. 34, pp. 187–191; May, 1953.) Analysis of turbulence measurements made in different weather conditions indicates that a 10 per cent reduction of the magnitude of the largest gust loads might be achieved, but that little if any reduction of loads due to the more numerous weak gusts is to be expected from the use of airborne radar to detect storm areas. Passenger comfort could be appreciably increased. The use of contour radar is desirable for selecting the smoothest path through storm areas that cannot be avoided.

MATERIALS AND SUBSIDIARY TECHNIQUES

535.371 1070

Fluorescence, Phosphorescence, and Photostimulation of NaCl(AgCl) with High-Energy Irradiation—M. Furst and H. Kallmann. (*Phys. Rev.*, vol. 91, pp. 1356–1367; Sept. 13, 1953.)

537.224 1071

Studies of the Heterocharges of Electrets—S. Wikström. (*Ericsson Tech.*, vol. 9, pp. 225–234; 1953.) Report of investigations of the dependence of the charging and discharging rate on temperature, the influence of a dc field at room temperature on a prepared electret, the influence of pressure during charging and discharging, and the influence of polarization on the capacitance and loss angle.

537.226:[546.48.882.5+546.817.882.5 1072

Dielectric Properties and Phase Transitions of $Cd_2Nb_2O_7$ and $Pb_2Nb_2O_7$—G. Shirane and R. Pepinsky. (*Phys. Rev.*, vol. 92, p. 504; Oct. 15, 1953.) Measurements were made at temperatures down to 40 degrees K; at a frequency of 10 kc and a field strength of 5 v/cm. The dielectric constant of cadmium niobate, which is 310 at room temperature, rises to a maximum of 1200 at 185 degrees K.; the Curie point. The corresponding constants for lead niobate are 185 and 360 (at 14 degrees K.) respectively. Ferroelectric-type hysteresis and a dielectric constant anomaly at 85 degrees K. were observed in the former material only.

537.226:[546.48.882.5+546.817.882.5 1073

Low-Temperature Dielectric Properties of Cadmium and Lead Niobates—J. K. Hulm. (*Phys. Rev.*, vol. 92, pp. 504–505; Oct. 15, 1953.) The results of measurements at temperatures down to 1.2 degrees K., at 1 kc and 10v/cm, are similar to those reported by Shirane and Pepinsky (1072 above).

537.311.1 1074

Changes in the Electrical, Thermal, and Thermoelectrical Properties of Monovalent Metals by Lattice Distortions—A. W. Saénz. (*Phys. Rev.*, vol. 91, pp. 1142–1151; Sept. 1, 1953.) A transport equation is derived for the distribution of conduction electrons in monovalent metals with static lattice distortions, simple assumptions regarding interaction energy being made. Using an iteration method, and considering temperatures very much higher than the Debye temperature of the lattice, a set of integral equations is obtained and solved explicitly. General expressions are found for the resistivity, thermal conductivity and absolute thermoelectric power tensors and are applied to specific calculations for a metal having an array of parallel positive-negative edge dislocations.

537.311.33:538.632 1075

Current Carrier Lifetimes Deduced from Hall Coefficient and Resistivity Measurements—L. P. Hunter, E. J. Huibregtse, and R. L. Anderson. (*Phys. Rev.*, vol. 91, pp. 1315–1320; Sept. 15, 1953.) A method for determining the lifetime of charge carriers in semiconductors, based on theory presented by Landauer and Swanson (165 of January), gives results in reasonable agreement with those obtained on the same samples by injection methods.

537.311.33:546.23 1076

Anisotropic Resistivities of Selenium Crystals at High Frequencies—H. W. Henkels and J. Maczuk. (*Phys. Rev.*, vol. 91, pp. 1562–1563; Sept. 15, 1953.) Some measurements are reported; results are compared with values obtained previously.

537.311.33:546.23 1077

The Structure of Amorphous Selenium—H.

Richter, W. Kulcke and H. Specht. (*Z. Naturf.*, vol. 7a, pp. 511–532; Aug., 1952.) A study of the amorphous material, prepared in various ways, and of its transition to the crystalline state.

537.311.33:546.26-1 1078
The Electrical Resistance of Graphite at Low Temperatures—J. M. Reynolds, H. W. Hemstreet, and T. E. Leinhardt. (*Phys. Rev.*, vol. 91, pp. 1152–1155; Sept. 1, 1953.) The resistance of large-crystal natural graphite increases with increasing temperatures from 1.35 degrees K. to room temperature. That of polycrystalline graphite decreases over practically the whole of this range; at 5 degrees K. it is nearly independent of temperature.

537.311.33:[546.28+546.289 1079
Electron Multiplication in Silicon and Germanium—K. G. McKay and K. B. McAfee. (*Phys. Rev.*, vol. 91, pp. 1079–1084; Sept. 1, 1953.) "Electron multiplication in silicon and germanium has been studied in the high fields of wide *p-n* junctions for voltages in the prebreakdown region. Multiplication factors as high as eighteen have been observed at room temperature. Carriers injected by light, alpha particles, or thermal-generation are multiplied in the same manner. The time required for the multiplication process is less than 2×10^{-8} second. Approximately equal multiplication factors are obtained for injected electrons and injected holes. The multiplication increases rapidly as 'breakdown voltage' is approached. The data are well represented by ionization rates computed by conventional avalanche theory. In very narrow junctions, no observable multiplication occurs before Zener emission sets in, as previously reported. It is incidentally determined that the efficiency of ionization by alpha particles bombarding silicon is 3.6 ± 0.3 electron volts per electron-hole pair produced."

537.311.33:546.28 1080
Paramagnetic Resonance in N- and P-Type Silicon—F. K. Willenbrock and N. Bloembergen. (*Phys. Rev.*, vol. 91, p. 1281; Sept. 1, 1953.) Experiments were performed at 9 kmc at 78 degrees K. The absorption signal was roughly proportional to impurity concentration

537.311.33:546.289 1081
Absorption of Infrared Light by Free Carriers in Germanium—H. B. Briggs and R. C. Fletcher. (*Phys. Rev.*, vol. 91, pp. 1342–1346; Sept. 15, 1953.) Measurements were made of the radiation absorbed when free carriers were injected across a *p-n* junction; the absorption was proportional to the carrier concentration. The absorption/wavelength curve exhibits the maxima previously observed in ordinary *p*-type Ge; the positions of these maxima depend on temperature. An explanation of the observations based on degenerate energy bands is advanced.

537.311.33:546.289 1082
Infrared Absorption in P-Type Germanium—W. Kaiser, R. J. Collins and H. Y. Fan. (*Phys. Rev.*, vol. 91, pp. 1380–1381; Sept 15, 1953.) Measurements made at temperatures down to 5 degrees K. indicate that the absorption coefficient is proportional to the hole concentration. There is a strong absorption band above 10 μ, caused by excitations within the filled band due to the presence of holes, and two weaker bands at shorter wavelengths, one of which disappears at low temperatures.

537.311.33:546.289 1083
Optical Studies of Injected Carriers: Part 1—Infrared Absorption in Germanium—R. Newman. (*Phys. Rev.*, vol. 91, pp. 1311–1312; Sept. 15, 1953.) Measurements were made of the amount of radiation transmitted through a *p-n*-junction Ge diode into which carriers were injected in pulses. The results indicate that for injection into low-resistivity material (*p*-type) the injection current is proportional to the injected carrier density, while for injection into high-resistivity material (*n*-type) there is a departure from proportionality.

537.311.33:546.289 1084
Optical Studies of Injected Carriers: Part 2—Recombination Radiation in Germanium—R. Newman. (*Phys. Rev.*, vol. 91, pp. 1313–1314; Sept. 15, 1953.) Measurements were made of the radiation produced by direct recombination of electrons and holes in *p-n*-junction Ge diodes into which carriers were injected in pulses. The observed spectral distribution is in fair agreement with theory. The intensity of radiation is proportional to the injection current for both high- and low-resistivity material. Polarization of the radiation at oblique angles of emergence was observed.

537.311.33:546.289 1085
Surface Recombination in Germanium—W. N. Reynolds. (*Proc. Phys. Soc.*, vol. 66, pp. 899–901; Oct. 1, 1953.) The variation of the surface recombination velocity of injected minority carriers with temperature was determined for five different surfaces. The absolute values were found to be very sensitive to changes in the etching procedure.

537.311.33:546.289 1086
Electrical Conductivity of Mechanically Disturbed Germanium Surfaces—E. N. Clarke and R. L. Hopkins. (*Phys. Rev.*, vol. 91, pp. 1566–1567; Sept. 15, 1953.) Experiments indicate that the effect of polishing, sandblasting, etc., is to produce a surface layer of relatively high conductivity.

537.311.33:546.289 1087
Microwave Observation of the Collision Frequency of Holes in Germanium—T. S. Benedict. (*Phys. Rev.*, vol. 91, pp. 1565–1566; Sept. 15, 1953.) Using the technique previously described [2330 of 1953 (Benedict and Shockley)], measurements were made of the permittivity and conductivity of *p*-type Ge in order to determine the effective mass of the holes and the relaxation time. Results are shown graphically and discussed.

537.311.33:546.289 1088
The Conductivity of Germanium at 2.4×10^{10} c/s—Y. Klinger. (*Phys. Rev.*, vol. 92, pp. 509–510; Oct. 15, 1953.) The conductivity of a highly purified Ge sample over the temperature range 0–100 degrees C. was determined from the results of permittivity and loss-tangent measurements. At 100 degrees C. the hf conductivity is approximately half the dc conductivity; the hf and dc activation energies are 0.60 and 0.775 ev respectively.

537.311.33:546.289 1089
Electrical Properties of Gold-Germanium Alloys—W. C. Dunlap, Jr. (*Phys. Rev.*, vol. 91, p. 1282; Sept. 1, 1953.) Au acts as an acceptor capable of taking up electrons at two distinct energy levels, the first 0.15 ev above the filled band, the second 0.2 ev below the conduction band. At 77 degrees K. resistivities up to 5×10^{7} Ω.cm have been obtained.

537.311.33:546.289 1090
The Atomic Heat of Germanium below 4 degrees K—P. H. Keesom and N. Perlman. (*Phys. Rev.*, vol. 91, pp. 1347–1353; Sept. 15, 1953.)

537.311.33:546.289:537.323 1091
Measurement of the Thermoelectric Power of Germanium at Temperatures above 78 degrees K—A. E. Middleton and W. W. Scanlon. (*Phys. Rev.*, vol. 92, pp. 219–226; Oct. 15, 1953.) An experimental investigation on pure and impure Ge over the temperature range 78 degrees to 925 degrees K. At low temperatures, where conduction is due to impurity-introduced carriers, the sign of the thermoelectric power is the same as that of the Hall coefficient, the magnitudes of the thermoelectric power, resistivity and Hall coefficient increase with a decrease in the impurity content. With rising temperature the magnitude of the thermoelectric power increases and passes through a maximum, the value for *p*-type material subsequently passing through zero. The results are presented graphically.

537.311.33:546.289:537.323 1092
Theory of Thermoelectric Power in Semiconductors with Applications to Germanium—V. A. Johnson and K. Lark-Horovitz. (*Phys. Rev.*, vol. 92, pp. 226–232; Oct. 15, 1953.) An expression for the thermoelectric power is obtained in terms of the Fermi level, width of forbidden band, temperature, ratio of electron to hole mobility, and effective electron and hole masses. For the impurity range, the formula can be simplified to involve only the Hall coefficient and the temperature. Consideration is given to the case when charge carriage by both holes and electrons must be taken into account. The results are compared with those obtained experimentally by Middleton and Scanlon (1091 above).

537.311.33:546.289:537.323 1093
Thermoelectric Power of Germanium below Room Temperature—H. P. R. Frederikse. (*Phys. Rev.*, vol. 92, pp. 248–252; Oct. 15, 1953.) Measurements were made on *n*-type Ge samples in the range 10 degrees to 300 degrees K. The temperature dependence at temperatures above 200 degrees K. is in good agreement with conventional theory but below this temperature the thermoelectric power rises sharply above the predicted value and reaches a maximum of several millivolts per degree (with Cu) at 15 degrees K. This deviation is due to the disturbance of the phonon equilibrium.

537.311.33:546.289:539.3 1094
Germanium under Ultrasonic Stress: Part 1—Anelastic Effects—G. S. Baker, L. M. Slifkin, and J. W. Marx. (*Jour. Appl. Phys.*, vol. 24, p. 1331; Oct., 1953.) Five samples of 30-Ω.cm single crystals, in the form of right prisms with the major axis along the [100] direction, were driven in fundamental longitudinal vibration at 111 kc. Young's modulus E_{100} was measured at temperatures between 26 degrees and 876 degrees C. Anomalies were observed in two samples.

537.311.33:546.289:539.3 1095
Germanium under Ultrasonic Stress: Part 2—Dynamic Yield Point—G. S. Baker, L. M. Slifkin, and J. W. Marx. (*Jour. Appl. Phys.* vol. 24, pp. 1331–1332; Oct., 1953.) The transient and the residual effects on the yield stress were investigated at temperatures up to 614 degrees C. for mechanically normal Ge specimens of conductivity ~20Ω.cm. but with the major axis parallel to the [111] direction, and driven in fundamental longitudinal vibration at ~ 37 kc.

537.311.33:546.472.21 1096
Some Electrical and Optical Properties of Synthetic Single Crystals of Zinc Sulfide—W. W. Piper. (*Phys. Rev.*, vol. 92 pp. 23–27; Oct. 1, 1953.) A report of measurements of hexagonal *ZnS* (wurtzite) to determine the energy band gap.

537.311.33:546.472.21 1097
Perfect Crystals of Zinc Sulfide—W. W. Piper and W. L. Roth. (*Phys. Rev.*, vol. 92, p. 503; Oct. 15, 1953.) Differences between the properties of *ZnS* crystals grown by sublimation in an atmosphere of hydrogen and those of ordinary *ZnS* crystals are discussed.

537.311.33:546.682.86 1098
Optical Properties of Indium Antimonide—N. Tanenbaum and H. B. Briggs. (*Phys. Rev.*, vol. 91, pp. 1561–1562; Sept. 15, 1953.) Measurements indicate that the intrinsic limit of absorption lies at 7 μ, and that the presence of Ni as an impurity is at least partly responsible for anomalous transmission in the range 3–7 μ.

537.311.33:546.86.48 1099

Impurity and Intrinsic Semiconduction of Intermetallic Compounds: Part 2—E. Justi and G. Lautz. (*Z. Naturf.*, vol. 7a, pp. 602–613; Sept., 1952.) An experimental investigation is made of the temperature dependence of the electrical conductivity, the variation of resistivity in a magnetic field, the differential thermoelectric power and the rectifying properties of CdSb. The results show that, as predicted by the theory presented in part 1 (3619 of 1953), CdSb is an intrinsic semiconductor in the stoichiometric state, and can be turned into an impurity semiconductor by departure from stoichiometric composition or by adding other metals.

537.311.33:621.396.822 1100

A Possible Mechanism for 1/f Noise Generation in Semiconductor Filaments—L. Bess. (*Phys. Rev.*, vol. 91, p. 1569; Sept. 15, 1953.) Analysis is presented based on the Brownian motion of a surface atom displaced from its equilibrium position.

538.221 1101

Orientational Superstructures—L. Néel. (*Compt. Rend. Acad. Sci.*, (*Paris*), vol. 237, pp. 1613–1616; Dec. 21, 1953.) Discussion of phenomena induced in solid solutions of ferromagnetic materials with crystals of cubic structure, when cooled so as to retain permanent anisotropy. Magnetic anisotropy of various ferromagnetic alloys is attributed to this phenomenon.

538.221 1102

Definition of a Ferromagnetic State with Maximum Stability—J. Creusot and A. Langevin. (*Compt. Rend. Acad. Sci.*, (*Paris*), vol. 237, pp. 1499–1500; Dec. 9, 1953.) B/H measurements were made on various materials with the object of finding the most stable equilibrium state corresponding to each value of H, so that B is a two-valued function of H. Superposed on the magnetizing field H is an alternating field whose strength is varied continuously down to zero; B is then measured. The formula obtained for B as a function of H is in agreement with that proposed by P. Langevin in 1911.

538.221 1103

Surface Anisotropy of Ferromagnetic Substances—L. Néel. (*Compt. Rend. Acad. Sci.*, (*Paris*), vol. 237, pp. 1468–1470; Dec. 9, 1953.)

538.221 1104

Study of the Ferromagnetism of Multicrystal and Single-Crystal Cementite [Fe_3C]—P. Blum and R. Pauthenet. (*Compt. Rend. Acad. Sci.*, (*Paris*), vol. 237, pp. 1501–1502; Dec. 9, 1953.) Measurements over the temperature range 20.4 degrees to 500 degrees K. are reported.

538.221 1105

Theory of Saturation Magnetization in Binary Ferromagnetic Alloys—H. Statz. (*Z. Naturf.*, vol. 7a, pp. 506–511; Aug., 1952.) The electron structure of the alloys Ni-Cu, Ni-Zn, Ni-Co, Fe-Ni, Fe-Co, Fe-V and Fe-Cr is investigated in relation to the saturation magnetization.

538.221 1106

Variation of the Magnetic Permeability of a Mild Steel subjected to a Periodic Stress—G. Vidal and P. Lanusse. (*Compt. Rend. Acad. Sci.*, (*Paris*), vol. 237, pp. 1213–1215; Nov. 16, 1953.) Continuation of work noted previously (446 of February).

538.221:621.318.124 1107

Dielectric Behaviour of Granular Semiconducting Aggregates, with Special Reference to Some Magnesium Ferrites—A. Fairweather and E. J. Frost. (*Proc. IEE*, part IIA, vol. 100, no. 3, pp. 15–22; 1953.) "Aggregated granular semiconductors can display high permittivities and dispersion effects which vary with temperature and voltage. This behavior need not be characteristic of the granule material: it can be a consequence of its conductivity and of a particular kind of inhomogeneity of the aggregate arising from the contact structure of the intergranular boundaries. The dielectric properties of certain sintered magnesium ferrites can be accounted for in this way."

538.221:621.318.124:538.66 1108

On the Magnetization of Magnesium Ferrite—N. Sakamoto, T. Asahi, and S. Miyahara. (*Jour. Phys. Soc.*, (*Japan*), vol. 8, pp. 677–678; Sept./Oct., 1953.) An experimental investigation of the effect of heat treatment is reported. The results are summarized in a table showing the rates of cooling from 100 degrees C. and the saturation magnetization per mol at room temperature. A table is also given of the characteristic temperatures, derived from Néel's chemical-equilibrium condition (3159 of 1949), corresponding to the various quenching temperatures.

538.221:621.318.134 1109

Influence of the Ionic Diameters of the Rare Earths on the Ferromagnetic Properties of their Ferrites—G. Guiot-Guillain. (*Compt. Rend. Acad. Sci.*, (*Paris*), vol. 237, pp. 1654–1656; Dec. 21, 1953.) The ferromagnetic Curie points for three groups of rare-earth ferrites exhibiting distinct properties are shown to be related to the ionic diameters by simple expressions. An explanation is advanced of the observations of two Curie points in one group.

538.23 1110

Some New Relations connecting the Magnitudes of Losses with the Coercive Force—N. S. Akulov and T. A. Elkina. (*Compt. Rend. Acad. Sci.*, (*U.R.S.S.*, vol. 83, pp. 377–379; March 21, 1952. In Russian.) The irreversible magnetization (retentivity) and the losses in weak fields can be calculated by assuming ferromagnetic materials to consist of elementary regions with rectangular hysteresis loops [Preisach (3691 of 1935)], the coercive force varying from region region to region. The experimental and calculated losses for several materials are in good agreement.

538.632:[546.817.241+546.331.31 1111

Sensitive Hall Measurements on NaCl and on Photoconductive PbTe—J. L. Levy. (*Phys. Rev.*, vol. 92, pp. 215–218; Oct. 15, 1953.) Measurements were made on NaCl over the temperature range 650–795 degrees C. and on PbTe at 190 degrees C. and at 30 degrees C. Because of the high noise level, only an upper limit for the Hall mobility could be determined for NaCl. The results for PbTe support Simpson's theory (861 of 1952) that an increase in photoconductivity is primarily due to an increase in carrier density and only secondarily to a change of mobility.

539.153:546.281.26 1112

Energy of Trapped Electrons in Ionic Solids—K. Lehovec. (*Phys. Rev.*, vol. 92, pp. 253–258; Oct. 15, 1953.) The energy of an electron localized in a crystal at an impurity carrying multiple positive charges is calculated. The results are applied to a tentative interpretation of activation energies of SiC crystals observed by Busch (745 of 1947).

546.26:539.215.3:537.311.4 1113

The Properties of Carbon Contacts—R. O. Grisdale. (*Jour. Appl. Phys.*, vol. 24, pp. 1288–1296; Oct., 1953.) The correlation between the contact properties and the crystallographic structure of the surfaces is relatively little influenced by the intercrystal boundaries. The contact properties are determined by those of the small number of crystals in contact. The physical properties of these crystals are studied.

546.431–31 1114

Oxygen Vacancies in Barium Oxide—R. L. Sproull, R. S. Bever, and G. Libowitz. (*Phys. Rev.*, vol. 92, pp. 77–80; Oct. 1, 1953.) The diffusion of the blue coloration produced when BaO crystals are heated in certain metal vapours (1115 below) was measured at temperatures between 800 degrees and 1300 degrees C. Comparison of the results with those obtained by Redington (432 of 1953) indicates that the diffusion process associated with the coloration does not involve the transport of Ba, and that the principal lattice defects in BaO with excess are 0 vacancies.

546.431–31:[535.215+535.343.2 1115

Optical Absorption and Photoconduction in the Visible and Near Infrared in Single Crystals of BaO—W. C. Dash. (*Phys. Rev.*, vol. 92, pp. 68–76; Oct. 1, 1953.) Optical-absorption measurements indicate the presence of absorption bands at 0.8 and 1.4 ev, induced by ultraviolet or X-ray irradiation at about −160 degrees C. These bands are enhanced about fivefold on heating a crystal in air to about 1600 degrees C. and quenching. Absorption curves for crystals heated in Ba, Al, Mg or Ca vapor all exhibit a maximum at 2 ev, which is attributed to interstitial Ba or O vacancies. Photoconduction studies indicate the presence of energy levels at 2 and 2.6 ev as well as those at 0.8 and 1.4 ev.

546.431.824–31 1116

Temperature Changes of Single Crystals of $BaTiO_3$—I. N. Belyaev, N. S. Novosil'tsev, E. G. Fesenko, and A. L. Khodakov. (*Compt. Rend. Acad. Sci.*, (*U.R.S.S.*), vol. 83, pp. 675–676; April 11, 1952. In Russian.) The temperature dependence of the dielectric constant and of the crystal cell parameters was determined for two crystals prepared from a solution and one by double decomposition. The results are shown graphically.

546.431.824–31:537.226.33 1117

Time Effects in the Hysteresis Loop of Polycrystalline Barium Titanate—M. McQuarrie. (*Jour. Appl. Phys.*, vol. 24, pp. 1334–1335; Oct., 1953.) The maximum and the remanent polarization decrease with aging by approximately the same amount until the loop shows a constriction at the center. This effect can be partially annulled by application of a strong alternating field, and completely annulled by heating to above the Curie temperature.

621.383.49:546.482.21:535.215.5 1118

Photoresistivity and Photoactivation of CdS Single Crystals—V. E. Lashkarev, V. S. Medvedev, A. I. Skopenko, G. A. Fedorus, and M. K. Sheinkman. (*Compt. Rend. Acad. Sci.* (*U.R.S.S.*), vol. 86, pp. 905–907; Oct. 11, 1952. In Russian.) The increase of light sensitivity due to background illumination was investigated. Graphs are shown of the variation in sensitivity (photoelectric-current/quantum) with the wavelengths of the background illumination λ_b and impulse illumination λ_i, in the visible-light range. The spectral sensitivity in the 2100–4000-Å range is also shown.

546.561.221 1119

On the Electrical Conductivity of Cuprous Sulfide: a Diffusion Theory—I. Yokota. (*Jour. Phys. Soc.* (*Japan*), vol. 8, pp. 595–602; Sept./Oct., 1953.) Cuprous sulphide, a p-type semiconductor, exhibits mixed electronic and ionic conduction in the 110–470 degrees C. β phase. A theory of conduction is developed which gives values for the potential distribution in the material in good agreement with values observed by Miyatani and Suzuki (1120 below), for both transient and steady-state conditions. The temperature dependence of the ionic and hole conductivities and of the mobility of vacant copper ion lattice points is shown graphically.

546.561.221 1120

On the Electric Conductivity of Cuprous Sulfide: Experiment—S. Miyatani and Y. Suzuki. (*Jour. Phys. Soc.* (*Japan*), vol. 8, pp. 680–681; Sept./Oct., 1953.) Yokota's theory of

conduction (1119 above) is used to calculate the mobility and the concentration of lattice defects from the experimental results for β-phase cuprous sulphide.

546.817.831.4.03+546.817.824.03 1121
Ferroelectricity versus Antiferroelectricity in the Solid Solutions of $PbZrO_3$—E. Sawaguchi. (*Jour. Phys. Soc. (Japan)*, vol. 8, pp. 615–629; Sept./Oct., 1953.) The phase diagram of the solid solutions was investigated by measuring the variations with temperature of the dielectric constant, thermal expansion, specific heat and lattice constants. In $Pb(Zr97\text{-}Ti3)O_3$, two antiferroelectric phases, one ferroelectric phase and one paraelectric phase were observed. The free energies in the four states are compared as functions of the temperature and the Ti ion concentration. The results are shown graphically.

548.0:539.378.3 1122
Some Predicted Effects of Temperature Gradients on Diffusion in Crystals—W. Shockley. (*Phys. Rev.*, vol. 91, pp. 1563–1564; Sept. 15, 1953.) Experiments are outlined by means of which it is possible to distinguish between diffusion due to interstitial atoms and diffusion due vacancies.

621.315.61:061.3 1123
Symposium of Papers on Insulating Materials—(*Proc. IEE*, part IIA, vol. 100, pp. 1–308; 1953.) Full report of the proceedings at the symposium held in March, 1953, with index.

MATHEMATICS

517.9 1124
The Eigenvalues of $\nabla^2 u+\lambda u=0$ when the Boundary Conditions are given on Semi-infinite Domains—D. S. Jones. (*Proc. Camb. Phil. Soc.*, vol. 49, part 4, pp. 668–684; Oct., 1953.) Investigation of an equation which represents the small oscillations of many physical systems.

MEASUREMENTS AND TEST GEAR

621.317:537.71(083.74) 1125
The Accuracy of Measurement of Electrical Standards—A. Felton. (*Proc. IEE*, part II, vol. 100, pp. 543–544; Oct., 1953.) Discussion on 1952 of 1952.

621.317:621.383.2 1126
The Photodianode—Deloffre, Pierre, and Roig. (*See* 1249.)

621.317.3:621.396.67.029.64 1127
Measurements on Aerials for Centimetre Waves—M. Bouix. (*Ann. Télécommun.*, vol. 8, pp. 314–326; Oct., 1953.) Methods are discussed for the measurement of (*a*) swr (to test the matching between antenna and feeder), (*b*) phase of primary radiators, (*c*) radiation patterns both for primary radiators and complete aerial systems, (*d*) gain. Equipment described includes transmitter and sensitive receiver for the radiation-pattern and gain measurements. Attention is drawn to special requirements as regards the terrain over which the tests are made.

621.317.3.029.63/.64+534.62 1128
A New Anechoic Chamber for Sound Waves and Short Electromagnetic Waves—Meyer, Kurtze, Severin, and Tamm. (*See* 942.)

621.317.31:621.375.2 1129
Measurement of Weak Electric Charges by means of Pulse Technique—Study of a Fast-Acting Preamplifier with High Input Impedance and Low Noise—H. Guillon. (*Onde élect.*, vol. 33, pp. 627–636; Nov., 1953.)

621.317.31:621.387.4 1130
Measurement of Weak Electric Charges by means of Pulse Technique—α-Radiation Spectrometry—G. Valladas. (*Onde élect.*, vol. 33, pp. 615–626; Nov., 1953.)

621.317.32:537.1 1131
The Static and Dynamic Measurement of Electromotive Forces—W. J. Poppelbaum. (*Helv. Phys. Acta*, vol. 26, pp. 489–498; Sept. 15, 1953. In French.) Generalization of Meixner's equation (*Ann. Phys., Lpz.*, vol. 35, p. 701; 1939.) Leads to a formulation including within its scope phenomena as varied as induction, thermoelectricity and the chemical processes in a cell. Consideration of the distinction between voltaic and galvanic potential difference results in a simple method of taking account of contact potential difference in static measurements.

621.717.335 1132
Precise Measurement of the Complex Dielectric Constants of Liquids by Voltage Resonance—C. Abgrall. (*Compt. Rend. Acad. Sci., (Paris)*, vol. 237, pp. 1650–1652; Dec. 21, 1953.) A substitution method is used in which either a standard capacitor or the cell containing the liquid is connected in an oscillatory circuit in series with a high impedance. Analysis of the system is based on a linear formulation, and the effect of the impedances of the connections is examined.

621.317.355.029.64 1133
Method for the Measurement of the Dielectric Constant of Gases at Ultrahigh Frequencies—A. Gozzini and E. Pollaco. (*Compt. Rend. Acad. Sci., (Paris)*, vol. 237, pp. 1497–1499; Dec. 9, 1953.) A klystron generator, operated at about 10 kmc, is frequency modulated between the limiting frequencies ν_0 and ν_1 by application of a sawtooth voltage to the repeller. Two cavities tuned to different frequencies between ν_0 and ν_1 are associated with the system, one of the cavities being filled with the gas under test. The dielectric constant is determined from the time interval between the instants when the two cavities resonate.

621.317.335.029.64 1134
Complementary Note on the Method of Measurement of the Dielectric Constant of Gases at U.H.F.—A. Gozzini and E. Pollaco. (*Compt. Rend. Acad. Sci., (Paris)*, vol. 237, pp. 1652–1654; Dec. 21, 1953.) A modification of the method previously described (1133 above) to give improved sensitivity is proposed.

621.317.335.3.029.65 1135
A Spectrometer for Millimetre Wavelengths—W. Culshaw. (*Proc. IEE*, part IIA, vol. 100, pp. 5–14; 1953. Discussion, pp. 54–60.) A free-space method is described for the measurement of permitivity at mm λ, using the microwave analogue of the optical spectrometer, horns with lenses taking the place of the optical collimator and telescope. The reflectivity of sheets of material is investigated at different angles of incidence for different polarizations; interference minima and Brewster angle are determined and accurate values of permittivity are hence deduced. The effects of dielectric loss and of diffraction are discussed. Results substantiating the theory are presented. An extension of the method to deal with liquids is described.

621.317.34.018.75:621.315.212 1136
Echo Meter with Very Short Pulse Duration for Investigation of Coaxial-Pair Television Cables—G. Comte, M. Bouderlique, and G. Thévenet. (*Câbles & Transm.*, vol. 7, pp. 263–269; July, 1953.) A description is given, with circuit details, of an instrument designed for production testing of television cables, with application for research at higher frequencies. Features of the equipment are the short pulse duration (0.02 μs) and the use of distributed amplification in the echo amplifier to give a wide frequency band.

621.317.34.029.5/.6]:621.315.212 1137
Tests at Very High Frequencies on Production Lengths of Coaxial Cable—J. Lorrin. (*Câbles & Transm.*, vol. 7, pp. 218–241; July, 1953.) A null method of measurement is described for determination of propagation coefficient and characteristic impedance in production testing of coaxial cables. Measurement principles are based on transmission-line theory taking account of irregularities; this is dealt with in an appendix. Bridge balance is obtained by adjustment of both reactance and frequency, and the design of an oscillator/receiver unit with a single frequency control is discussed. Illustrated descriptions are given of two instruments. Results of measurements on cables of length 100–300 m are shown.

621.317.341.029.62/.65 1138
A.M. System measures Microwave Attenuation—J. Korewick. (*Electronics*, vol. 27, pp. 175–177; Jan., 1954.) See 803 of March.

621.317.35:621.3.018.78 1139
Measuring Non-linearity—D. C. Pressey. (*Wireless World*, vol. 60, pp. 60–62; Feb., 1954.) A simple method is presented which is supplementary to that described by Wigan (2383 of 1953) and which uses a frequency-insensitive element to perform the subtraction of the fundamental from the composite wave.

621.317.38.029.63 1140
Force on a Shorted Ring in a U.H.F. Field in a Coaxial wavity—S. N. Kalra. (*Jour. Appl. Phys.*, vol. 24, pp. 1339–1340; Oct., 1953.) The change in potential energy of a metal ring, the dimensions of which are small compared with λ, is made the basis of power measurement. Accuracy for powers of a few watts is within $\sim$1 per cent.

621.317.44 1141
Dynamic Measurements on Electromagnetic Devices—M. A. Logan. (*Bell. Sys. Tech. Jour.*, vol. 32, pp. 1413–1467; Nov., 1953.) An electronic timing system controlling mercury-contact relays is used to switch in the measuring instrument (e.g. fluxmeter) for a very short period only at a preselected instant of each cycle of operation of the device tested. The system and the circuits required for the dynamic measurement of magnetic flux, current, displacement and velocity response are described.

621.317.444 1142
A Small Sensitive Magnetometer—T. M. Palmer. (*Proc. IEE*, part II, vol. 100, pp. 545–550; Oct., 1953.) The measuring head contains a solenoid through which is passed a dc balancing out the field under test. A saturating current of frequency 5 kc is passed through a fine mumetal wire within the solenoid, causing its effective longitudinal permeability to alternate, so that a field along the wire produces an alternating emf in the solenoid; this emf serves as an indication of the field. A change of 2×10^{-6} oersted can be detected. The measurement range is limited to about 50 oersted due to the heating of the solenoid.

621.317.7 1143
Device for Measurement of the Time Constants of Indicating Instruments—J. Mey and H. Thibert. (*Ann. Télécommun.*, vol. 8, pp. 327–334; Oct., 1953.) The time constant for indication, the time for restoration to zero, and the "integrating" time are considered, particularly in relation to volume and modulation meters for telephone testing. Apparatus devised for measuring these time constants comprises an electro-optical system for marking the passage of a galvanometer needle through a predetermined position, an electronic chronometer, and a "chronotome" relay.

621.317.7.088 1144
Performance Limits in Electrical Instruments—A. H. M. Arnold. (*Proc. IEE*, part II, vol. 100, pp. 543–544; Oct., 1953.) Discussion on 1971 of 1952.

621.317.714.024 1145
Non-Contact D.C. Ammeter—W. H. Bailey. (*Elect. Rev.*, vol. 153, pp. 397–400; Aug. 21,

1953.) Description of a device which permits remote indication and recording of the current in a cable.

621.317.729 1146
The Accurate Mapping of Electric Fields in an Electrolytic Tank—K. F. Sender and J. G. Yates. (*Proc. IEE*, part II, vol. 100, pp. 569–570; Oct., 1953.) Discussion on 2739 of 1953.

621.317.756+621.317.77 1147
A Harmonic-Response-Testing Equipment for Linear Systems—D. O. Burns and C. W. Cooper. (*Proc. IEE*, part II, vol. 100, p. 467; Oct., 1953.) Discussion on 3374 of 1953.

621.317.76+531.771]:621.387 1148
A High-Speed Precision Tachometer—Bland and Cooper. (*See* 1153.)

621.373.43:621.317.3 1149
High-Voltage Sawtooth and Rectangular-Wave Pulse Generator—W. D. Edwards. (*Electronic Eng.*, vol. 26, pp. 36–39; Jan., 1954.) Descriptions are given of a square-wave generator giving an output of 0–20 kv with pulse duration continuously variable between 0.05 and 80 μs, and of a sawtooth generator giving a peak output of 18 kv with rise time of 15 μs-10ms. The generators are specially designed for measurements of dielectric strength.

621.397.5:535.623].001.4 1150
Methods of Verifying Adherence to the N.T.S.C. Color Signal Specifications—A. C. Luther, Jr. (Proc. I.R.E., vol. 42, pp. 235–240; Jan., 1954.) The composition of the color signal is determined by measuring the amplitudes and relative phases of a group of bar signals corresponding to saturated primary colors and their complements. The use of a simple demodulator for accurate measurement of phase at the 3.58-mc subcarrier frequency is described.

621.397.5:535.623].001.4 1151
A Versatile Approach to the Measurement of Amplitude Distortion in Color Television—J. A. Bauer. (Proc. I.R.E., vol. 42, pp. 240–246; Jan., 1954.) The importance of system linearity for maintaining the selected value of gamma is emphasized, and an instrument developed for checking the degree of nonlinearity is described.

621.397.5:535.623].001.4 1152
Test Instruments for Color Television—W. C. Morrison, K. Karstad, and W. L. Behrend. (Proc. I.R.E., vol. 42, pp. 247–258; Jan., 1954.) Equipment for use in measuring frequency response, differential gain and differential phase, and sound-to-picture frequency separation is described.

OTHER APPLICATIONS OF RADIO AND ELECTRONICS

531.771+621.317.76]:621.387 1153
A High-Speed Precision Tachometer—W. R. Bland and B. J. Cooper. (*Electronic Eng.*, vol. 26, pp. 2–8; Jan., 1954.) An integrating-type electronic tachometer with ranges of 0–8,000 and 10–80,000 rpm and a reading accuracy within 0.01 per cent uses cold-cathode glow-discharge tubes to give a visual display. The instrument can be used with a variety of pickups, and can also be used as a frequency meter.

551.508.1:621.396.91 1154
The Ground Equipment for the F-Type Radioisonde—B. B. Huddar, K. Nagarajan, N. C. Dhar, and S. P. Venkiteshwaran. (*Indian Jour. Met. Geophys.*, vol. 4, pp. 347–352; Oct., 1953.) Equipment in use at Poona is described; the receiver is of superregenerative type, and the recorder is of moving-coil type.

621–52:621.389 1155
An Electronic Process Controller—J. R. Boundy and S. A. Bergen. (*Proc. IEE*, part II, vol. 100, pp. 561–562; Oct., 1953.) Discussion on 1386 of 1952.

621.316.722:621.383.27:621.387.464 1156
Device for the Stabilization of Photomultipliers—R. Ascoli. (*Nuovo Cim.*, vol. 9, pp. 615–617; July 1, 1952.) A stabilizer for scintillation counters makes use of the constancy of the emission from a radioactive source.

621.385.832/.833]:538.691 1157
Image Distortion due to Pole-Piece Asymmetry in Electron Lenses—W. Glaser and P. Schiske. (*Z. angew. Phys.*, vol. 5, pp. 329–339; Sept., 1953.)

621.385.833 1158
Electron-Optical Properties of Electrostatic Lenses: Part 2—W. Lippert and W. Pohlit. (*Optik, Stuttgart*, vol. 10, pp. 447–454; 1953.) Further results are given of the work on unipotential lenses previously noted (795 of 1953).

621.385.833 1159
Trajectories in the Symmetrical Electron Lens—L. Jacob and J. R. Shah. (*Jour. Appl. Phys.*, vol. 24, pp. 1261–1266; Oct., 1953.) Numerical evaluation of trajectories in a strong unipotential lens. See also 1396 of 1952 (Shah and Jacob).

621.387.464 1160
The Efficiency of the Anthracene Scintillation Counter—D. K. Butt. (*Proc. Phys. Soc.*, vol. 66, pp. 940–944; Oct. 1, 1953.) Investigation of a counter having a crystal about 1 cm^2×2 mm placed on the end of an uncooled E.M.I. Type-5060 photomultiplier tube.

621.387.464:550.835 1161
A Scintillation Counter for Radioactivity Prospecting—D. H. Pierson and J. Pickup. (*Jour. Brit. IRE*, vol. 14, pp. 25–32; Jan., 1954.) A simple and compact arrangement is obtained by using a scintillation counter in conjunction with a cold-cathode tube in the counting-rate meter circuit and a specially designed recording microammeter.

PROPAGATION OF WAVES

538.566+621.372.2 1162
The Effect of the Radiation Condition in the case of Complex Wave Number, and its Significance in the Problem of Surface Waves—A. Haug. (*Z. Naturf.*, vol. 7a, pp. 501–505; Aug., 1952.) It is shown that Sommerfeld's radiation condition (i.e. that the wave vanishes at infinity) guarantees the uniqueness of the solutions of the wave equation $\Delta u+k^2u=0$ in the whole space, not only for real but also for complex values of the wave number k. It follows from the treatment presented that the existence of Sommerfeld's surface wave is not inconsistent with the radiation condition.

538.566+621.372.2]:621.3.012 1163
New Chart for the Solution of Transmission-Line and Polarization Problems—Deschamps. (*See* 965.)

621.396.11 1164
Nonstandard Radio Propagation—P. G. Forsyth. (*Nature*, (*London*), vol. 172, p. 966; Nov. 21, 1953.) Brief report of long-range reception observed at Ottawa on frequencies between 60 and 72 mc.

621.396.11:551.5 1165
Ultra-Short-Wave Field Strength in a Ground-Based Radio Duct—R. S. Unwin. (*Nature*, (*London*), vol. 172, pp. 856–857; Nov. 7, 1953.) Discussion of field-strength observations at λ 3 cm, 9 cm, and 60 cm, made at various heights over the sea under widely differing superrefractive conditions. A rough estimate can be made of the maximum field strength in the duct at distances up to 100 km from the transmitter, if the average width of the ground-based duct can be assumed from meterological data to be greater than a critical value for the particular wavelength. See also 1471 of 1953 (Hay and Unwin) and back references.

621.396.11:551.510.535 1166
Coupling and Conditions for Reflection of the Ordinary and Extraordinary Electromagnetic Waves in an Inhomogeneous Anisotropic Plasma (Ionosphere)—R. Jancel and T. Kahan. (*Compt. Rend. Acad. Sci.*, (*Paris*), vol. 237, pp. 1657–1659; Dec. 21, 1953.) The propagation equations are solved, taking the plasma as characterized by a dielectric tensor which does not vary with height except insofar as the electron concentration varies, and taking account of the earth's magnetic field. The conditions giving rise to triple splitting are established.

621.396.11:551.510.535 1167
A New Phenomenon of Interaction between Waves and Free Electrons subjected to the Terrestrial Magnetic Field—M. Cutolo. (*Nuovo Cim.*, vol. 9, pp. 687–698; Aug. 1, 1952.) Experiments indicate that modulated waves of carrier frequency equal to the frequency of rotation of free electrons suffer appreciable demodulation in traversing the ionosphere, due to the action of the terrestrial magnetic field. The magnitude of the effect increases with the modulation frequency. Various names, eg self-demodulation, are suggested for the effect. The physical process involved is discussed and the importance of the effect in broadcasting and in the study of the ionosphere E layer is indicated. A long summary in English is included. See also 1758 of 1951.

621.396.11:621.317.353.3:551.510.535 1168
Experimental Determination of the Resonance Curves in the Motion of the Slow Electrons in the Upper Atmosphere—M. Cutolo. (*Nuovo Cim.*, vol. 9, pp. 391–407; May 1, 1952.) Theory of gyrointeraction between waves in an ionized medium is reviewed. Double-hump or single-hump resonance curves are obtained, depending in the experimental conditions. Experiments made in Italy in 1950 are described, when by varying the frequency of the unmodulated wave the transition from the double-hump to the single-hump curve was observed. *See* also 720 of March.

621.396.812:621.396.621.59 1169
Long-Range Communication Trends—Crosby. (*See* 1172.)

RECEPTION

621.396.62/.63 1170
Conelrad Receiver With Built-In Alarm—R. E. Quenstedt. (*Electronics*, vol. 27, pp. 156–157; Jan., 1954.) Description of a FM broadcast receiver with carrier-failure alarm, for use in controlled broadcast stations required to monitor a regional parent station of a civil defence scheme.

621.396.621+621.397.621 1171
Radio and Television at the 16th Salon National—P. A. François. (*TSF et TV*, vol. 29, pp. 359–364 and 393–394; Nov./Dec., 1953.) General report of the Paris exhibition, Sept. 25–Oct. 5, 1953, with tabulated details of the television receivers shown. Reports of the exhibition are also given in *Toute la Radio*, vol. 20, pp. 410–413; Nov., 1953, and *Télévision*, pp. 264–266; Nov., 1953.

621.396.621.59:621.396.812 1172
Long-Range Communication Trends—M. G. Crosby. (*Trans. I.R.E.*, vol. CS-1, pp. 41–53; July, 1953.) Various types of fading due to multipath ionospheric transmission are reviewed and methods of reducing fading effects are discussed. The interference rejection and IF/AF selectivity transfer obtained with exalted-carrier detection are noted. The advantage of frequency-shift keying for telegraphy is pointed out. The application of exalted-carrier and ssb methods in a triple-diversity system is illustrated.

STATIONS AND COMMUNICATION SYSTEMS

621.376.2:621.396.5:621.396.822 1173
Noise and Radiotelephony with Amplitude

Modulation—P. Marcou. (*Ann. Télécommun.*, vol. 8, pp. 339–351; Oct., 1953.) Dsb and ssb AM systems are compared, using as criterion the signal/noise ratio at the receiver output when white noise is mixed with the rf signal. The effect of different detecting processes on the signal/noise ratio is analyzed. The ssb system makes better use of available power inasmuch as (*a*) the carrier is suppressed or reduced, (*b*) the rf noise is halved, and (*c*) reception is linear; the advantage due to this last feature only becomes significant, however, for output signal/noise ratios too low to give good intelligibility even with this improvement. With ssb transmissions, the ear accepts quasi-synchronous demodulation, not encessarily in phase.

621.39.001.11:621.376 1174
Comparative Study of Modulation Methods—R. M. Page. (*Trans. I.R.E.*, vol. CS-1, pp. 13–22; July, 1953.) To express the results of information theory in practical terms, direct, coded and modulated-carrier coded transmission systems are analysed on the basis of Shannon's equation. Basic transformer theory is applied to the action of coding and an expression is derived in relating power ratio to bandwidth ratio in terms of the dynamic range of the information. Curves showing signal/noise ratio at a receiver output as a function of signal input power for ssb AM, dsb AM, wide-band FM and binary-code pcm systems are compared.

621.39.001.11:[621.394+621.396.3 1175
The Transmission of Intelligence in Typescript—I. S. Coggeshall. (*Trans. I.R.E.*, vol. CS-1, pp. 4–13; July, 1953.) Practical aspects of the application of communication theory in telegraphy are reviewed.

621.396.41 1176
Properties of a Multiplex Signal with Pulses of Alternate Sign—L. Le Blan. (*Compt. Rend. Acad. Sci.*, (*Paris*), vol. 237, pp. 1662–1664; Dec. 21, 1953.) Analysis of a two-channel system in which one channel uses positive AM pulses while the other uses negative AM pulses. Simple methods can be used to separate the two signals at the receiver; in general, arrangements must be provided for eliminating cross-talk.

621.396.44 1177
Simplified Carrier-Current System for Short Distances—G. H. Bast and J. L. Hurault. (*Câbles & Transm.*, vol. 7, pp. 185–217; July, 1953.) The paper is based on an investigation for the Netherlands Administration. A detailed study is made of design requirements for a multiplex system operating over distances of 20–300 km on any suitable existing lines. In a discussion of economic aspects, the attenuation required to avoid cross-talk is calculated for different channel widths, and the choice of a 6-kc bandwidth is explained. Full descriptions are given of terminal equipment for a 32-channel system on symmetrical-pair cables and the main oscillators and control circuits at the central station.

621.396.662.029.5/.6 1178
Design Trends in Communication Equipment—L. M. Craft. (*Trans. I.R.E.*, vol. CS-1, pp. 22–30; July, 1953.) General discussion of omnichannel operation of usw and sw equipment. Outline descriptions are given, with illustrations, of automatic tuning and channel-selection units incorporated in Collins transmitters and receivers.

621.396.712 1179
La Maison de la Radio, Paris—H. Testemale. (*Télév. Franç.*, no. 100, pp. 23–25; Nov., 1953.) A competition has been held for the design of a building to serve as a comprehensive broadcasting centre; the winning design is described.

621.396.712.029.62:621.396.81 1180
U.S.W. Planning in Austria—J. Burgstaller. (*Radio Tech.* (*Vienna*), vol. 29, pp. 340–345; Oct., 1953.) Formulas and curves for the theoretical field strength of horizontally polarized waves received at a height 30 feet above ground, and for the correction factor for the attenuation due to diffraction in mountainous districts are given and discussed. The frequency considered is 98 mc. Maps show the estimated area coverage provided by six 25-kw transmitting stations and, in the case of the Klagenfurt transmitter, the effects of different aerial heights and radiated powers.

621.396.932 1181
Aspects of Naval Communications Systems—J. A. Krcek. (*Trans. I.R.E.*, vol. CS-1, pp. 54–58; July, 1953.) General review of problems and developments in the design and arrangement of radio equipment aboard a communications vessel which may carry about 150 receivers and 60 transmitters.

621.396.933 1182
A Discussion of United Air Lines V.H.F. Network Developments—K. J. Rhead. (*Trans., I.R.E.*, no. PGAE-8, pp. 9–11; June, 1953.) The network comprises five transmitter-receiver stations, with automatic locking out of either transmitter of receiver when the other unit is in operation. Preliminary experience indicates that intermodulation effects are considerably reduced by using a stable oscillator, by sub-channelling arrangements, and by restriction of the speech band to the range 200–2,500 cps.

SUBSIDIARY APPARATUS

621–526 1183
Analogue Methods for Optimum Servomechanism Design—F. C. Fickeisen and T. M. Stout. (*Trans. AIEE*, vol. 71, part II, pp. 244–250; 1952. Discussion, p. 250.)

621–526 1184
Feedback Control Systems with Dead-Time Lag or Distributed Lag by Root-Locus Method—Y. Chu. (*Trans. AIEE*, vol. 71, part II, pp. 291–296; 1952.) Generalization and application of the method originated by Evans (2337 of 1952).

621–526 1185
Synthesis of Feedback Control System by Phase-Angle Loci—Y. Chu. (*Trans. AIEE*, vol. 71, part II, pp. 330–339; 1952.) A modification of Evans' root-locus method (2337 of 1952).

621–526 1186
The Use of Nonlinear [tachometric] Feedback to Improve the Transient Response of a Servomechanism—J. B. Lewis. (*Trans. AIEE*, vol. 71, part II, pp. 449–453; 1952. Discussion, p. 453.)

621–526:621.3.016.35 1187
Stabilization of a Servomechanism subject to Large Amplitude Oscillation—E. S. Sherrard. (*Trans. AIEE*, vol. 71, part II, pp. 312–324; 1952.) The frequency-response method of analysis developed by Kochenburger (224 of 1951) is applied to a system containing a saturable amplifier and in determining the required nonlinear stabilizing filter.

621–526:621.3.016.35 1188
Stability Limits for Third-Order Servomechanisms—T. J. Higgins and J. G. Levinthal. (*Trans. AIEE*, vol. 71, part II, pp. 459–466; 1952. Discussion, pp. 466–467.) Curves are derived giving directly the roots and coefficients of specified cubic characteristic equations. Their use in facilitating analysis and design of servo systems is explained and illustrated.

621–526:621.316.7 1189
The Analysis of Sampled-Data Systems—J. R. Ragazzini and L. A. Zadeh. (*Trans. AIEE*, vol. 71, part II, pp. 225–232; 1952. Discussion, pp. 232–234.) Full paper; summary abstracted in 1801 of 1953.

621.311.6 1190
Stable Power Supplies for Microwave Standards—W. P. Ernst. (*Electronics*, vol. 27, pp. 168–171; Jan., 1954.) A description of the power-supply system for operating the klystron frequency standards at the National Bureau of Standards.

621.314.634.004 1191
Selenium Rectifiers—Factors in their Application—J. Gramels. (*Bell. Sys. Tech. Jour.*, vol. 32, pp. 1469–1492; Nov., 1953.)

621.318.5 1192
Design of Relays—(*Bell Sys. Tech. Jour.*, vol. 33, pp. 1–259; Jan., 1954.) A symposium of papers, dealing with design, production, service, and measurements.

TELEVISION AND PHOTOTELEGRAPHY

621.397.2:535.623 1193
The Colorplexer—a Device for Multiplexing Color Television Signals in accordance with the N.T.S.C. Signal Specifications—E. E. Gloystein and A. H. Turner. (Proc. I.R.E., vol. 42, pp. 204–212; Jan., 1954.) The functions of matrixing, band limiting, delay correction, modulating, burst generating, and mixing are described in detail. Waveforms photographed at several points in the colorplexer are shown for one standard color-bar signal.

621.397.2:535.623 1194
Transmission of Color over Inter-City Television Networks—J. R. Rae. (Proc. I.R.E., vol. 42, pp. 270–273; Jan., 1954.) The additional gain and delay equalizers required for color transmission and the method for shifting the color information within the pass band of coaxial cable are discussed.

621.397.24:621. 315.212.4 1195
The Birmingham-Manchester-Holme-Moss Television-Cable System—R. J. Halsey and H. Williams. (*P.O. Elec. Eng. Jour.*, vol. 46, parts 3 and 4, pp. 118–121 and 171–176; Oct., 1953 and Jan., 1954.) Shortened version of paper noted in 1147 of 1953.

621.397.24:621.372.55 1196
Adjustment of Bridged-T Phase-Correction Networks used in Television Equipment—H. Martin. (*Câbles & Trans.*, vol. 7, pp. 175–184; July, 1953.) Properties of equalizing networks for television cables are reviewed. A practical method is described, based on measurement of impedance at a frequency for which the phase shift is 180 degrees, by which network impedance can be adjusted to within $\pm 2\Omega$ of the correct value throughout the pass band. Propagation-time/frequency characteristics are shown for two types of network.

621.397.5 1197
Fundamental Problems of Subscription Television: the Logical Organization of the Telemeter System—L. N. Ridenour and G. W. Brown. (*Jour. Soc. Mot. Pict. & Telev. Engrs.*, vol. 61, pp. 183–194; Aug., 1953.) The basic principles of a satisfactory "pay-as-you-view" system are considered. Essential requirements are (*a*) provision of a coin-actuated mechanism controlling a signal decoder, (*b*) program sales on a variable-price basis, the unit being a single programme rather than a unit of time, (*c*) announcement of details of the available program on the channel carrying the programme, for the benefit of anyone tuning in to that channel, and (*d*) an accurate record of each show purchased.

621.397.5 1198
Camera Adapter for TV Receivers—L. E. Flory, W. S. Pike, and G. W. Gray. (*Electronics*, vol. 27, pp. 141–143; Jan., 1954.) A com-

plete television system for domestic or business use comprises a vidicon camera unit, the picture signals from which are transmitted to the receiver on a vhf carrier via cable; operating voltages for the camera tube are derived from the receiver circuits.

621.397.5:535.623 1199
Second Color Television Issue—(Proc. I.R.E., vol. 42, pp. 1–348; Jan., 1954.) The whole of this issue is devoted to an account of developments in color television since 1951. The first group of papers comprises 22 contributions from the National Television System Committee; the titles of these papers, with the page references and the authors' names, are listed below. The second group comprises 27 contributions from individual workers on various branches of the subject; these papers are abstracted separately.
The Future of Color Television (pp. 5–7).—W. R. G. Baker.
Basic Concepts and Evolution of Color Television (pp. 7–9).—E. W. Engstrom.
Psychophysical and Electrical Foundations of Color Television (pp. 9–11).—A. V. Loughren.
The NTSC—An Exercise in Technical Coordination (pp. 11–14).—D. B. Smith.
NTSC Signal Specifications (pp. 17–19).
NTSC Color Television Field Test (pp. 20–43).—R. DeCola, R. E. Shelby, and K. McIlwain.
The NTSC Monographs (pp. 43–45).—D. G. Fink.
The NTSC Color Television Standards (pp. 46–48).
Colorimetry in Television: Parts 2 and 3 (pp. 48–57).—F. J. Bingley.
The Choice of Axes and Bandwidths for the Chrominance Signals in NTSC Color Television (pp. 58–59).—G. H. Brown.
The Constant Luminance Principle in NTSC Color Television (pp. 60–66).—W. F. Bailey
Mathematical Formulations of the NTSC Color Television Signal (pp. 66–71).—G. H. Brown.
Transfer Characteristics in NTSC Color Television (pp. 71–78).—F. J. Bingley.
Choice of Chrominance Subcarrier Frequency in the NTSC Standards (pp. 79–80).—I. C. Abrahams.
The "Frequency Interleaving" Principle in the NTSC Standards (pp. 81–83).—I. C. Abrahams.
Quandrature Cross Talk in NTSC Color Television (pp. 84–90).—W. F. Bailey and C. J. Hirsch.
Narrow-Band Transmission of the NTSC Color Signal (pp. 90–91).—J. G. Reddeck.
System Delay Characteristics in NTSC Color Television (pp. 92–95).—R. C. Palmer.
Effect of Transmitter Characteristics on NTSC Color Television Signals (pp. 95–105).—G. L. Fredendall and W. C. Morrison.
Color-Carrier Reference Phase Synchronization Accuracy in NTSC Color Television (pp. 106–133).—D. Richman.

621.397.5:535.623 1200
The Concept of Transmission Primaries in Color Television—P. W. Howells. (Proc. I.R.E., vol. 42, pp. 134–138; Jan., 1954.) The concept of color space is discussed. The location of the three transmission primaries of the NTSC system (namely, the monochrome signal and the two color signals) in the CIE color space is indicated.

621.397.5:535.623 1201
Colorimetric Analysis of the NTSC Color Television System—D. C. Livingston. (Proc. I.R.E., vol. 42, pp. 138–150; Jan., 1954.) Formulas are developed which enable color fidelity, compatibility, and adherence to the constant-luminance principle to be assessed. Three different forms of the NTSC system are compared; the superiority of the standard form is indicated. See also *Convention Record I.R.E.*, part 4, pp. 51–55; 1953.

621.397.5:535.623 1202
Quantitative Spectral Measurements in Color Television—J. A. Rado and W. L. Hughes. (Proc. I.R.E., vol. 42, pp. 151–156; Jan., 1954.) Measurements of the spectral sensitivities of the pickup devices and of the colorimetric characteristics of the display devices are discussed. Techniques are described using narrow-band optical-interference filters in conjunction with "red-pass" gelatin filters.

621.397.5:535.623 1203
Wide-Range Chromaticity Measurements with Photoelectric Colorimeter—J. B. Chatten. (Proc. I.R.E., vol. 42, pp. 156–160; Jan., 1954.) Detailed description of a tristimulus colorimeter sufficiently accurate for general work in color television.

621.397.5:535.623 1204
Reproduction of Colors in Outdoor Scenes—D. L. MacAdam. (Proc. I.R.E., vol. 42, pp. 166–174; Jan., 1954.) Available spectrophotometric and colorimetric data on the colors of skin, hair, grass, foliage, sky, and earth are summarized and measurements of typical color-film reproductions are presented. Subjective assessments of color-reproduction quality are discussed on the basis of these data.

621.397.5:535.623 1205
The Use of Electronic Masking in Color Television—R. P. Burr. (Proc. I.R.E., vol. 42, pp. 192–200; Jan., 1954.) A discussion of the use of electrical networks for reducing the effects of photographic cross-coupling introduced when using subtractive color transparencies as material for television transmission.

621.397.5:535.623 1206
Matrix Networks for Color TV—W. R. Feingold. (Proc. I.R.E., vol. 42, pp. 201–203; Jan., 1954.) Summation-type circuits required for color-television transmitters and receivers are described.

621.397.5:535.623 1207
Transients in Color Television—P. W. Howells. (Proc. I.R.E., vol. 42, pp. 212–220; Jan., 1954. *Convention Record I.R.E.*, part 4, pp. 24–34; 1953.) When a color transient occurs, the three components of the transmitted signal respond in a manner determined by the characteristics of the individual channels. The system response is characterized by the resulting path along which the reproduced color point moves through the color space from its initial to its final location. An analysis is made of such paths, and the corresponding subjective effects are discussed.

621.397.5:535.623 1208
Transition Effects in Compatible Color Television—J. B. Chatten. (Proc. I.R.E., vol. 42, pp. 221–228; Jan., 1954.) An experimental video-frequency television system is described for investigating the best choice of the three independent color-signal components to be transmitted. A description is given of an arrangement for predistorting the luminance signal to compensate for luminance variations arising in the reproduction of color transients.

621.397.5:535.623 1209
Reproduction of Luminance Detail by NTSC Color Television Systems—D. C. Livingston. (Proc. I.R.E., vol. 42, pp. 228–234; Jan., 1954.) Three different forms of the NTSC system are compared to determine the relative merits of different methods of gamma correction. The system using a luminance corrector in the transmitter is judged to be the best.

621.397.5:535.623 1210
Delay Equalization in Color Television—G. L. Fredendall. (Proc. I.R.E., vol. 42, pp. 258–262; Jan., 1954.) An equalizer designed by the potential-analog method consists of four conventional all-pass filter sections.

621.397.5:535.623.778.5 1211
Brightness Modification Proposals for Televising Color Film—W. L. Brewer, J. H. Ladd, and J. E. Pinney. (Proc. I.R.E., vol. 42, pp. 174–191; Jan., 1954.) It is recommended that (*a*) the effective luminance range should be modified, either in the film or in the television system, to conform to that available on the kinescope, and (*b*) kinescope reproductions of film color should be made lighter with increasing saturations. Alternative methods of achieving the color-dependent brightness compensation are discussed.

621.397.5:535.623].001.4 1212
Methods of Verifying Adherence to the NTSC Color Signal Specifications—Luther. (*See* 1150.)

621.397.5:535.623].001.4 1213
A Versatile Approach to the Measurement of Amplitude Distortion in Color Television—Bauer. (*See* 1151.)

621.397.5:535.623].001.4 1214
Test Instruments for Color Television—Morrison, Karstad, and Behrend. (*See* 1152.)

621.397.61:535.623 1215
Alignment of a Monochrome TV Transmitter for Broadcasting NTSC Color Signals—J. F. Fisher. (Proc. I.R.E., vol. 142, pp. 263–270; Jan., 1954.) A detailed account of the procedure used to adapt the WPTZ (Philadelphia) station.

621.397.611:535.623 1216
Image Orthicons for Color Cameras—R. G. Neuhauser, A. A. Rotow, and F. S. Veith. (Proc. I.R.E., vol. 42, pp. 161–165; Jan., 1954.) Operating conditions and performance characteristics of the R.C.A. Type-1854 are given.

621.397.611.2 1217
The Flying-Spot Scanning Systems by means of the MC 13-16—J. J. P. Valeton and F. H. J. van der Poel. (*Elec. Appl. Bull.*, vol. 14, pp. 77–90; June/July, 1953.) A detailed description is given of the flying-spot scanning system for deriving a picture signal for modulating a television transmitter from transparent positives or negatives. The MC 13-16 cr tube used has magnetic focusing and deflection systems. Circuit means are indicated for correcting the signal for the afterglow of the fluorescent screen.

621.397.62+621.396.621 1218
Radio and Television at the 16th Salon National—François. (*See* 1171.)

621.397.62 1219
Technical Description of a Television Receiver—A. Bilotti. (*Rev. Telegr. Electronica*, (*Buenos Aires*), no. 492, pp. 555–560; Sept., 1953.) Detailed circuit diagram and component values are given for a receiver operating on the intercarrier system, to suit the Argentina television standards, namely, negative picture signal, line frequency 15,625 per second, frame frequency 50 per second, aspect ratio 3:4, FM sound, separation of 4.5 mc between picture and sound carriers, channel width 6 mc. Reception on all twleve vhf channels is possible.

621.397.62 1220
The D.C. Component in Television—W. T. Cocking. (*Wireless World*, vol. 60, pp. 63–66; Feb., 1954.) Birkinshaw's views (552 of February) are generally supported, but criticism is made of the well known circuit in which the cathode of the cr tube is fed from the anode of the video stage through a voltage divider. The low input resistance of the tube thus connected may lead to a 50 per cent reduction of the dc component. Various modifications of the coupling circuit are suggested for eliminating this defect.

621.397.62 1221
A Printed-Circuit Television Receiver—

G. Székely. (*Télévision*, no. 37, pp. 230–232; Oct., 1953.) Brief description of a new French receiver. All circuits, excepting time bases and supply, but including deflection coils, are printed.

621.397.62:535.623 1222
Improving the Transient Response of Television Receivers—J. Avins, B. Harris, and J. S. Horvath. (PROC. I.R.E., vol. 42, pp. 274–284; Jan., 1954.) An examination is made of the delay distortion produced in the IF amplifier, and of the extent to which this is compensated by the peaking circuits of the video detector and amplifier. Factors determining cross-talk between the two color-difference signals in the NTSC system are discussed. The use of a standard transmission monitor throughout the industry is recommended.

621.397.62:535.623 1223
Theory of Synchronous Demodulator as used in NTSC Color Television Receiver—D. C. Livingston. (PROC. I.R.E., vol. 42, pp. 284–287; Jan., 1954.) See 3129 of 1953.

621.397.62:535.623 1224
The D.C. Quadricorrelator: a Two-Mode Synchronization System—D. Richman. (PROC. I.R.E., vol. 42, pp. 288–299; Jan., 1954.) A description is given of an automatic-phase-control circuit adapted for color synchronization in the NTSC system. The arrangement is simple, reliable, and noise-immune. The pull-in and hold-in modes of operation are made independent of each other by means of an automatic switch. See also *Convention Record I.R.E.*, part 4, pp. 13–23; 1953.

621.397.62:535.623 1225
Processing of the NTSC Color Signal for One-Gun Sequential Color Displays—B. D. Loughlin. (PROC. I.R.E., vol. 42, pp. 299–308; Jan., 1954.) Receiver arrangements of two types are described, namely those in which the color subcarrier signal is modified, and those in which the luminance signal is modified. Circuits of the latter type generally demodulate the color signal to some extent and correct the luminance signal to produce constant luminance.

621.397.62:535.623 1226
Compatible Color Picture Presentation with the Single Gun Tricolor Chromatron—J. D. Gow and R. Dorr. (PROC. I.R.E., vol. 42, pp. 308–315; Jan., 1954.) Techniques are discussed for operating the tube previously described [3131 of 1953 (Dressler)]. Circuit nonlinearity is corrected so as to ensure color balance at different brightness levels.

621.397.62:535.623 1227
Improvements in the R.C.A. Three-Beam Shadow-Mask Color Kinescope—M. J. Grimes, A. C. Grimm, and J. F. Wilhelm. (PROC. I.R.E., vol. 42, pp. 315–326; Jan., 1954.) Advances made since the original design [844 of 1952 (Law)] are described.

621.397.62:535.623 1228
The C.B.S.-Colortron: a Color Picture Tube of Advanced Design—N. F. Fyler, W. E. Rowe, and C. W. Cain. (PROC. I.R.E., vol. 42, pp. 326–334; Jan., 1954.) Full details are given of a three-gun tricolor tube in which the phosphor dots are applied directly to the curved endplate, and a self-supporting curved mask is provided. See also *Tele-Tech.*, vol. 12, pp. 73, 150, 165; Nov., 1953.

621.397.62:535.623 1229
A Laboratory Receiver for Study of the NTSC Color Television System—C. Masucci, J. J. Insalaco, and R. Zitta. (PROC. I.R.E., vol. 42, pp. 334–343; Jan., 1954.)

621.397.62:621.314.632 1230
Some High Frequency Effects in Germanium Diodes with Special Reference to Television Sound Detectors—D. D. Jones and B. C. Brodribb. (*Electronic Eng.*, vol. 26, pp. 33–35; Jan., 1954.) The measured performance of Ge diodes at HF may differ appreciably from that predicted from the static characteristics. The deviations, which may give rise to highly nonlinear detection, are shown to be consistent with hole-storage effects in the Ge; these effects can be eliminated by production techniques designed to reduce the lifetime of the holes. An indication is given of suitable test methods.

621.397.62:621.375.2 1231
Video Amplifier Design Charts—W. K. Squires and H. L. Newman. (*Electronics*, vol. 27, pp. 190–192, 194; Jan., 1954.) Charts based on transient-response analysis are presented.

621.397.8 1232
Image Gradation, Graininess and Sharpness in Television and Motion-Picture Systems: Part 2—The Grain Structure of Motion Picture Images—An Analysis of Deviations and Fluctuations of the Sample Number—O. H. Schade. (*Jour. Soc. Mot. Pict. & Telev. Engrs.*, vol. 58, pp. 181–222; March, 1952.) A treatment of aperture-response theory as applied to the evaluation of relative deviation in motion-picture processes. Part 1: 2293 of 1951.

621.397.8 1233
Image Gradation, Graininess and Sharpness in Television and Motion-Picture Systems: Part 3—The Grain Structure of Television Images—O. H. Schade. (*Jour. Soc. Mot. Pict. & Telev. Engrs.*, vol. 61, pp. 97–164; Aug., 1953.) An adequate description of granularity in television and motion-picture frames requires specification of the sine-wave spectrum and signal/deviation ratio in the retinal image as a function of luminance and for a specified viewing distance. An assessment of the perception of deviations can be made by introducing the characteristic of threshold signal/deviation ratios as the reference level. Part 2: 1232 above.

621.397.8 1234
Note on the Influence of Phase in Television—A. Dubec. (*Onde élect.*, vol. 33, pp. 606–614; Nov., 1953.) Methods of calculating phase /frequency characteristics are reviewed in relation to the design of television receivers.

621.397.8 1235
Eye Movements in Connexion with Television Viewing—M. P. Lord. (*Nature, (London)*, vol. 172, pp. 964–965; Nov. 21, 1953.) A comparison of movements made when viewing interlaced and sequentially scanned pictures indicates little difference; hence instability and "crawling" associated with interlaced scan cannot be immediately explained in terms of the pattern of eye movements.

TRANSMISSION

621.375.227.029.4:621.3.018.783 1236
Nonlinear Distortion in a Class-B Push-Pull Amplifier with Transformer Output—F. Böttcher. (*Telefunken Ztg.*, vol. 26, pp. 313–322; Aug., 1953.) Expressions for the distortion, up to the 4th harmonic, are derived for the case of imperfect magnetic coupling between the halves of the primary winding of the output transformer. An approximate numerical calculation is made for the modulator output stage of an anode-modulated transmitter.

621.396.664 1237
Speech Clippers and their Operation—W. Schreuer. (*Short Wave Mag.*, vol. 11, pp. 465–469; Oct., 1953.) Average modulation depth is increased without introducing overmodulation on voice peaks by using an amplitude limiter followed by a low-pass filter, with a cut-off frequency ~3kc. High-level and low-level systems are discussed.

TUBES AND THERMIONICS

621.314.63:546.289 1238
Recovery Currents in Germanium *p-n* Junction Diodes—R. G. Shulman and M. E. McMahon. (*Jour. Appl. Phys.*, vol. 24, pp. 1267–1272; Oct., 1953.) When a diode biased in the forward direction has a reverse voltage suddenly applied, a large transient ("recovery") current flows. Values of this current based on a proposed mechanism involving diffusion of stored minority carriers to the barrier are compared with experimental results. The application to variable time-delay devices is discussed.

621.314.632:546.289 1239
Theory of the Forward Characteristic of Injecting Point Contacts—P. C. Banbury. (*Proc. Phys. Soc.*, vol. 66, pp. 833–840; Oct. 1, 1953.) curacy of 0.01 per cent uses cold-cathode glow-The voltage/current relation for forward voltages across a metal/semiconductor point contact is derived for the case when the current is carried by injected minority carriers. Unit injection ratio is assumed and recombination effects are neglected. Results of measurements on contacts under illumination and experimentally determined forward-voltage/current characteristics of a Ge-W rectifier are in satisfactory agreement with theory.

621.314.632:546.289 1240
Forward Characteristic of Injecting Area Contacts on Germanium—H. K. Henisch and F. D. Morten. (*Proc. Phys. Soc.*, vol. 66, pp. 841–844; Oct. 1, 1953.) Measurements were made on 4.2-Ω. cm *n*-type and on 12-Ω. cm *p*-type Ge illuminated with white light. Au or Ag rectifying electrodes were used. The results found for the floating potential, the contact capacitance and the derived current/voltage barrier characteristic are given and compared with those derived from theory.

621.314.632:546.289:537.312.6 1241
Thermal Effects at Point-Contact Diodes—P. M. Tipple and H. K. Henisch. (*Proc. Phys. Soc.*, vol. 66, pp. 826–832; Oct. 1, 1953.) Measurements, by a thermoelectric method, have been made of the contact temperature of a Ge diode, at several points on the voltage/current characteristics. The voltage turnover occurs at a critical contact temperature, which is constant for a given specimen. Also, the temperature gradients in the vicinity of a hot point contact may increase the contact resistance. This observation is discussed.

621.314.632:546.289:537.312.6 1242
Thermal Effects in Point-Contact Rectifiers—H. L. Armstrong. (*Jour. Appl. Phys.*, vol. 24, pp. 1332–1333; Oct., 1953.) The relation between thermal effects and the shape of the reverse characteristic is investigated for a Type-1N34 Ge rectifier.

621.314.632:621.397.62 1243
Some High Frequency Effects in Germanium Diodes with Special Reference to Television Sound Detectors—Jones and Brodribb. (*See* 1230.)

621.314.7 1244
Design Theory of Junction Transistors—J. M. Early. (*Bell Sys. Tech. Jour.*, vol. 32, pp. 1271–1312; Nov., 1953.) "The small signal ac transmission characteristics of junction transistors are derived from physical structure and bias conditions. Effects of minority carrier flow and of depletion layer capacitances are analyzed for a one dimensional model. The ohmic spreading resistance of the base region of a three dimensional model is then approximated. Short circuit admittances representing minority carrier flow, depletion layer capacitances, and ohmic base resistance elements are then combined into an equivalent circuit. Theoretical calculations are compared to observations for two typical designs."

621.314.7:[546.28+546.289 1245

A Study of Carrier Injecting Properties of Emitter Contacts and Light Spots at Normal and Moderately Elevated Temperatures—C. A. Hogarth. (*Proc. Phys. Soc.*, vol. 66, pp. 845–858; Oct. 1, 1953.) "The method of determining the injection efficiency γ of an emitter contact (described previously by Shockley, Pearson and Haynes [380 of 1950]) is discussed and some of the difficulties of the method and the interpretation of experimental results are described. The effects of a fine spot of white light when used as an emitter and hence as a conductivity modulator are investigated and an equivalent minority carrier current for a given optical assembly can be determined. This procedure suggests an experimental method for the determination of γ as a function of temperature. A simple theory for γ in terms of forbidden energy gap, depth of Fermi level, and barrier height ϕ is given and applied to the determination of barrier heights at room temperature, at elevated temperatures, and when an emitter contact on germanium is illuminated. Theoretical curves relating γ and ϕ are presented for Ge and Si of various impurity concentrations. The experiments carried out at elevated temperatures suggest that the surface states on germanium lie at the top of the filled band and that their density is of order 10^{10}–10^{11} cm^{-2}."

621.314.7:546.28 1246

Forming Silicon Point-Contact Transistors—H. Jacobs, F. A. Brand, and W. Matthei. (*Jour. Appl. Phys.*, vol. 24, p. 1340; Oct., 1953.) Particles of a suitable impurity material were pressed between a pointed tungsten electrode and the Si sample, and a high-density current was passed through the junction in a N_2 atmosphere. Best results were obtained using Sb to form *p-n-p* transistors and Al to form *n-p-n* transistors. Some test results are tabulated.

621.314.7:621.375.4.029.3 1247

Power Transistors for Audio Output Circuits—L. J. Giacoletto. (*Electronics*, vol. 27, pp. 144–148; Jan., 1954.) Junction transistors cooled by liquids and by metal radiators are described. Analysis previously developed for small-signal operation is adapted to power operation, and the influence of finite base-lead resistance, temperature, frequency and source resistance is examined. Biasing arrangements for typical audio output circuits are discussed.

621.314.7:621.396.822 1248

Shot Noise in Junction Transistors—H. C. Montgomery and M. A. Clark. (*Jour. Appl. Phys.*, vol. 24, pp. 1337–1338; Oct., 1953.) The noise figure calculated from Montgomery's formula (643 of 1953) is in fair agreement with the values measured in a *p-n-p* alloy transistor at frequencies above those for which the noise is inversely proportional to frequency.

621.383.2:621.317 1249

The Photodianode—G. Deloffre, É. Pierre, and J. Roig. (*Compt. Rend. Acad. Sci.*, (*Paris*), vol. 237, pp. 1507–1509; Dec. 9, 1953.) The "photodianode" is a tube with a plane photoemissive cathode and two linear anodes arranged parallel to and equidistant from the cathode. The anodes are connected via separate equal resistors to the positive supply terminal, thus there is a potential difference between them depending on the position of the illuminated area; this potential difference is determined by measuring the current through a galvanometer connected between the anodes. The arrangement responds to variations of the direction of a light beam, and can be used for accurate measurements of displacement and rotation.

621.383.4 1250

On the Relation between the Speed of Response and the Detectivity of Lead Sulfide Photoconductive Cells—R. C. Jones. (*Jour. Opt. Soc. Amer.*, vol. 43, pp. 1008–1013; Nov., 1953.)

621.383.5:546.23 1251

A New Type of Fatigue of Photocells—G. Blet. (*Compt. Rend. Acad. Sci.* (*Paris*), vol. 238, pp. 228–230; Jan. 11, 1954.) The impedance of a Se barrier-layer cell was found to drop immediately on illumination, returning to its original value only very slowly when the illumination was cut off. The suggested explanation is that the liberated photoelectrons suffer collisions which reduce their energy and hence delay their return to their original positions.

621.385:681.142 1252

Valve Reliability in Digital Calculating Machines—Knight. (*See* 984.)

621.385.029.6 1253

Instabilities in the Smooth-Anode Cylindrical Magnetron—L. A. Harris. (*Jour. Appl. Phys.*, vol. 24, p. 1335; Oct., 1953.) Correction to paper noted in 3293 of 1952. An error in equation (38) invalidates the previous conclusions. The system is stable, with a set of real characteristic frequencies.

621.385.032.213 1254

The Motion of Idealized Vacuum Tube Filaments under Shock—H. S. Thomas. (*Jour. Appl. Phys.*, vol. 24, pp. 1341–1342; Nov., 1953.) An analysis is made of the motion of an ideal clamped-end string; approximate formulas are derived for the maximum displacement for accelerations of the supports corresponding to three different types of shock. The results are applicable to the investigation of filament-to-grid shorting.

621.385.032.216 1255

Molded Thermionic Cathodes—D. MacNair, R. T. Lynch, and N. B. Hannay. (*Jour. Appl. Phys.*, vol. 24, pp. 1335–1336; Oct., 1953.) The cathodes are made of sintered mixtures of Ni powder and (Ba, Sr)CO_3 on an Ni base. Test results indicate a very high resistance to deactivation by gases and ion bombardment. A comparison with sprayed oxide and other cathodes is made.

621.385.032.216 1256

Deterioration of Valve Performance due to Growth of Interface Resistance—G. H. Metson and M. F. Holmes. (*P.O. Elect. Eng. Jour.*, vol. 46, part 4, pp. 198–199; Jan., 1954.) Results of measurements on a high-slope pentode indicate that the development of interface resistance between core and coating of the oxide cathode [2378 of 1952 (Child)] can be eliminated by replacing the usual Ni core by a material free from Si, e.g. commercially pure Pt. The lifetime of common receiving tubes might be doubled.

621.385.032.216:546.831 1257

Thermionic Emission and Optical Emissivity of Zirconium—G. Mesnard. (*I e Vide*, vol. 8, pp. 1392–1399; July/Sept., 1953.) Measurements were made of the properties of cathodes prepared by coating W or Mo wires with zirconium oxide. The activation by thermal treatment and the development of emission with time were studied. Spectral emissivity, resistance and required heating power are discussed. Similarities are noted between these cathodes and those using other materials, particularly thoria.

621.385.032.216:546.841.4-31 1258

Thermionic Properties of Thoria Coatings on Thoriated Molybdenum—G. Mesnard. (*Le Vide*, vol. 8, pp. 1377–1383; July/Sept., 1953.) Measurements of the emission from (*a*) thoriated molybdenum, (*b*) thoria coatings on nonthoriated molybdenum, and (*C*) thoria coatings on thoriated molybdenum indicated that no appreciable advantage is obtained by thoriating the molybdenum base. The effects of treatment at high and at medium temperatures are discussed.

621.385.032.24 1259

Gold as a Grid Emission Inhibitor in the Presence of an Oxide-coated Cathode—B. O. Baker. (*Brit. Jour. Appl. Phys.*, vol. 4, pp. 311–315; Oct., 1953.) "Emission measurements have been made on gold-plated molybdenum and gold-plated manganese-nickel grids in the presence of an oxide-coated cathode. For grids which cannot be designed to operate below 350 degrees C. a minimum thickness of 1 μ of gold will suppress grid emission. Silver is not so reliable, but is effective in some cases."

621.385.2:621.316.722.1 1260

Characteristics of the Temperature-Limited Diode Type 29C1—V. H. Attree, F. A. Benson and M. S. Seaman. (*Electronic Eng.*, vol. 26, p. 42; Jan., 1954.) Comment on 595 of February and authors' reply.

621.385.3:621.396.822 1261

Experimental Investigation of Grid Noise—N. Houlding and A. E. Glennie. (*Wireless Eng.*, vol. 31, pp. 35–42; Feb., 1954.) "Results are given of a detailed investigation of triode noise factor, with particular reference to correlation of induced grid noise with shot noise. It is deduced that correlation is very slight and that, although the optimum value of noise factor can be calculated fairly accurately from the values of shot noise and optimum source resistance, the latter must be found by experiment and therefore the theory is not of major practical importance."

621.385.832 1262

Theory of the Triode Electron Gun with Cylindrical Symmetry—É. Labin. (*Onde élect.*, vol. 33, pp. 597–605 and 713–719; Nov./Dec., 1953.) Systems with nonrotational symmetry are considered, such as those used for producing ribbon beams. The method of conformal representation is used and formulae are derived based on purely electrostatic conditions; the validity of these formulas for operating conditions is examined.

621.385.832:681.142 1263

Recent Advances in Cathode-Ray-Tube Storage—Williams, Kilburn, Litting, Edwards, and Hoffman. (*See* 985.)

621.378 1264

Investigation of the use of Caesium Vapour in Thyratron-Type Gas-Filled Valves—R. Coquerel. (*Le Vide*, vol. 8, pp. 1384–1391; July/Sept., 1953.)

621.387:621.318.57 1265

A New Cold-Cathode Decade Counter Tube—H. v. Gugelberg. (*Helv. Phys. Acta*, vol. 26, pp. 586–588; Nov. 16, 1953. In German.) Counting rates up to 10^5 pulses/second have been obtained with experimental tubes using asymmetrical molybdenum cathodes.

621.387:621.318.572 1266

Cold-Cathode Tubes for Transmission of Audio-Frequency Signals—M. A. Townsend and W. A. Depp. (*Bell. Sys. Tech. Jour.*, vol. 32, pp. 1371–1391; Nov., 1953.) The desirable transmission properties are discussed of tubes for use as switching elements in series with telephone circuits; performance curves for an experimental diode are given.

MISCELLANEOUS

621.39 1267

Some Recent Developments in Electronic Engineering—(*Electronic Eng.*, vol. 26, pp. 20–26; Jan., 1954.) Developments discussed include telemetering in guided missiles, radar, and navigation aids and associated equipment for radio communication, equipment for meteorological investigations, computers, and apparatus for electromedical and industrial uses.

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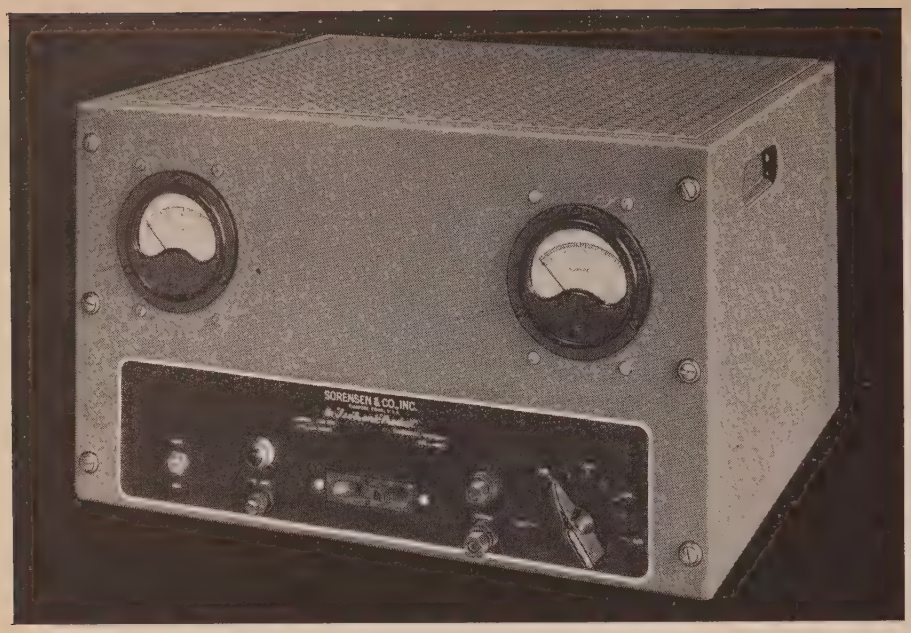

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Output voltage	115VAC, 1Ø, adjustable between 110-120 volts		
Output frequency	400～, adjustable ±10%	400～, adjustable ±10%	60～, adjustable between 45 and 65
Output voltage regulation	±1.0%	±1.0%	±1.0%
Output frequency regulation	±1.0% in standard models; ±0.01% with auxiliary frequency standard (output frequency is fixed when using frequency standard)		
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Load range	0.1 to full load		
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P. F. range	Down to 0.7 F		
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Envelope modulation	2% maximum		

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Industrial Engineering Notes

TECHNICAL*

A new radio-operated remote control set which makes practicable the operation of a radio, radio teletypewriter, or radio telephone miles from the sending unit, has been developed by a Signal Corps noncommissioned officer, it was recently announced. The equipment is considered so promising that it has been submitted to the Signal Corps Engineering Laboratories for further tests, and, perhaps, general adoption by the Army, it was reported. The unit, about the size of a cigar box, is inexpensive to manufacture according to members of the 301st Signal Group, Seventh U. S. Army in Europe. Used in conjunction with the radio AN/GRC-26A, it reportedly can control a set two miles away by the "one-way reversible" and "duplex operational" methods as easily as one can make adjustments on radio or television set just a few feet away. The unit is capable of selecting combinations to alternately send and receive at the same time and can cut teletype tape as easily as sitting at the machine itself. In addition, radio telephone and telegraph are available by a simple flick of a switch. The remote control unit, the Signal Corps said, eliminates the use of an operator at the radio site and frees manpower that heretofore was needed to lay wire from the communications center to the site of the equipment. The device also saves the cost of wire that normally is involved in a metallic circuit from the communications center to the equipment being operated at a distance. **. . . The National Bureau of Standards has announced the development of an altimeter which will measure altitudes as low as two feet, attained by using appropriate frequency shifting and by modifying existing techniques for recognizing short-distance information.** The instrument, called the nonquantized frequency-modulated altimeter, was designed by H. P. Kalmus, J. C. Cacheris, and H. A. Dropkin as part of the Bureau's research in ordnance electronics. Since, heretofore, conventional altimeters gave information only to the closest 10 or 20 feet, the new NBS altimeter makes it possible for a helicopter pilot to know when his craft is within several feet of the landing surface, it was reported. This is accomplished through an operating frequency in the X-band of the radio spectrum (10,000 mc). **. . . The Office of Technical Services, Department of Commerce, in its February issue of the "Bibliography of Technical Reports," lists a number of important advances in electronics.** The following government-sponsored research studies can be purchased from the Photoduplication Section, Library of Congress, Washington 25, D. C., for the reported price: "Radio Noise Consideration," PB 112552, microfilm, $2.75, photostat, $7.50; "Cooling De-

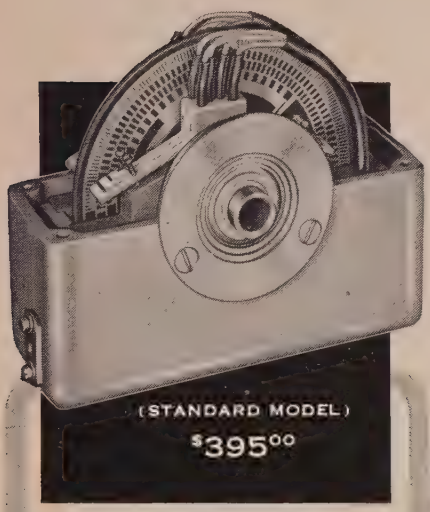
(Continued on page 84A)

* The data on which these NOTES are based were selected by permission from *Industry Reports*, issues of February 22, March 2 and 8, published by the Radio-Electronics-Television Manufacturers Association, whose helpfulness is gratefully acknowledged.

Giannini

DIGITAL COMMUTATORS

Codes: Binary, Binary-Grey, Binary-Decimal with non-ambiguous outputs



Immediate Delivery from stock on standard models having following characteristics:

- 1,000 counts decimal (300° to 360°)
- 1024 counts—binary, grey (300° to 360°)
- torque—0.5 oz. in ball bearings
- inertia—400 gm. cm^2
- micrometer zero adjustment
- automatic alignment—no gears

CUSTOM COMMUTATORS

Commutators can be furnished to fit specific applications by either modification of standard models or wholly new designs. Some variations now available are:

- non-linear coding
- high-count multi-turn units
- ultra low torque models
- miniature size, geared units
- direct decimal coding

G. M. GIANNINI & CO., INC.
LABORATORY APPARATUS DIVISION
PASADENA 1, CALIFORNIA

Offices: New York: Phone Judson 6-7500
Los Angeles: Phone Ryan 1-7152

Giannini

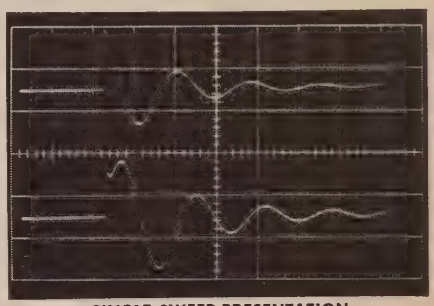
Dual-Trace Applications

WITH THE TEKTRONIX TYPE 535 OSCILLOSCOPE AND TYPE 53C DUAL-TRACE PLUG-IN PREAMPLIFIER

Here is a combination ideally suited to most applications involving accurate comparisons of two signals.

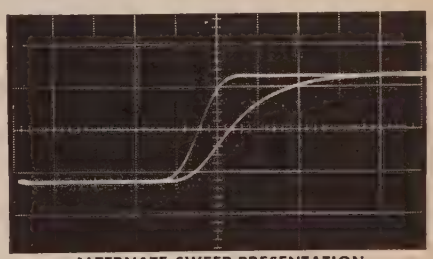
The Type 53C Dual-Trace Unit contains two identical amplifier channels that can be electronically switched either by the oscilloscope sweep or at a free-running rate of approximately 100 kc. When amplifier switching is triggered by the sweep, the two signals to be compared appear on alternate sweeps. Because the sweeps are identical, and time-delay characteristics of the two amplifier channels are closely controlled, time comparisons accurate within 1 mμsec can be made. Two simultaneous transients may be viewed by free-running the switching. Transients of as little as 1 msec duration are well delineated, having about 100 elements in each trace. For many purposes, shorter transients can be adequately observed.

The Type 535 Oscilloscope is designed to use plug-in preamplifiers. It has an exceptionally wide sweep range, high accelerating potential, new accurate sweep-delay circuitry, and many other important features. Four Plug-In Preamplifiers have been developed for use with the Type 535, to provide an unusually high degree of flexibility in a single oscilloscope.



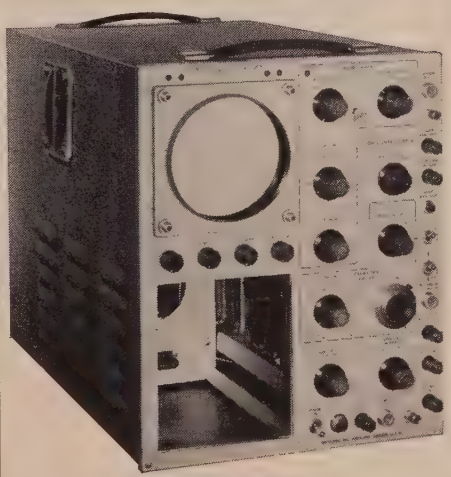
SINGLE-SWEEP PRESENTATION

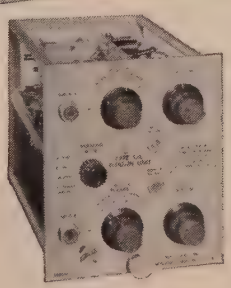
Response of two networks excited by a single pulse shows free-running operation of the Dual-Trace Unit in a one-shot application. A single 200-μsec/cm sweep is used for this display.



ALTERNATE-SWEEP PRESENTATION

Output of an RC network superimposed on the input pulse. Both waveforms appear on alternate 0.04 μsec/cm sweeps, accurately measuring the risetime deterioration caused by passage through the network.





MAIN OSCILLOSCOPE FEATURES

600,000,000 to 1 Sweep Range—0.02 μsec/cm to 12 sec/cm, continuously variable. Calibrated—0.02 μsec/cm to 5 sec/cm, accurate within 3%.

10 KV Accelerating Potential—Brighter display at low repetition rates.

Flexible Sweep Delay—1 μsec to 0.1 sec, jitter-free, incremental accuracy within 0.2% of full scale.

Type 535 Oscilloscope—$1300 plus price of desired plug-in units.

DUAL-TRACE PLUG-IN PREAMPLIFIER

Type 53C Specifications

Two Identical Amplifier Channels
Frequency Response—DC to 8.5 mc.
Risetime—0.04 μsec.
Sensitivity—0.05 v/cm to 20 v/cm calibrated, continuously variable to 50 v/cm.

Electronic Switching
Triggered—actuates alternate sweeps.
Free-running rate—100 kc, approximately.

Type 53C Dual-Trace Unit—$275.

OTHER PLUG-IN PREAMPLIFIERS

Type 53A Wide-Band DC Unit—$85.
Type 53B Wide-Band High-Gain Unit—$125.
Type 53D High-Gain Differential Unit—$145.

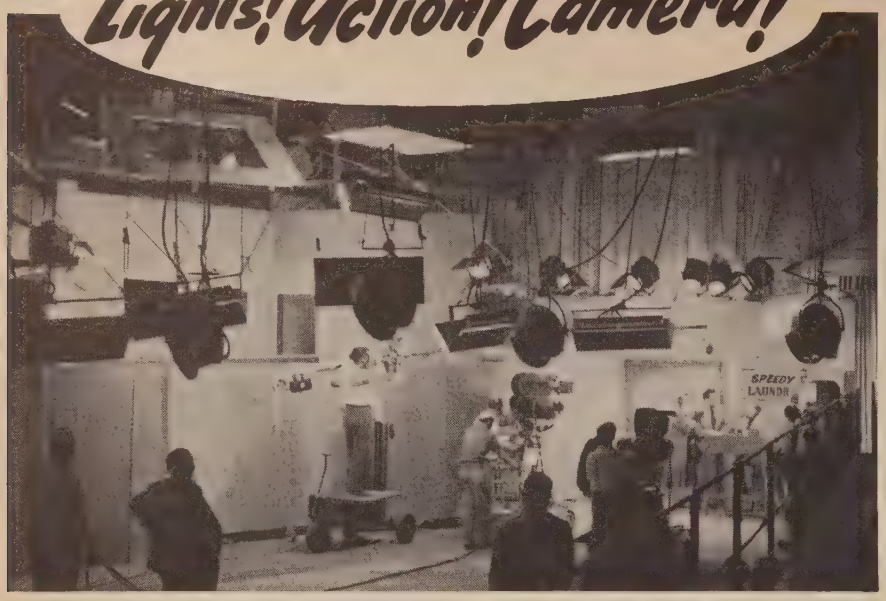
Prices f.o.b. Portland (Beaverton), Oregon.

Please write for complete specifications

Tektronix, Inc.

P. O. Box 831B • Portland 7, Oregon • Cable: TEKTRONIX

Lights! Action! Camera!

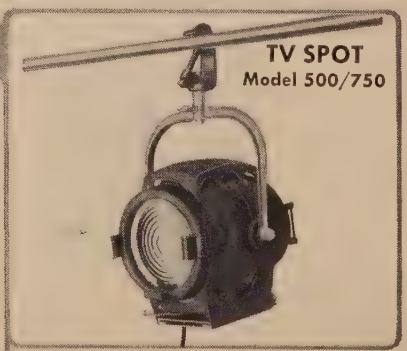


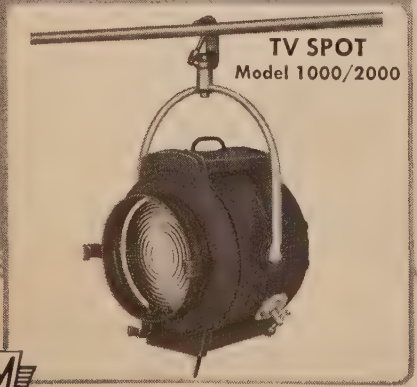
Typical Installation of B & M Spotlights in "I Love Lucy" TV Show

"Paint With Light" with

BARDWELL & McALISTER'S Engineered Television Lights TV SPOTS

This complete line of television spotlights was especially designed and engineered for television stage lighting. The many years of experience in the production of studio lights for stage and motion pictures was drawn upon to produce the B & M TV Spots which make possible painting with light. Painting with light is the ability to control the light source in order to emphasize the necessary highlights and all the important shadows. Only through controlled light can the scene and subject be given the desired brilliance, beauty, and third dimensional effects to produce ideal screen pictures.





Our specialists are always ready to assist and advise your engineering staff so that your studios and stages will be fully equipped to properly paint with light. Write for complete specifications and prices of these B & M TV Spots.

BARDWELL & McALISTER, Inc. 2950 ONTARIO STREET BURBANK, CALIFORNIA

Industrial Engineering Notes

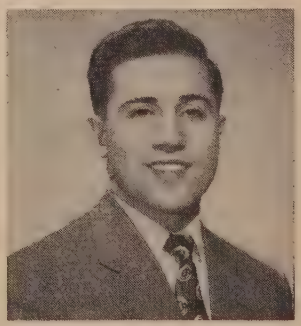
(Continued from page 82A)

sign of Transceiver Unit AN/ARC-21 Radio Equipment," PB 112527, microfilm, $2.25; photostat, $5; "Electron Impact Study of Nitrogen in the Kinetic Energy Range 400 to 600 Volts," PB 112474, microfilm, $3, photostat, $8.75; "Diffraction of Microwaves by Slits and Elliptical Apertures," PB 112421, microfilm, $2.50; photostat, $6.25" "Free Convective Cooling in Air of Confined Small Bodies Similar to Electronic Components," PB 112402, microfilm, $2.75, photostat, $7.50; "Generation of Millimeter Electromagnetic Waves—Progress Report," PB 112364, microfilm, $1.25, photostat, $1.25. "Induction Method of Measuring Electrical Resistivity," PB 112524, microfilm, $2, photostat, $3.75" "Microwave Components," PB 112451, microfilm, $2.25, photostat, $5; "Microwave Noise Study—Quarterly Report," PB 112376, microfilm, $2.75, photostat, $7.50; "Scattering and Attenuation of Microwave Signals by Hydrometeors," PB 112565, microfilm, $6.75, photostat, $23.75; "Study of Heat Dissipation from Electronic Devices at High Altitudes," PB 112625, microfilm, $1.25, photostat, $1.25; "Signal Corps Electronic Computer Research and Development," PB 112629, microfilm, $2.25, photostat, $5. . . . **The Civil Aeronautics Administration has announced publication of "Operation of the Air Traffic Control System."** The new booklet tells how the federal airways work, describing the communications network and electronic equipment that provides navigational aid enroute and at terminals. The new booklet may be obtained from the Superintendent of Documents, Government Printing Office, Washington 25, D. C., at 25 cents per copy. . . . **The Radio Technical Commission for Aeronautics announced recently that it has developed procedures for calibrating audio- and radio-frequency signal generators to produce signals acceptable for testing and adjusting airborne VOR and ILS receivers.** The procedures

(Continued on page 86A)

Waveline Appoints Del Vento

John M. Del Vento has been appointed Chief Engineer of **Waveline, Inc.**, Caldwell, N. J.



Prior to joining the company he served in various capacities in the field of microwave research and development.

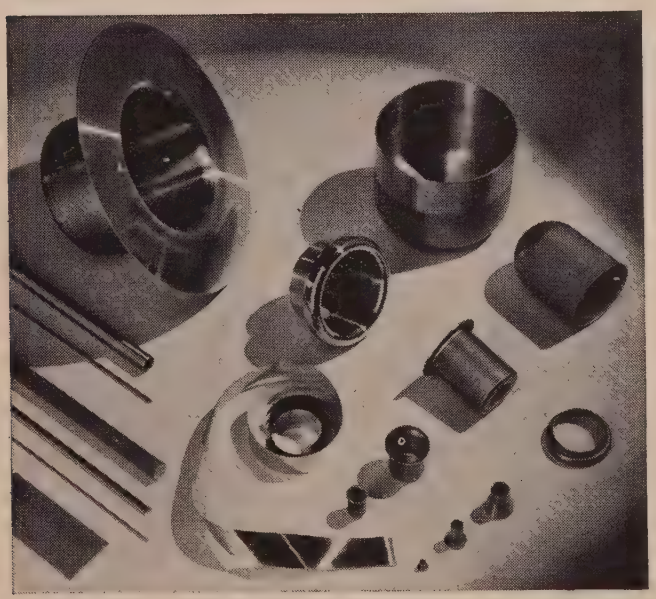
Years of experience

proves to users . . .

the dependability of

KOVAR*

Glass-sealing Alloy



* Westinghouse Trade Mark No. 337,962

The ideal alloy for glass sealing, Kovar matches the expansivity of certain hard glasses over the entire working temperature range. It resists mercury attack, has ample mechanical strength and seals readily. A permanent and impervious bond is obtained by a closely controlled thickness of oxide on Kovar alloy interfused with hard glass.

Kovar is a cobalt, nickel, iron alloy, manufactured under very carefully controlled conditions, and supplied by Stupakoff in the form of: SHEET, ROD, WIRE, FOIL, TUBING, EYELETS, LEADS and FABRICATED SHAPES. The prominent users of KOVAR and the length of time they have employed this metal are convincing proof of satisfaction.

Full information on the use of Kovar is given in Stupakoff Bulletin 145, which we will send upon request.

Stupakoff CERAMIC & MFG. CO.

Latrobe, Pennsylvania

18 YEARS

18 YEARS

Western Electric

16 YEARS

AMPEREX

12 YEARS

RCA

18 YEARS

Eimac

TUBES

13 YEARS

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11 YEARS

MACHLETT

15 YEARS

Bomac

6 YEARS

WESTINGHOUSE

18 YEARS

VARIAN associates

5 YEARS

RAYTHEON

18 YEARS

7 YEARS

SPERRY

13 YEARS

Announcing
PRECISION D-C VOLTMETER
Model 124

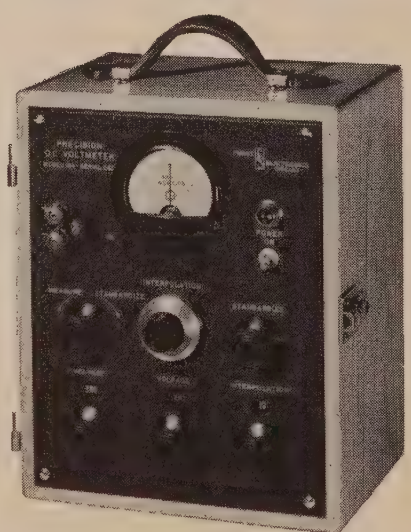
The Model 124 Precision D-C Voltmeter produces an accurately adjustable reference voltage for comparison with the voltage to be measured. A null indicating meter is used to indicate equality of the two voltages. Voltages between 0 and 510 Volts can be measured. A very stable regulated power supply circuit is used as the internal voltage source. It can be standardized against a built-in standard cell by a switch and control on the front panel. The switch also controls the sensitivity of the null indicator when measurements are made.

Two sensitivity ranges are provided and are selected by a switch on the front panel.

On special order, the terminals of the reference voltage may be connected to a suitable receptacle on the front panel bypassing the null indicator.

Also available for relay rack mounting.

Voltmeters for other voltage ranges supplied on special order.



Specifications:

- **VOLTAGE RANGE:** 0 to 510 Volts in steps of 10 Volts or 0 to 500 Volts in steps of 100 Volts; each step is subdivided by a vernier dial reading 10 Millivolts or 100 Millivolts, resp., per division.
- **ACCURACY:** When properly standardized, voltage indications are accurate within 0.1% in the "10 Volt" position of the range switch. Small voltage differences can be measured with an accuracy of better than 0.1% or 5 Millivolts, whichever is greater.
- **INPUT IMPEDANCE:** Infinite, after null balance is obtained.
- **DIMENSIONS:** 9½" wide by 12" high by 8" deep, excluding carrying handle and rubber feet. Front panel: 8½" wide by 11" high.

FURST ELECTRONICS
3326 W. Lawrence Ave., Chicago 25, Illinois

Industrial Engineering Notes

(Continued from page 84A)

were developed primarily to aid the operators of aircraft radio service stations. They were prepared for the more commonly-used types of signal generators, RTCA said, but also may be used to determine whether signals produced by other types are within acceptable tolerances. The calibration study was performed by RTCA Special Committee 61. Copies of the 80-page report (Paper 208-53/DO-52) may be obtained from the RTCA Secretariat, 1724 F Street, N.W., Washington 25, D. C. In addition to signal generators, standard omni-bearing selectors are used in testing VOR receivers at bearing settings other than the zero indexing position. Calibration procedures for omni-bearing selectors are being prepared for later publication, RTCA reported.

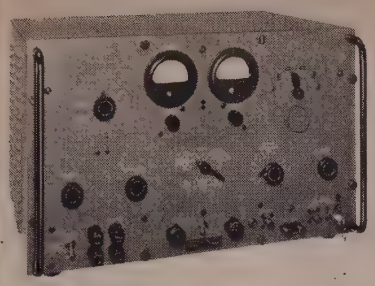
Industry Statistics

Over 7.2 million radios, excluding automobile sets, were shipped to dealers during 1953, according to a report issued recently from the RETMA Statistical Department. This was about two per cent above the level of shipments in 1952, but slightly below the 1951 figure. . . . **During the eight-year period 1946–1953, over 28.5 million television receivers were shipped to dealers in the United States, Hawaii, and Alaska, according to a tabulation compiled recently by the RETMA Statistical Department.**

Federal Personnel

Capt. Rawson Bennett reported for duty recently as Assistant Chief of the Bureau of Ships for Electronics. He succeeds Admiral F. R. Furth, the new Chief of Naval Research. Capt. Bennett was returned to Washington from Schenectady, N. Y., where he has been assigned as Supervising Inspector of Naval Machinery and Naval Ordnance. Prior to that he was head of the Minesweeping Branch of the Bureau for Ships. Capt. T. W. Rogers, who has served as Acting Assistant Chief of the Bureau for Electronics since Admiral Furth's transfer, will resume his duties as Assistant Director for Electronics. . . . **The formal appointment of Dr. Robert D. Huntoon as Associate Director of the National Bureau of Standards for Physics has been announced by Dr. A. V. Astin, NBS Director.** Dr. Huntoon also is Acting Chief of the Electronics Division, and the Central Radio Propagation Laboratory. Wide recognition has been accorded Dr. Huntoon for his research on atomic-beam measurement, special amplifiers, atomic particle counting, electronic ordnance devices, the phasing of oscillators and deuteron-deuteron nuclear reaction studies. While Director of the NBS Corona, Calif., Laboratories, he was responsible for a research and development program that included guided missiles, electronic ordnance devices, digital computers and infrared measurements.

(Continued on page 88A)



Complete Testing Equipment for

OMNI and LOCALIZER RECEIVERS

A.R.C. Type H-14 Signal Generator

For a quick and accurate check by pilot before take-off, or for maintenance on the bench, this is the favored and dependable instrument. Checks up to 24 omni courses, omni course sensitivity, to-from and flag-alarm operation, and left-center-right on localizer. For ramp check, RF output 1 volt into 52 ohm line; for bench checks, 0-10,000 microvolts.

The H-16 Standard Course Checker is a companion instrument to the H-14. It makes possible a precise check on the course-accuracy of the H-14 or of any other omni signal generator. Just as a frequency meter is necessary in connection with a variable frequency signal generator, the H-16 Standard Course Checker is required in connection with a VOR signal generator for a precise measurement of phase accuracy.

These instruments sold only direct from factory.

Write for detailed literature

Dependable Airborne Electronic Equipment Since 1928

Aircraft Radio Corp
BOONTON, NEW JERSEY

NOW

NEW YORK—
73 RUTLEDGE ST., BROOKLYN 11, N. Y.

CARGO PACKERS EAST & WEST COAST PACKAGING SERVICE!

CALIFORNIA—
2050 BROADWAY, SANTA MONICA

Individually-engineered
CLIMATE-PROOF, SHOCK-PROOF PACKAGING

Three years in making—
three days in breaking!

12-PAGE BROCHURE tells how to avoid the pitfalls, gives complete detailed information on C-P methods of protective packaging.

To meet the increased, country-wide demands of the electronics industry. Cargo Packers has now inaugurated complete facilities in the West Coast area also, for packaging protection of delicate instruments, equipment and components, including

- **Special Assembly Equipment**
- **Experts in Military Specifications**
- **Economical Production line methods**
- **Full Compliance in Every Detail**
- **Competent Consulting Service**

Whether in storage or in transit anywhere, absolute protection against strain, stress, shock, vibration, temperature or humidity fluctuations, moisture—or even skin damage—is assured by Cargo Packers' individual attention to every job. For recommendations on specific packaging problems, call or write our Sales-Engineering Department today.

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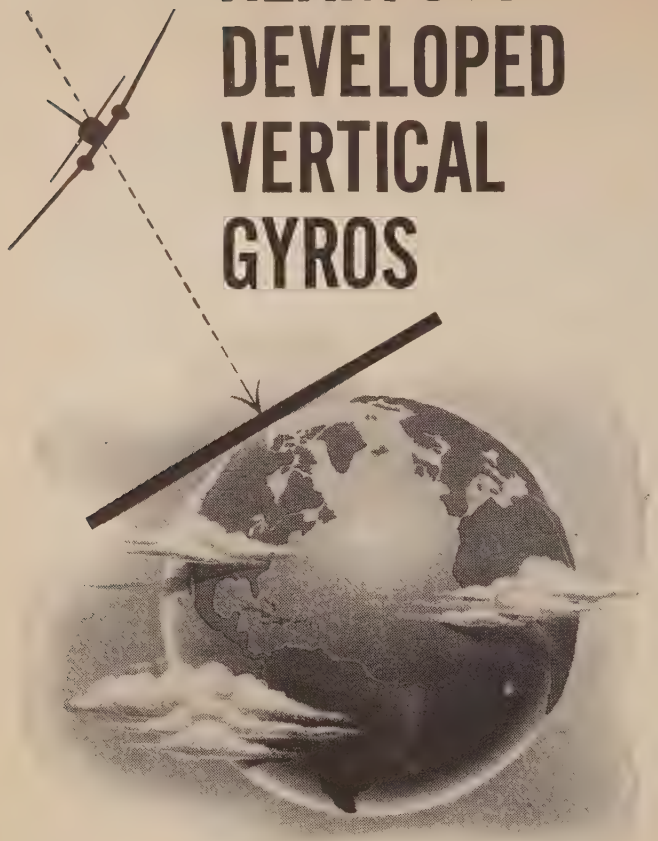
SPECIALISTS IN
CLIMATE-PROOF, SHOCK-PROOF
PACKAGING FOR ORDNANCE AND INDUSTRY

KEARFOTT DEVELOPED VERTICAL GYROS



...in production

KEARFOTT vertical gyros providing continuous vertical reference within two minutes of arc under bench conditions typify the engineering knowhow and production skills available to you in the field of precision gyros for airborne control applications.

Designed for particular applications with stringent performance requirements, a wide variety of vertical gyros now in production are being used extensively in aircraft and missile control systems demanding the most precise gyro reference obtainable.

Kearfott gyros incorporate many unique features permitting operation under extreme operational or environmental conditions. A true hermetic seal in dry inert gas provides positive environmental protection. Synchro pick-offs and rigid structural elements assure performance during adverse conditions of vibration or shock.

TECHNICAL DATA SHEETS

Complete technical data in tabular form on Kearfott Precision Vertical Gyros are available on request. Send for copies for your files. Write today.

KEARFOTT COMPONENTS INCLUDE:

Gyros, Servo Motors, Synchros, Servo and Magnetic Amplifiers, Tachometer Generators, Hermetic Rotary Seals, Aircraft Navigational Systems, and other high accuracy mechanical, electrical and electronic components.

Technical Data on these and other components is available on request.

Kearfott SINCE 1917

KEARFOTT COMPANY, INC., LITTLE FALLS, N. J.

Sales and Engineering Offices: 1378 Main Avenue, Clifton, N. J.
Midwest Office: 188 W. Randolph Street, Chicago, Ill. South Central Office: 6115 Denton Drive, Dallas, Texas
West Coast Office: 253 N. Vinedo Avenue, Pasadena, Calif.

A GENERAL PRECISION EQUIPMENT CORPORATION SUBSIDIARY

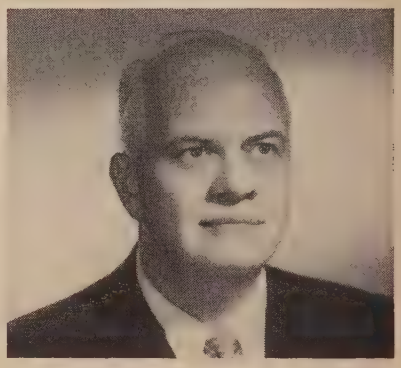
Industrial Engineering Notes

(Continued from page 86A)

National Television System Committee

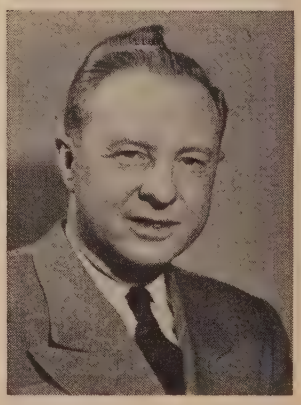
Dr. W. R. G. Baker, Chairman of the National Television System Committee and Director of the RETMA Engineering Department, at the close of the RETMA Board of Directors meeting received an award from the Academy of Television Arts and Sciences for the color television work of the NTSC. The "Emmy" statuette was presented to Dr. Baker by Director H. L. Hoffman who had previously accepted the award for Dr. Baker from the ATAS in Los Angeles. NTSC was given the "distinguished contribution" award "for evolving the present-day method of attaining compatible color."

Eaton Appointed VP by National



Raymond C. Cosgrove, Chairman of the Board of the **National Co., Inc.,** Malden and Melrose, Mass., has announced the appointment of Lynn Eaton as Vice-President in charge of sales.

New Sales Director for Servomechanisms



Servomechanisms, Inc., 500 Franklin Ave., Garden City, L. I., N. Y., has announced the appointment of C. H. Hartley to the position of Sales Director. Mr. Hartley is now located in the company's Corporate Offices in Garden City, New York.

A Revolutionary New Relay Development

of utmost importance to electrical and electronic design engineers

The Mullenbach Capaswitch uses an entirely new and different concept in relay design to transfer the contacts; *provides extreme sensitivity, low power requirements, high current-carrying capacity.*

The revolutionary new Capaswitch is basically an ultra-sensitive relay with unusual current carrying capacity. It will perform all of the jobs of conventional magnetic-coil relays, within the same current carrying capacity, plus many jobs that magnetic-coil relays cannot do. However, in design it departs radically from conventional relays. Instead of the usual electromagnetic armature, *a unique electrostrictive capacitive element provides the mechanical energy to open and close the contacts.* Only 0.5 milliwatt-seconds of operating power (150 volts d.c.) is required to close the contacts. To keep them closed requires less than 0.1 milliwatt, or *less than one-hundredth the power required to keep a conventional magnetic-coil relay closed!* This low power requirement opens up a vast new field of applications, eliminating need for much pre-amplified equipment.

How the Capaswitch works—Application of an actuating voltage creates a bending moment in the electrostrictive capacitive element, closing the contacts. Removal of the actuating voltage and discharge of the electrostatic element through external circuits or through a resistor, removes the bending moment, opening the contacts.

Time Delay Function—If appropriate resistances are applied in the circuit, the Capaswitch will function as a time delay relay to open or close the contacts. For longer time delays a larger condenser may be paralleled to the capacitive element.

Pulse Characteristics—Initial closing time of the Capaswitch is 10 milliseconds. However, it can be actuated by pulses as short as 10 microseconds or less. The electrostatic element may also be used to store low power pulses until sufficient voltage has been accumulated to operate the relay. However, present models cannot be used for accurate counting.

High Overvoltage Capacitance—Absence of a coil enables the Capaswitch to withstand wide voltage variations. As much as 200% overvoltage may be applied to the electrostatic element without damage. Low power requirements virtually eliminate heat and resulting dissipation problems, and reduce the number of needed components, saving space and weight.

Available now—Until recently the Capaswitch has been available only in limited quantities. Now, however, stepped-up production schedules assure increasing supplies.

OVERALL DIMENSIONS:
Length: 3½"
Width: 1-7/16" including solder terminals
Depth: 11/16"
Weight: 1.7 ounces

THE MODEL A-150 CAPASWITCH
a single pole, double throw relay, rated at 1 amp., 110 v., A.C. non-inductive load. Normal operating voltage 150 volts D.C.

Write today for complete specifications and prices!

Mullenbach
ELECTRICAL MANUFACTURING CO.
Established in 1927
2300 East 27th Street • Los Angeles 58, Calif.

New Multiple Function MICROWAVE VSWR AMPLIFIER

FEATURES:

- Crystal current and power monitoring
- Two channel input
- VSWR to 60 db
- Set-up signal sources

for faster, more accurate readings... easier operation

WAVELINE, precision leader in the manufacture of microwave test equipment, now offers you the ultimate in advance-design VSWR AMPLIFIERS.

This NEW Standing Wave Amplifier is the culmination of extensive research and testing . . . directed at developing new concepts of instrument function and design for microwave test equipment.

Exclusive combinations of features make WAVELINE microwave instruments the most valuable test equipment available today!

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SPECIFICATIONS

New **MODEL 2000**
Standing Wave Amplifier

Crystal Current Measurement – a feature is incorporated making the meter available for monitoring crystal current and power.

Two Channel Input – provides in one instrument:

1. By alternate use of two channels a pulsed oscillator in combination with a calibrated attenuator provides a substitute for a costly signal generator.

2. Monitoring crystal current and measuring VSWR.

3. Both channels measure VSWR.

4. Monitoring power with bolometer and measuring VSWR.

5. Monitoring power at two points.

Sensitivity – Full scale deflection; minimum 0.3 microvolts; maximum 0.3 volts.

Selectivity – Overall Q of approx. 20.

Calibration – Calibrated for use with a square law detector. 60 db over-all range in 6 steps. Accuracy ±0.1 db per 10 db.

Detector – Crystal rectifier or bolometer with 8.75 Ma. or 4.0 Ma. bolometer bias for standard 200 ohm bolometer, barretter or 1/100 amp instrument fuse.

Modulation Requirements – For VSWR measurement the RF source must be modulated at 1000 CPS ±20 CPS. Plug in units for frequencies 250 to 2500 CPS available.

Price – $200. F.O.B. Caldwell, N. J.

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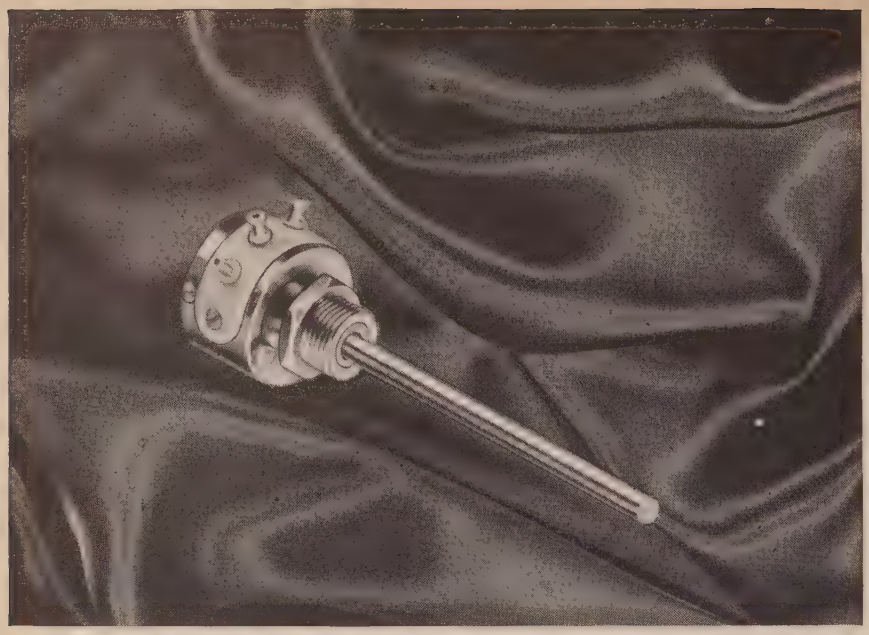

(Potentiometer is shown in actual size)

Another **NEW**

Fairchild Precision Potentiometer

The FilmPot TYPE 771

This metallic film potentiometer offers infinite resolution, high temperature operation (225° C.), high wattage dissipation, and 100 to 200,000 ohms resistance range in a case only ¾" in diameter and ½" long. The infinite resolution of a metallic film resistance element in servo applications limits hunting and oscillating. Available with servo flange or threaded bushing mounting. Gold-plated terminals. Now manufactured to target specifications for engineering evaluation; sample orders are accepted in standard resistance values only.

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Fairchild makes a complete line of precision potentiometers to fill all your needs—linear and nonlinear potentiometers, singly or in ganged combinations . . . single-turn and helical . . . with servo or threaded bushing mounts . . . and with resistance elements to meet your requirements.

Fairchild guarantees accuracy of ±1% or better in nonlinear types and ±0.5% or better in linear types. Highly accurate production methods and close mechanical tolerances, plus thorough type-testing and quality control, provide high resolution, long life, low torque and low electrical noise level in every Fairchild potentiometer. For more information, or for help in meeting your potentiometer problems, call on Fairchild Camera and Instrument Corp., Potentiometer Division, 225 Park Avenue, Hicksville, L. I., N. Y., Department 140-51H.

FAIRCHILD

PRECISION POTENTIOMETERS

Reports from Chapters

Aeronautical and Navigational Electronics

The Dayton Chapter of the Professional Group on **Aeronautical and Navigational Electronics** met on February 4 at the Engineers' Club of Dayton under the chairmanship of Paul Wiegert, Jr. E. G. Fubini of the Airborne Instrument Laboratory presented a paper entitled "Magnetic Airborne Detectors." Mr. Fubini reviewed briefly the problem involved in detecting and measuring very small variations in magnetic fields, such as are caused by the presence of a small ship or submarine in the earth's field. He discussed in some detail the various systems used in the past for measuring both the angle and the magnitude of the earth's field, indicating the advantages and faults of each. The system offering the greatest advantages and the one now used in detection of submarines from aircraft, utilizes small saturable reactors which generate a second harmonic component to an exciting carrier frequency when placed in a steady magnetic field. This provides by far the most sensitive system yet developed and will, under the service conditions, detect submarines at distances ranging up to 400 feet below the aircraft.

The Los Angeles Chapter of the Group met on February 2 at the IAS Building in Los Angeles, in a joint meeting with the Los Angeles Section. Dean E. Wooldridge of the Ramo-Wooldridge Corp., and Leslie Cromwell, Assistant Professor at Los Angeles State College, spoke on "Some Characteristics of Military Research and Development," and "A Look at Engineering Education in Great Britain," respectively.

Antennas and Propagation

The Los Angeles Chapter on the Professional Group on **Antennas and Propagation** met on December 15 at the IAS Building Los Angeles, under the chairmanship of M. J. Ehrlich. Carl Linnes, Jet Propulsion Laboratories, California Institute of Technology, presented a paper entitled "Some Influences of Over-all System Performance on Antenna Systems." Bob Stevens, also of the Jet Propulsion Laboratories, spoke on "Problems in Design of Low Side Lobe Level Antennas."

Audio

The Albuquerque-Los Alamos Chapter of the Professional Group on **Audio** met on February 4 at the Radiation Therapy Building, Lovelace Clinic, Albuquerque, under the chairmanship of A. M. Garblik. F. G. Hirsch, Medical Director of Sandia Corp., spoke on "Some Physiological and Psychological Considerations in High Fidelity Audio Reception."

(Continued on page 94A)

DIGEST

TIMELY HIGHLIGHTS ON G-E COMPONENTS

New electronic relays have high sensitivity

This new electronic resistance-sensitive relay is able to amplify minute currents carried by very delicate contacts. Even a wet thread will provide enough signal for it to operate.

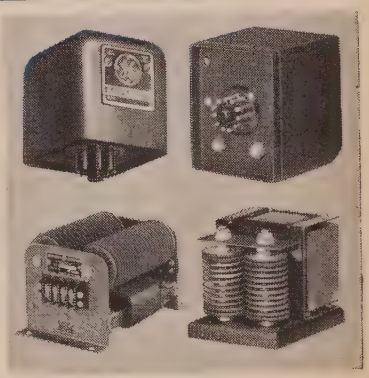
Sensitivity level is set by adjusting dial, which can be locked in place. The relay may be remotely controlled from as far away as 500 feet. Each can be set for either "normal" (relay "drops-out") or "reverse" (relay "picks-up") operation of the magnetic relay included in the device.

Built for long life, its enclosure is weather-resistant and dust-tight. Terminals are easily accessible; all components of this G-E relay are open for ease in servicing. For further information send for Bulletin GEA-5893.

Fast, accurate circuit analysis

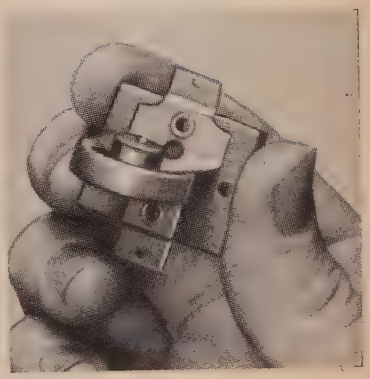
This self-contained, highly stable G-E self-balancing potentiometer rapidly converts small d-c voltages to measureable currents—*without* loading the measured circuit—for analysis of electronic circuits. It is consistently accurate because simple controls, and automatic, rapid circuit balance minimize operator errors. Easily changed resistor permits selection of input ranges from 100 microvolts to one volt d-c full scale with 5-milliampere d-c output. See Bulletin GEC-367.

Tiny signals amplified



Combining amplifying and rectifying elements in a unit, G-E amplistats (self-saturating magnetic amplifiers) "sense" small signal changes, amplify them greatly, and impart the amplified signal to a system to obtain the desired control. They give you the practical advantages of virtually instantaneous response, low power consumption, long life, and electrical signal isolation. Obtain assistance in applying G-E amplistats at your G-E Apparatus Sales Office. See Bulletin GEA-5950.

Small rectifier has high output



G-E germanium rectifiers offer the highest output in the smallest of rectifiers. For example, the dime-sized, sealed, air-cooled type is available in ratings up to 50 volts, 0.4 amperes d-c. Germanium rectifiers have these advantages: ***high efficiency***—operate 98% to 99% efficient; ***compactness***—small size and weight per watt output means you can build more compact assemblies; and ***long life***—two-year life tests show no detectable aging. Write for Bulletin GEA-5773.

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- ☐ **GEA-5893 Electronic Resistance Sensitive Relay**
- ☐ **GEA-5950 Amplistats**
- ☐ **GEA-6065 Micro-miniature Tantalytic Capacitors**
- ☐ **GEC-367 Self-balancing Potentiometer**

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PROFESSIONAL GROUP MEETINGS

(Continued from page 90A)

The Chicago Chapter of the Group met on February 19 at the Western Society of Engineers Building in Chicago. P. Williams and J. T. Williams were co-chairmen. J. T. Clark, El-Rad Manufacturing Co., presented a paper entitled "Engineer's Approach to Hi-Fi." It was a joint meeting with the Broadcast and Television Receivers' Chapter.

The Houston Chapter on the Group met on February 11 at the Southern States Life Insurance Co. Auditorium in Houston, under the chairmanship of L. A. Geddes. Rod Parvin and Don Freeman, sales engineers of Reeves Soundcraft Ltd., discussed the manufacturing techniques and general properties of the new Mylar magnetic recording tape. Following the discussion, demonstrations were given which illustrated the tensile strength temperature stability and fineness of finish now possible with this new recording medium. Commercial shapes and sizes were also exhibited and samples of the tape were given to the audience. This was a joint meeting with the Professional Group on Instrumentation.

The Philadelphia Chapter met on February 18 at WCAU Studios in Philadelphia. Dr. Daniel W. Martin, Acoustical Consultant for the Baldwin Co., spoke on "Musical Tone Radiation."

Circuit Theory

The Albuquerque-Los Alamos Chapter of the Professional Group on **Circuit Theory** met on January 27 at the University of New Mexico. Leo V. Skinner, Captain in the U.S.A.F. spoke on "Kirchoff vs. Lagrange" and "Stability Criteria." Col. Skinner presented the first part of his paper on "Kirchoff vs. Lagrange" at the December 9th meeting of the Group.

The Los Angeles Chapter of the Group met on January 14 at the Institute for Numerical Analysis in Los Angeles under the chairmanship of John Aseltine. Louis Weinberg, Research Physicist at Hughes Aircraft Co., spoke on "Transistors and Their Applications."

Communications Systems

The Washington, D. C. Chapter of the Professional Group on **Communications Systems** met on January 25 at the National Academy of Sciences Building in Washington, under the chairmanship of C. L. Engleman. Sidney T. Fisher, President of the Northwest Telephone Co., and Charles B. Fisher, President of Radio Engineering Products Ltd., spoke together on "Arctic VHF Radio Relay System Problems." The paper was discussed by Messrs. Jones, Toth, Wallace, Ould, Hessel and others. A number of color lantern slides were shown after the talk which not only illustrated the rather arduous conditions encountered in the Arctic, but also were found very entertaining.

(Continued on page 96A)

(Continued from page 94A)

Component Parts

The Philadelphia Chapter of the Professional Group on **Component Parts** met on February 15 at Towne Scientific School in Philadelphia, under the chairmanship of D. C. Bowen. Dr. W. T. Sackett of the Battelle Memorial Institute spoke on "Problems and Procedures in Obtaining Short-Term Life Ratings for Composition Resistors." Dr. Sackett pointed out that some military applications of electronic components require operation at higher than normal ambient temperature and loads. If this is done, the penalty of shorter lives must be paid. Before such applications can be made on a sound basis, procedures for obtaining short-term life ratings as a function of temperature, load, and altitude must be obtained. A previous paper given at the Eighth National Electronic Conference described a proposed step method of evaluating components. Dr. Sackett's talk presented an experimental evaluation of the proposed method, together with a discussion of problems of life definition, variability, conditioning, and sampling. The paper was discussed by all present.

Electron Devices

The San Francisco Chapter of the Professional Group on **Electron Devices** met on February 17 at Stanford University. Stanley F. Kaisel was acting chairman. Walter Haack of Western Gold and Platinum Co. of San Francisco spoke on "Brazing for Vacuum Tubes." The paper was discussed by H. A. Folgner, Manager of the Los Angeles Plant of Handy & Harman, Oakland, Calif.

Electronic Computers

The Washington, D. C. Chapter of the Professional Group on **Electronic Computers** met on January 6, at the PEPCO Auditorium in Washington under the chairmanship of C. V. L. Smith, Dr. C. E. Miller, of the ERA Division of Remington-Rand presented a paper entitled "The ERA 1103 Computer." This Chapter also met on February 3 to hear a paper presented by Joe Siller of the Diamond Ordnance Fuze Laboratory.

The Albuquerque Chapter of the Group met on January 26 for the election of officers. The following officers were elected: Chairman, J. P. Shoup, Vice-Chairman, Karl Ball, and Secretary-Treasurer, Julian E. Gross. Ralph McGehee was appointed Meetings and Papers representative. This Chapter also met on February 17 at Mitchell Hall at the University of New Mexico under the chairmanship of J. P. Shoup. Frank O. Lane presented a paper entitled "The Analog Computer as an Automatic Computing Machine."

The New York and Long Island Chapter of the Group met on December 1 at

(Continued on page 96A)

A NEW SUB-MINIATURE PRECISION QUARTZ CRYSTAL THAT USES ONLY 1/5th THE SPACE FORMERLY REQUIRED

The McCoy M-20 "McMite" has a frequency range of 5.0 mc. to 110 mc. It's a sub-miniature hermetically sealed unit, adaptable to multi-channel design for communications and frequency control equipment. Can be plugged into a sub-miniature tube socket, wired into miniature selector switch assembly or can be soldered to a printed circuit terminal board. Meets Military Specification requirements for fundamental operation above 5 mc. and overtone operation above 15 mc. The "McMite" does not sacrifice stability or dependability but meets same characteristics and performance requirements as larger crystal units used by the Military.

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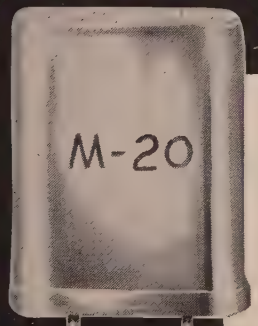
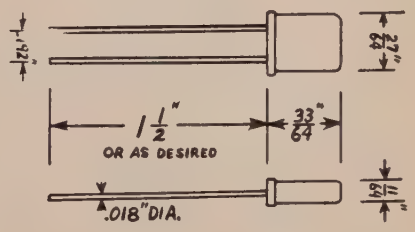
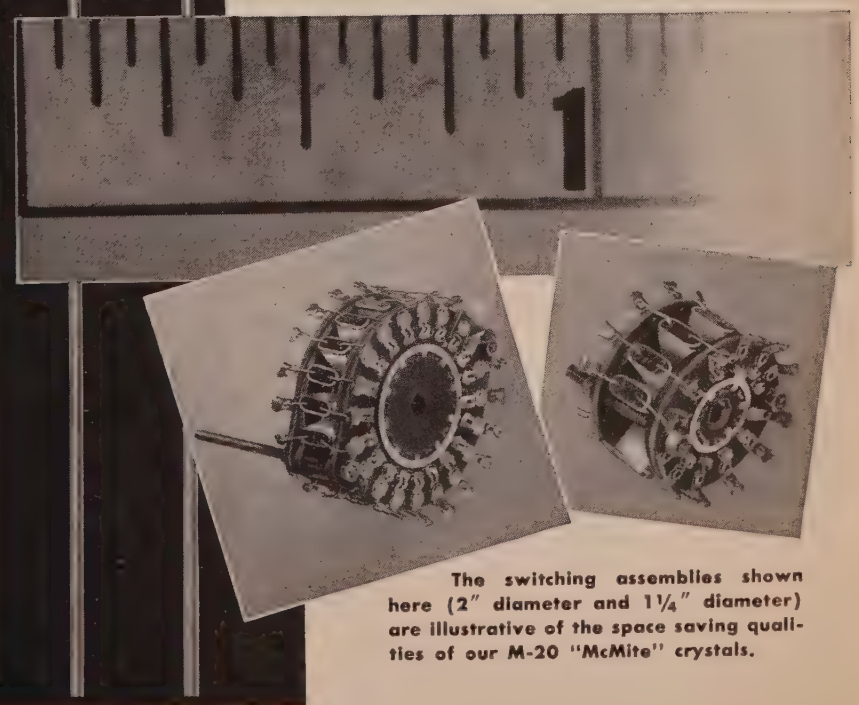
McCoy M-20 "McMite"







The switching assemblies shown here (2" diameter and 1¼" diameter) are illustrative of the space saving qualities of our M-20 "McMite" crystals.

• The largest of the porcelain rod insulators shown in the illustration develops 12,000 lb. strength. It is available, if specified, with rain shield and/or corona rings. All hardware is silicon aluminum alloy. Smaller units, in porcelain or steatite, are suited to lighter duty for strain or spreader use. Engineering and production facilities are also available for design and manufacture of units to almost any performance specification. Write for Bulletin 301, with complete description and specification data. Lapp Insulator Co., Inc., Radio Specialties Division, 213 Sumner St., Le Roy, N. Y.

PROFESSIONAL GROUP MEETINGS

(Continued from page 95A)

IBM World Headquarters in New York. John Bartelt, Engineer for IBM, spoke on the IBM 701 Computer. Election of officers and adoption of By-Laws was held at this meeting. This Chapter also met on January 26 at La Guardia Airport under the chairmanship of Daniel Haagens to hear Craig Andrews, E. L. Schmidt, and Charles Ammond of the Teleregister Corp. and American Air Lines speak on "The American Airlines Reservisor."

Information Theory

The Albuquerque-Los Alamos Chapter of the Professional Group on **Information Theory** met on January 13 at Mitchell Hall, University of New Mexico to hear Alex Fursa of the Sandia Corp. speak on Maxwell's Demon Cannot Work." Mr. Fursa also spoke at the February 10th meeting of the Chapter. His paper was entitled "Application of Information Theory and Negentropy to Physical Measurements."

The Washington, D. C. Chapter of the Group met on February 15 at the Engineers Club of Washington. The interim chairman was Charles R. Tieman. Dr. W. G. Tuller, Vice President of Engineering at Melpar Corp. presented a paper entitled "Information Theory Today." Permanent officers of the Chapter were elected to serve until June 1955: Chairman, Robert M. Page, Vice-Chairman, Ben C. Melton, and Secretary, Charles R. Tieman.

Radio Telemetry and Remote Control

The Los Angeles Chapter of the Professional Group on **Radio Telemetry and Remote Control** met on February 16 at the IAS Building in Los Angeles, under the chairmanship of J. R. Kauke. F. M. Riddle of the Jet Propulsion Laboratory, California Institute of Technology, spoke on "Transistors in Telemetry." Ed Perce, Chief Engineer of the Ralph M. Parsons Co., spoke on "Improved Decommutation Equipment."

Vehicular Communications

The Chicago Chapter of the Professional Group on **Vehicular Communications** met on February 19 at the Western Society of Engineers Building in Chicago, under the chairmanship of William J. Weisz. Edwin L. White, Chief of the Safety & Special Services Bureau of the Federal Communications Commission presented a paper entitled "Can Mobile Radio Grow?" The paper was followed by a general discussion of the question.

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★ ULTRA LOW capacitance & attenuation

TYPE	μμ F/ft	IMPED.Ω	O.D.
C 1	7.3	150	.36"
C 11	6.3	173	.36"
C 2	6.3	171	.44"
C 22	5.5	184	.44"
C 3	5.4	197	.64"
C 33	4.8	220	.64"
C 4	4.6	229	1.03"
C 44	4.1	252	1.03"

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AKRON

"Evolution of the Universe," by Dr. R. A. Alpher, Johns Hopkins University; March 16, 3954.

ATLANTA

Discussion and demonstration of precautionary measures taken to prevent failure of service at a broadcast station, by Ivan Miles, WGST; election of officers; February 19, 1954.

Films: "The Guiding Hand" and "For Progress—Call Dixie"; March 19, 1954.

BEAUMONT-PORT ARTHUR

Field inspection of KPRC-TV; February 24 1954.

"Capacitors," by W. M. Allison, Sprague Electric Company; February 25, 1954.

BINGHAMTON

"Color Television," by Lou Stantz, W.N.B.F.-AM-TV; February 17, 1954.

"Cathode Ray and TV Phosphors," by A. B. Davis, Sylvania Electric Products Inc.; March 8, 1954.

BUFFALO-NIAGARA

"Color Television—Its Challenge to the Engineer,"—written by Donald G. Fink. Speaker: Karl Wendt, Wendt-Squires, Inc.; February 17, 1954.

"Methods of Cooling Electronic Equipment," by J. P. Welsh, Cornell Aeronautical Labs. Inc.; March 17, 1954.

CEDAR RAPIDS

"Some Applications of the High Speed Differential Analyzer," by V. C. Rideout, Faculty, University of Wisconsin; February 17, 1954.

CHICAGO

"Operation of a Computing Center," by David Rubinfien, Computing Center; February 19, 1954.

"Color Television Receivers," by W. O. Swinyard, Hazeltine Research Inc.; March 4, 1954.

CLEVELAND

Symposium on Cleveland Industry—Dr. Begun, Chairman.

COLUMBUS

Panel on "Professional Ethics and Attitudes": Dr. R. C. McMaster, Battelle Memorial Institute; W. P. Corcoran, Ohio Bell Telephone Company; Paul Mazuzan, National Electric Coil Company; February 9, 1954.

"Ferrites in Microwave Applications," by J. H. Rowen, Bell Telephone Labs.; March 9, 1954.

DALLAS-FORT WORTH

"Development and Industrial Applications of the Transistor," by Harold Danchik; "Block Diagrams for TV Receivers," by E. M. Martin, Jr.; and "Investigation of Premature Firing of Blasting Caps," by Kenneth Zerbe; all students, Southern Methodist University; March 19, 1954.

DAYTON

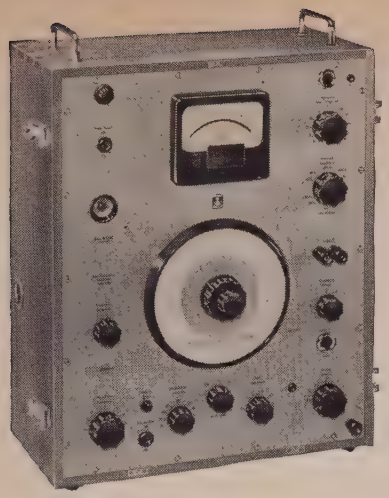
"Color Television," by C. N. Hoyler, David Sarnoff Research Center, RCA; March 17, 1954.

DES MOINES-AMES

"Getting Acquainted with Single-Sideband," by Warren Bruene, Collins Radio Co.; March 3, 1954.

"Computer Applications of Square Loop Magnetics," by Dr. S. M. Rubens, Engineering Research Associates; March 9, 1954.

(Continued on page 100A)



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This Beat Frequency Oscillator is designed to incorporate the many features required in a complete and flexible signal source such as — continuously tuned frequency range, metered output, frequency modulating circuit, variable compression and noiseless oscillator stop switch.

When coupled with the Bruel & Kjaer Level Recorder, the Oscillator automatically sweeps through the audio frequency range. This permits fast, accurate recording of frequency response, sound insulation properties, reverberation times, energy decay, etc. The Oscillator has a high frequency accuracy; 1%± one cycle per second. An incremental tuning adjustment and scale permit very accurate frequency setting. An automatic gain control circuit is utilized to maintain the desired output voltage for constant current, voltage, or sound pressure.

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BL-2002 Heterodyne Voltmeter
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BL-2109 Audio Frequency Spectrometer
BL-2304 Level Recorder
BL-2423 Megohmmeter and D. C. Voltmeter
BL-3423 Megohmmeter High Tension Accessory
BL-4002 Standing Wave Apparatus
BL-4111 Condenser Microphone
BL-4120 Microphone Calibration Apparatus and Accessory
BL-4708 Automatic Frequency Response Tracer

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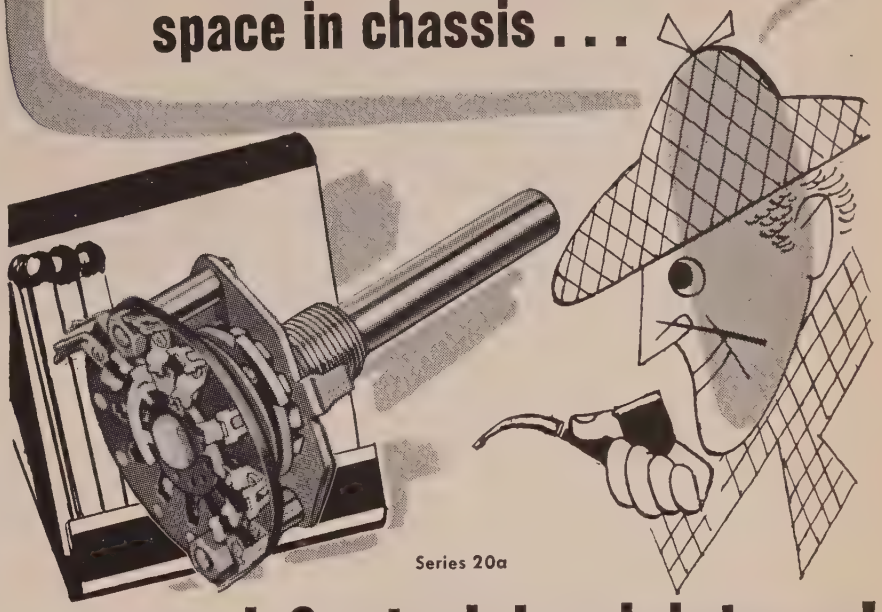

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- Standard or special combinations — up to 12 positions or up to 6 poles per section.
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- Phenolic insulation—only high grades used. NEMA Grade XXXP. Mil grade P3115B.
- Indexing 30° or 60° (standard or miniature). 90° (standard only).
- Shorting and non-shorting types.
- Now available—new Series 100 Sub-Miniature for military application only (1⅛" dia.).

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Series 20a

Series 20b

Series 30a

Series 30b

Series 30c

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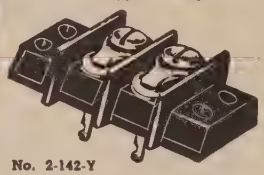
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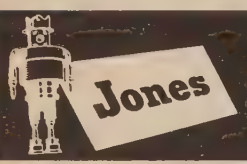

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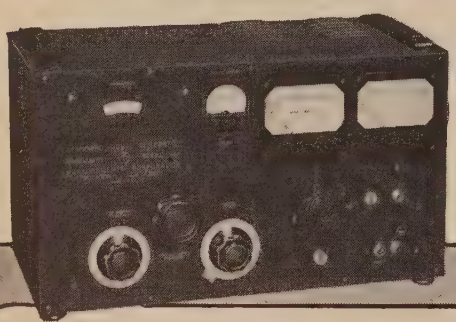

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BOONTON NEW JERSEY

SECTION MEETINGS

(Continued from page 98A)

DETROIT

"Instrumentation," by Seymour Sterling, Sterling Instrument Company; January 22, 1954.

"Fool Proof Electronic Control," by Stephen Hart, Electronic Control Corp.; February 19, 1954.

ELMIRA-CORNING

"Project Tinkertoy," by R. L. Henry, Nat onal Bureau of Standards; February 18, 1954.

"UNIVAC Electronic Computers," by J. D. Chapline, Remington Rand, Inc.; March 15, 1954.

EL PASO

"Our Expanding Technology," by Dr. J. W. McRae, Sandia Corporation; February 26, 1954.

EMPORIUM

"An Electronic Micrometer for Production Measurements," by Messrs. George Chilton and Ralph West, Sylvania Electric Products, Inc.; February 25, 1954.

EVANSVILLE-OWENSBORO

"Design of Microwave Systems," by A. E. Kowitz, Motorola, Inc.; March 10, 1954.

HAMILTON

"Human Relations in Industry," by L. C. Sentance, Canadian Westinghouse Co.; January 11, 1954.

"The Effects of Weather on the Performance of Marine Radar Over Lake Ontario," by Adam Hood, National Research Council of Canada, February 8, 1954.

HUNTSVILLE

"Tremendous Trifles," by Raymond Moeller, General Electric Company; February 16, 1954.

"Guided Missiles Are a Must," by D. E. Mullen, General Electric Company; March 2, 1954.

INYOKERN

"The Transistor"—demonstration-lecture by Charles Develon, Pacific Telephone Company of Los Angeles; March 1, 1954.

LITTLE ROCK

"The Use of Radio in Petroleum Operations," by Dr. W. M. Rust, Jr., Humble Oil Co.; and film, "Project Tinkertoy" March 9, 1954.

LONDON

"The Production of Oscillations with Gas Tubes," by Mr. Greenwood, Student, University of Western Ontario; "Improvements in Interlacing in TV Receivers," by Mr. Southern, Student, Ryerson Institute of Technology; "The Ripple Tank as an Aid to the study of Phase Fronts," by Mr. Clark, Student, University of Western Ontario; Problems Associated with Large Scale Public Address Systems," by Mr. Godfrey, Student, Ryerson Institute of Technology; "The Electric Resistance Strain Gage," by Mr. Broughton, Student, University of Western Ontario; and "Symbolic Logic in Practice," by Mr. Hewitt, Student, University of Toronto; February 22, 1954.

"Semi-Automatic Communications," by Sqdn. Leader M. Lemke, RCAF; and "Mobile Communications in the RCAF Today," by Sqdn. Leader Fitzgerald, RCAF; March 8, 1954.

(Continued on page 102A)

See pages 161A to 164A for IRE Directory Questionnaire

New converters extend range of -hp- 524A Frequency Counter to 220 mc!

- Direct readings 10 cps to 220 mc
- Increases sensitivity
- No loss in accuracy
- Instant, automatic measurement
- Easily used by non-technical personnel
- Low cost; eliminates complex setups

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Standard Mixing Frequency: Multiplied from 100 kc output of 524A. (Accuracy 1/1,000,000 short term.)

Output Voltage: 2 v minimum (sufficient to operate *-hp-* 524A over entire range).

Connectors: Standard BNC type.

Power Source: 115/230 v, 50/60 cps, 60 watts.

Size: 7" x 10¾" x 14" deep. Weight 20 lbs. Rack mount available.

Accessories Furnished: Cables to interconnect with *-hp-* 524A.

Price: $350.00.

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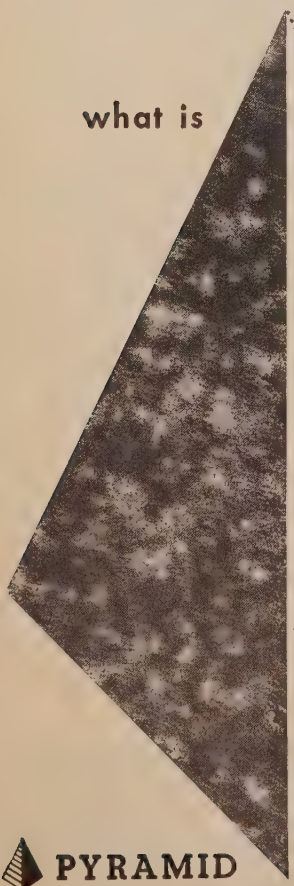

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SECTION MEETINGS

(Continued from page 100A)

LONG ISLAND

"The Surface-Barrier Transistor," by W. H. Forster and Dr. J. B. Angell, both of Philco Corp.; February 9, 1954.

Presentation of Fellow Awards. Speaker: Prof. J. R. Ragazzini, Columbia University. Also, "The Growth of the IRE and the Importance of the Fellow Award," by J. V. L. Hogan, Hogan Labs.; March 7, 1954.

LOS ANGELES

"Geophysics and Geophysical Exploration," by M. L. Swann, United Geophysical Company; "Cinemascope and Steriophonic Sound," by Lorin Grignon, Twentieth-Century Fox Corp.; dinner speaker: "Operation Deepfreeze," by E. R. Freeto, Lockheed Aircraft, Inc.; March 2, 1954.

LOUISVILLE

"Future Sources of Power," by R. W. Ferguson, Westinghouse Electric Company; February 26, 1954.

"Project Tinkertoy or Modular Design and Mechanized Production of Electronic Mechanisms," by Dr. B. L. Davis, National Bureau of Standards; March 11, 1954.

MONTREAL

"Radio Astronomy," by A. E. Covington, National Research Council; February 24, 1954.

NEW ORLEANS

"The Advent of Television," by G. Mayoral, WJMR-TV and Jack Petrik, Supreme Broadcasting Company; March 5, 1954.

NEW YORK

"Solid State Amplifiers," by Ernst Weber, Polytechnic Institute of Brooklyn; March 3, 1954.

NORTH CAROLINA-VIRGINIA

"NTSC Color Television Standards," by C. J. Simon, General Electric Company; March 5, 1954.

OMAHA-LINCOLN

Demonstration and discussion of the Scotchlock Wire Connectors and the resin products, by E. W. Bollmeir, Electro-Mechanical Products; discussion and demonstration of sound recording tape, by R. A. Von Behren, Magnetic Products Lab.; February 18, 1954.

OTTAWA

"Method for Time or Frequency Compression—Expansion of Speech," by Dr. W. L. Everitt, Dean, University of Illinois College of Engineering; February 18, 1954.

"Built-In Brains," by W. J. Turnbull, Deputy Postmaster General of Canada; March 11, 1954.

PHILADELPHIA

"Ferroelectric Storage Devices," by J. R. Anderson, Bell Telephone Labs.; February 9, 1954.

"Magnetic Drums and Tapes," by J. Rabinow and S. N. Alexander, National Bureau of Standards; February 16, 1954.

"Ferromagnetic Storage Devices," by T. H. Bonn, Remington Rand, Eckert Mauchly Division; February 23, 1954.

"Studio and Remote Pickup Equipment," by J. H. Roe, RCA; February 25, 1954.

"Electrostatic Storage Devices," by J. H. Pomerene, Institute for Advanced Study; March 2, 1954.

(Continued on page 103A)

SECTION MEETINGS

(Continued from page 102A)

"Acoustic Storage Devices," by J. F. Koch, Jr., Technitrol Engineering Co. Inc.; March 9, 1954.

"Film Problems in Color TV," by A. N. Goldsmith, Chairman, NTSC Panel II; "Color TV From Film," by E. H. Traub, Philco Corp.; "Color Signal Recording Techniques," by E. D. Goodale, N.B.C.; and "Network Transmission of Color Signals," by J. R. Rae, American Tel. and Tel. Company; March 11, 1954.

"Summary and Evaluation," by A. L. Samuel, IBM Corp.; March 16, 1954.

"Techniques of Color Broadcasting," by R. A. Monfort, N.B.C.; "Transmitter Specifications," by T. M. Gluyas, RCA; and "Specially Developed Test Equipment and Test Methods," by J. W. Wentworth, RCA; March 19, 1954.

PITTSBURGH

Inspection Trip of Bell Telephone 4A Toll Switching System, Automatic Message Accounting Center, and TV Control Center. This was a Joint Meeting with AIEE, Engrs. Soc. of W. Pa.; March 9, 1954.

ROME-UTICA

"Packaging of Speech Communications by Means of Pre-Emphasis and Peak Clipping Methods," by Dr. Oliver Straus, National Company; March 2, 1954.

SACRAMENTO

"Television Station Equipment," by Harry Bartolomei, KCCC-TV; February 26, 1954.

"Color Television Systems," by W. E. Evans, Stanford Research Institute; March 12, 1954.

(Continued on page 104A)

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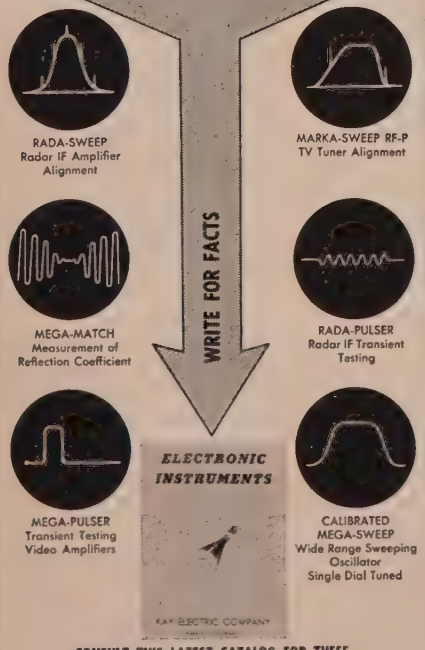

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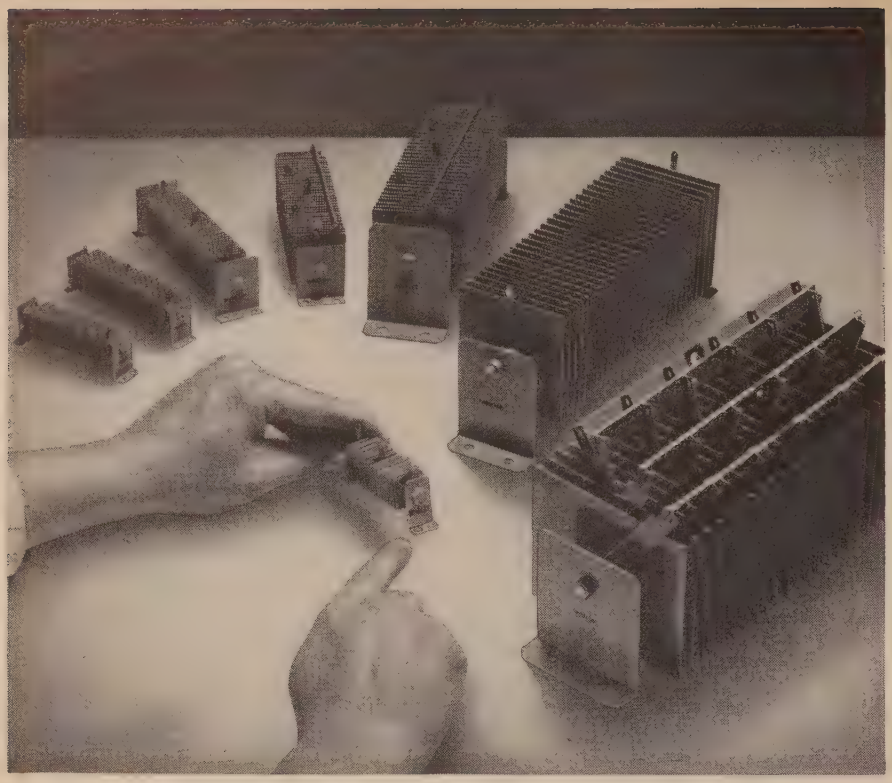

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SECTION MEETINGS

(Continued from page 103A)

SAN DIEGO

"Two Large-Scale Real-Time Computers," by Stanley Rogers, Consolidated Vultee Aircraft Corp.; March 9, 1954.

SYRACUSE

"Application of Correlation Function to Communications Problems," by Dr. R. E. Johnson, Syracuse University; February 4, 1954.

"Operations Research," by Sherman Kingsbury, Arthur D. Little, Inc.; February 23, 1954.

TULSA

"The Application of Matrix Algebra to Circuit Analysis," by R. N. Norman; "An Electronic Ignition System," by E. M. Barnes, Jr.; and "Can The Yagi Antenna Solve Your TV Reception Problem?" by T. B. Hall, all Students; March 11, 1954.

TWIN CITIES

"Problems in the Development of a Power Transistor," by Dr. Van W. Bearinger, Mineapolis-Honeywell; March 10, 1954.

VANCOUVER

"Theoretical Considerations in the Correction of Loudspeaker Distortion," by P. G. Scott; "The FTL-14 Direction Finder," by G. E. Forman; and "Distributed Amplifiers," by W. R. Tracy, all Students, University of B. C.; February 15, 1954.

WASHINGTON

Matching Opportunity with Need," by J. W. McRae, Sandia Corporation; March 8, 1954.

WINNIPEG

"Manitoba Power Commission VHF Radio System," by E. H. Edge, Rogers-Majestic Electronics Ltd.; and "Microwave Systems," by T. G. Lynch, Canadian General Electric Co. Ltd.; March 2, 1954

SUBSECTIONS

AMARILLO-LUBBOCK

"Application of Symbolic Logic in Design of Switching Circuits," by W. G. Breckenridge, Student winner of annual paper contest; January 26, 1954.

ITHACA

"Photoconductivity," by Albert Rose, RCA Labs.; January 22, 1954.

LANCASTER

"Transistor Circuit Applications," by L. G. Schimpf, Bell Telephone Labs.; March 10, 1954.

MID-HUDSON

"Construction, Characteristics and Performance of Capacitors in Conventional and High Reliability Circuits," by D. B. Peck, Sprague Electric Company; February 18, 1954.

MONMOUTH

"Kodachrome Travelogue of Europe," by G. W. Crawford, RCA Victor; February 17, 1954.

NORTHERN NEW JERSEY

"Electronic Digital Computers," by Dr. J. W. Mauchley, Remington Rand, Inc.; February 10, 1954.

TUCSON

"Transistors and Their Applications," by Tudor Finch, Bell Labs.; February 12, 1954.

USAFIT

"Applications of Transistors," by R. D. Alberts, Aircraft Radiation Lab.; February 26, 1954.

ALABAMA POLYTECHNIC INSTITUTE, IRE BRANCH

General meeting; February 22, 1954.

BROWN UNIVERSITY, IRE-AIEE BRANCH

General meeting; February 17, 1954.

UNIVERSITY OF CALIFORNIA, IRE-AIEE BRANCH

Executive Committee meeting; February 16, 1954.

General meeting; February 23, 1954.

Executive Committee meeting; March 2, 1954.

"Our Interesting Universe," by Dr. L. Reukema, Faculty, University of California; March 5, 1954.

"Transistors,"—demonstration by Ray Tiffany, Pacific Telephone; March 15, 1954.

Business meeting; March 16, 1954.

"The Swindletron"—talk and demonstration by Prof. J. Woodyard, University of California; March 19, 1954.

CALIFORNIA INSTITUTE OF TECHNOLOGY, IRE BRANCH

"High Frequency Heating," by W. A. Dent, Westinghouse Electric Corp; February 8, 1954.

CALIFORNIA STATE POLYTECHNIC COLLEGE, IRE BRANCH

"Engineering at Hewlett-Packard," by Noel Porter, Hewlett-Packard; "100 Megacycle Counter." by R. W. Forsblad and D. L. Palmer, Hewlett-Packard; March 4, 1954.

CARNEGIE INSTITUTE OF TECHNOLOLY, IRE-AIEE BRANCH

Film, "Voices in Paper" and Talk, "Position Interviews," by O. P. Robinson, Cutler-Hammer Company; December 10, 1953.

"The Future of the Electric Power Industry," by A. C. Montieth, Westinghouse Electric Corporation, December 16, 1953.

Election of officers; February 19, 1954.

CASE INSTITUTE OF TECHNOLOLY, IRE BRANCH

"Electronics in Nuclear Physics," by Prof. E. C. Gregg, Faculty, Case Institute of Technology; March 23, 1954.

COLORADO A & M COLLEGE, IRE-AIEE BRANCH

General meeting and film, "Steels Party Line"; February 17, 1954.

"History of AIEE," by Albert Anderson, AIEE and General Electric; March 3, 1954.

COLUMBIA UNIVERSITY, IRE-AIEE BRANCH

Business meeting; February 25, 1954.

UNIVERSITY OF CONNECTICUT, IRE-AIEE BRANCH

"Automatic Pilots," by E. J. Isbister, Sperry Gyroscope Company; January 14, 1954.

UNIVERSITY OF DAYTON, IRE-AIEE BRANCH

General meeting; January 5, 1954.

Talks by C. Marshall, Regional Director; A. Petit and A. B. Henderson, Chairman and Vice Chairman of the Dayton Section, and John Heyt, member of Professional Group, Dayton Section; February 11, 1954.

"Ionosphere and Its Effect on Radio Waves," by Wilbur Chang, Student; February 18, 1954.

"Color Television," by Mr. Dominic, Student; February 25, 1954.

(Continued on page 106A)

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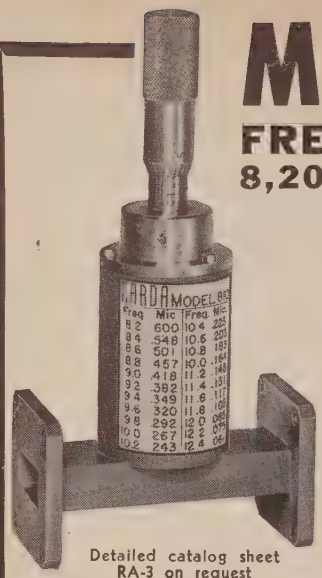

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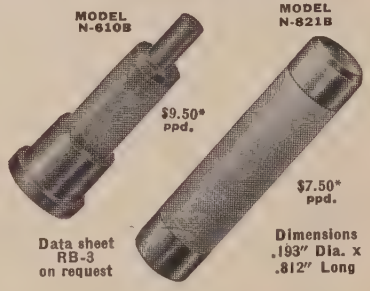

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STUDENT BRANCH MEETINGS

(Continued from page 105A)

UNIVERSITY OF DELAWARE, IRE-AIEE BRANCH

"Customer Toll Dialing," by W. F. Denkhouse, Plant Extension Engineer; January 13, 1954.

"Scientific Training in United States and Europe," by Dr. W. A. Mosher, Faculty, University of Delaware; January 18, 1954.

"Professional Registration for the Engineer," by Dean D. L. Arm, University of Delaware; March 1, 1954.

FENN COLLEGE, IRE BRANCH

Business meeting; January 26, 1954.

Business meeting and election of officers; February 14, 1954.

UNIVERSITY OF FLORIDA, IRE-AIEE BRANCH

"Flying Saucers," by C. B. Searles, Vice Chairman of the Student Branch; film, "Hidden World"; February 21, 1954.

General meeting and election of officers; March 8, 1954.

GEORGE WASHINGTON UNIVERSITY, IRE BRANCH

Quiz of six selected members for prize sponsored by Washington Section; March 3, 1954.

UNIVERSITY OF ILLINOIS, IRE-AIEE BRANCH

General Electric exhibit "House of Magic" with W. R. Whitmore and Richard Bogh, of General Electric; election of officers; February 24, 1954.

"Military Electronics at Bell Labs.," by W. H. C. Higgins, Bell Telephone Labs.; March 10, 1954.

STATE UNIVERSITY OF IOWA, IRE BRANCH

"Collins Intergrated Flight System," by Mr. Westcott, Collins Radio Company; IRE Paper Contest; March 3, 1954.

KANSAS STATE COLLEGE, IRE BRANCH

"Use of Precision Snap-Acting Switches in Safety Applications," by A. L. Riche, Micro Switch; February 25, 1954.

"Production of Synthetic Crystals for Telephone Work," by Fred Masek, Bell Telephone Labs.; March 11, 1954.

LAFAYETTE COLLEGE, IRE-AIEE BRANCH

"Sales Engineering," by Howard Lovett, of Roller Smith; January 11, 1954.

"Electronic Umpire," by Richard Soiler, Student, Lafayette College; February 29, 1954.

LEHIGH UNIVERSITY, IRE BRANCH

"The College Graduate in Public Utilities," by W. Astley, Philadelphia Electric Company; February 18, 1954.

UNIVERSITY OF MARYLAND, IRE-AIEE BRANCH

Demonstration: TV and Micro Wave Equipment, by Robert Gish, Philco Corp.; March 10, 1954.

MASSACHUSETTS INSTITUTE OF TECHNOLOGY, IRE-AIEE BRANCH

"Color Television," by C. N. Hoyler, David Sarnoff Research Labs., RCA; March 9, 1954.

MICHIGAN STATE COLLEGE, IRE-AIEE BRANCH

Election of officers; March 3, 1954.

(Continued on page 108A)

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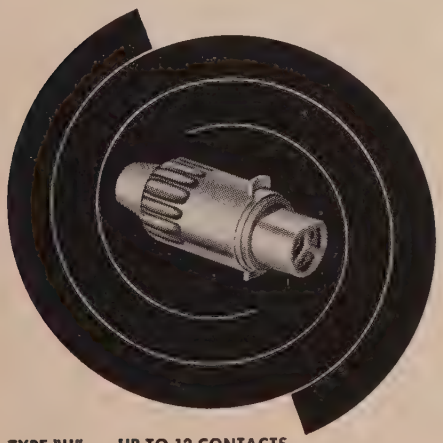
TYPE "D"... UP TO 50 CONTACTS

For Your Miniaturization Program!

Save space... save weight... rugged. 1000's of uses! Pressed steel shells. Heavily *gold-plated* copper-alloy 5-amp. contacts give untarnishing high surface conductivity and corrosion resistance. Take #20 solder-tipped wires. Polarization means. Salt-spray resistant shells. Flashover 1700v dc.



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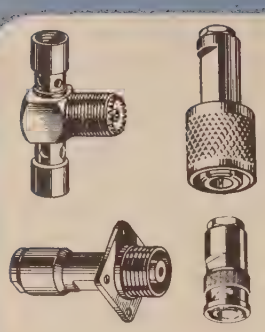
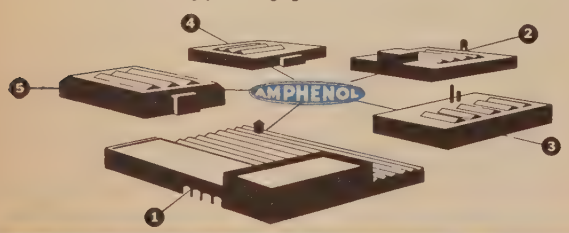


Yes, in a way, a second helping—for we are trying to express the factors which have helped raise the AMERICAN PHENOLIC CORPORATION by its boot straps in a little over twenty years. In the early thirties AMPHENOL was just another struggling industrial infant in one small factory and our main concern then was radio sockets. Today there are five modern plants producing over 11,000 cataloged products ranging all the way from AN connectors to television antennas.

"A little more of everything"—we think a little more courage, a little more imagination, a little more faith have been responsible for the growth of AMPHENOL during the last two decades. And we look forward to the same "little more" qualities, a second helping of the same, to keep AMPHENOL a vital member of the electronics industry for many, many years to come.



AMERICAN PHENOLIC CORPORATION — chicago 50

STUDENT BRANCH MEETINGS

(Continued from page 106A)

UNIVERSITY OF MINNESOTA, IRE-AIEE BRANCH

"What Can Computers Do For You," by P. N. Hess and D. W. Fenne, University of Minnesota Computing Center; February 3, 1954.

Inspection Trip to Minneapolis-Honeywell Company Research Lab.; February 18, 1954.

"Problems in the Development of a Power Transistor," by Dr. Van W. Bearinger, Minneapolis-Honeywell Regulator Company; March 10, 1954.

MISSOURI SCHOOL OF MINES & METALLURGY, IRE-AIEE BRANCH

General meeting and films, "The Microwave Oscillator" and "The Atom Goes to Work"; February 25, 1954.

MONTANA STATE COLLEGE, IRE-AIEE BRANCH

"Street Lighting," by Mr. Graff, General Electric Company; February 8, 1954.

"Television Broadcasting and Receiving," by R. C. Seibel, Faculty, Montana State College; February 25, 1954.

"Power Systems Analyzer," by J. F. Fuller, General Electric Company; March 4, 1954.

NEW MEXICO COLLEGE OF A & M ARTS, IRE-AIEE BRANCH

"What Management Expects of New Employees," by Jack Bowen, El Paso Electric Company; February 11, 1954.

Election of officers; March 12, 1954.

COLLEGE OF THE CITY OF NEW YORK, IRE BRANCH

General meeting; February 11, 1954.

"Digital Computers," by Donald Rosenheim, IBM; February 18, 1954.

"Regulating Systems," by Mr. Redmon, Westinghouse Electric Corp.; February 25, 1954.

"Problems in Power Generation," by M. F. Kent, General Electric Company; March 4, 1954.

"Wave Guides," by A. C. Beck, Bell Telephone Labs.; March 18, 1954.

UNIVERSITY OF NORTH DAKOTA, IRE-AIEE BRANCH

General meeting and film on transistors; March 10, 1954.

OKLAHOMA A & M COLLEGE, IRE-AIEE BRANCH

"Suicide of a High Voltage Arc," by G. H. Mahke, Line Material Company; February 18, 1954.

Student paper competition: "The Application of Matrix Algebra to Circuit Analysis," by R. N. Norman; "An Electronic Ignition System," by E. M. Barnes, Jr.; and "Can the Yagi Antenna Solve Your TV Reception Problem?" by T. B. Hall; March 11, 1954.

"Mechanized Intelligence," by W. Keister, Bell Telephone Labs.; Election of officers; March 15, 1954.

OREGON STATE COLLEGE, IRE BRANCH

Film, "Prelude to Kitimat"; March 2, 1954.

UNIVERSITY OF PENNSYLVANIA, IRE-AIEE BRANCH

Election of officers; February 12, 1954.

General meeting; February 21, 1954.

"The Physics of Music and Hearing," by Dr. W. E. Koch (on tapescript), Bell Telephone Laboratories; February 24, 1954.

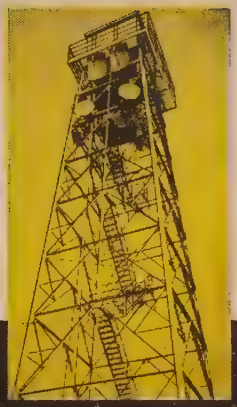
"Electric Cardiology," by E. Frank, Faculty, Moore School of Elec. Eng'g.; March 8, 1954.

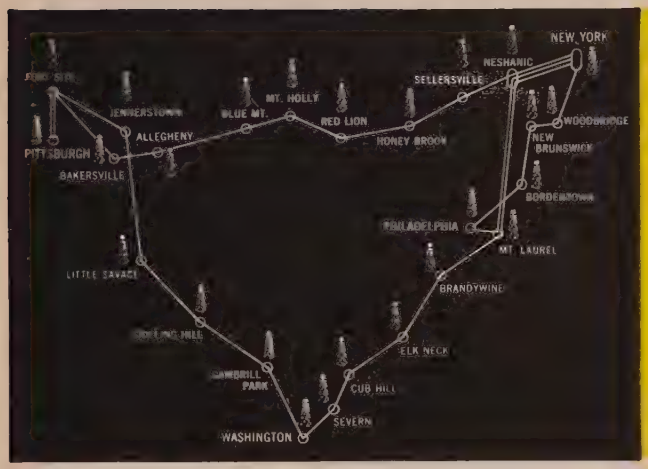
"Radiology," by Dr. S. R. Warren, Jr., Asst. Vice Pres. in Charge of Undergraduate Affairs; March 15, 1954.

(Continued on page 110A)

IN ITS 21 TOWER MICROWAVE SYSTEM

Western Union uses Sperry Klystrons



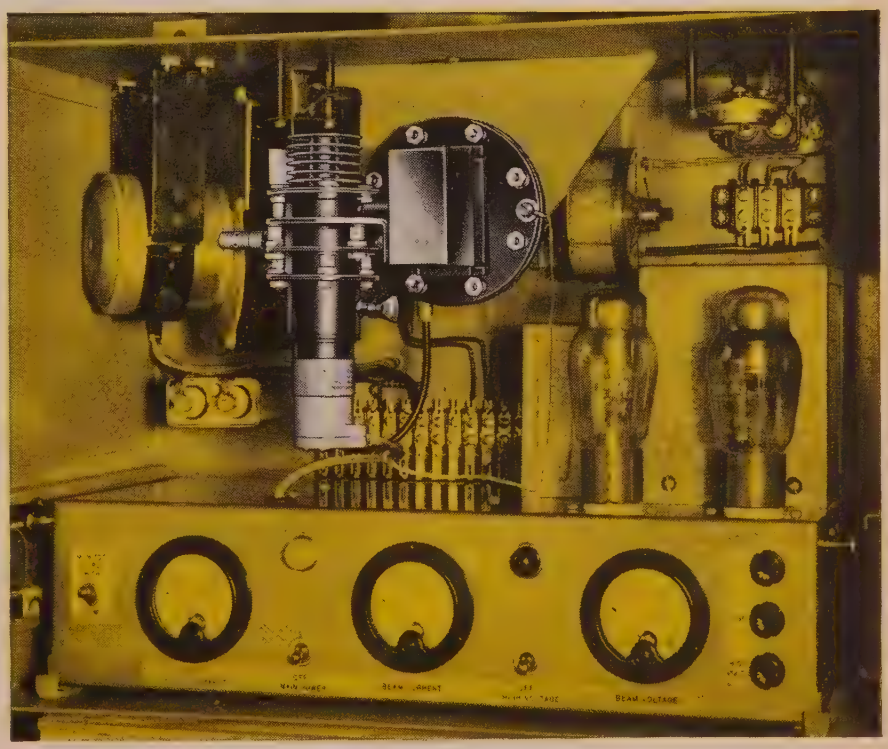


"Dependability has been excellent," writes H. P. Corwith, Vice President, Development and Research. "Average tube life has been more than 15 months, and some tubes have been in continuous service for almost 3 years."

■ Installed in 1948, the Western Union microwave system between New York, Philadelphia, Washington and Pittsburgh, consists of 21 towers, varying in height from 60 to 120 feet, and spaced up to 55 miles apart. The system handles hundreds of telegraph circuits—including important government and leased private wire systems as well as circuits for regular message traffic.

■ Through the use of Sperry SAC-41 Klystrons providing power output of 10 watts, Western Union has effectively reduced circuit outages due to fading, and provided dependable service under all conditions. Furthermore, as Mr. Corwith points out above, the average life of Sperry SAC-41 Klystrons has been 15 months—and some tubes have served continuously for almost 3 years.

■ Since 1938, when Sperry sponsored the development of the Klystron, this Company has extended its application to tubes for low, medium and high power applications—and in a frequency range from 750 to 40,000 megacycles. The research, development and specialized production facilities of Sperry in Klystrons — and in accompanying Microline* equipment—are at your disposal to help in solving your problem.



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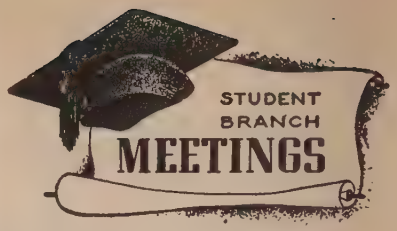

(Continued from page 108A)

RENSSELAER POLYTECHNIC INSTITUTE, IRE-AIEE BRANCH

"Engineering Considerations in the Design of Nuclear Power Reactors," by K. A. Kesselring, Knolls Atomic Power Lab.; film, "A is for Atom"; February 16, 1954.

UNIVERSITY OF RHODE ISLAND, IRE-AIEE BRANCH

General meeting; March 17, 1954.

RUTGERS UNIVERSITY, IRE-AIEE BRANCH

"Color Television," by Dr. Brown, RCA Labs.; February 4, 1954.

"Amateur Radio and Civil Defense," by R. Forsyth, Civil Defense; February 18, 1954.

Talk by Messrs. Woody and Frick, Westinghouse Electric Corp.; February 19, 1954.

Election of officers; March 4, 1954.

Field trip through New Jersey Bell Tel. Co., conducted by Robert Miller; March 17, 1954.

SAN DIEGO STATE COLLEGE, IRE BRANCH

Film, "Germanium, the Magic Metal"; March 16, 1954.

UNIVERSITY OF SOUTHERN CALIFORNIA, IRE-AIEE BRANCH

"The Training Program of General Electric," by Walter Scott, General Electric Company; March 3, 1954.

"The Ground Approach Control for Aircraft," by E. R. Swanson, Gillfillen Radio Corporation; March 10, 1954.

STANFORD UNIVERSITY, IRE-AIEE BRANCH

"Some Circuit Properties of Junction Transistors,"—Bell Telephone Lab. Tapescript; February 25, 1954.

TEXAS TECHNOLOGICAL COLLEGE,

"The Practical Aspects of Engineering," by Bob Wilson, Electrical Dept. of the City of Lubbock; March 1, 1954.

Election of officers; March 15, 1954.

UNIVERSITY OF TOLEDO, IRE-AIEE BRANCH

General meeting; January 13, 1954.

General meeting; March 3, 1954.

UNIVERSITY OF TORONTO, IRE-AIEE BRANCH

Student night sponsored by Toronto Section. Papers presented: "Symbolic Logic Applicable to Switches," by A. Hewitt; "Cathode Interface Resistance," by R. E. Hobson; "Ground Approach Control," by R. Richardson; and "Transistors," by R. Charette; February 15, 1954.

"Engineering Related to Microwave, Broadcast and T.V. Equipment," by R. A. Muller, Canadian General Electric Company; February 19, 1954.

"Voltage Regulators," by F. M. Squires, Ferranti Electric Ltd.; February 26, 1954.

"Design and Application of Electronic Computers," by Mr. Aitchson, IBM Co. Ltd.; film, "Piercing the Unknown"; March 5, 1954.

Two films, "Wires Under Water" and "Manufacture of Paper Insulated Cables." Speaker: H. C. Graham, Northern Electric Co. Ltd.; March 12, 1954.

UNIVERSITY OF UTAH, IRE-AIEE BRANCH

"Unified Engineering Program," by A. D. Hill Student, University of Utah; February 2, 1954.

Business meeting; February 9, 1954.

Film, "Okenite Plant"; February 16, 1954.

(Continued on page 112A)

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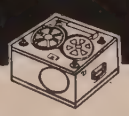

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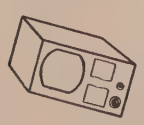

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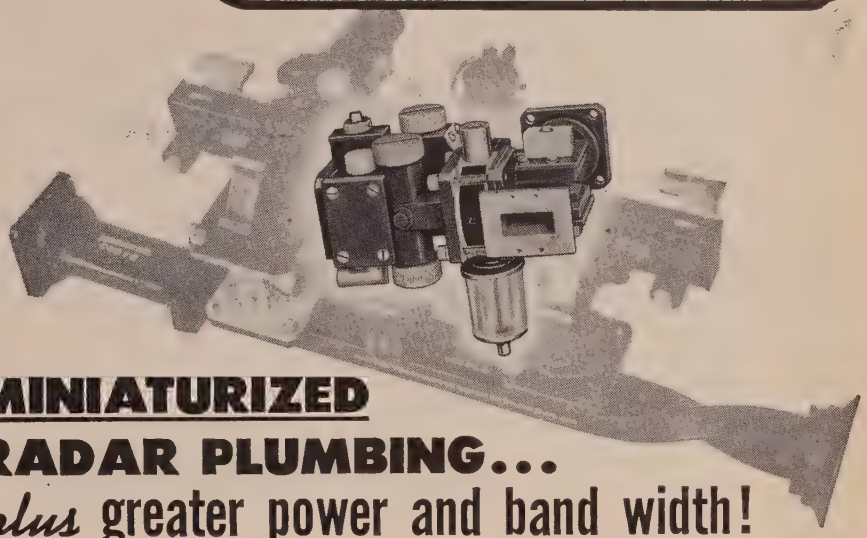

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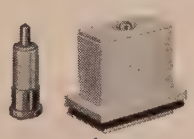

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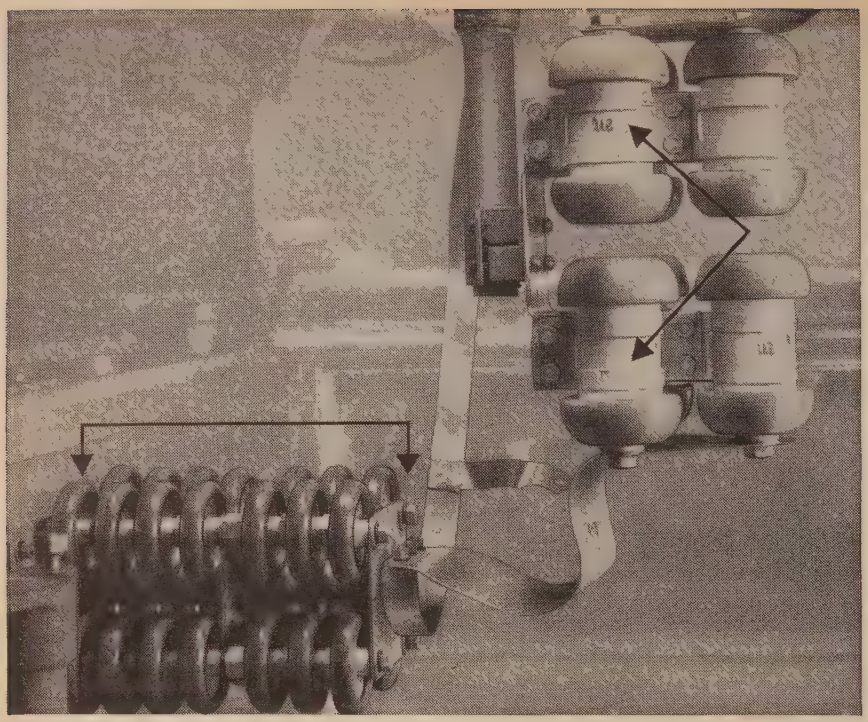

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(Continued from page 112A)

UNIVERSITY OF VERMONT, IRE-AIEE BRANCH

"Nuclear Reactors," by John Landis, Babcock and Wilcox; February 25, 1954.

VIRGINIA POLYTECHNIC INSTITUTE, IRE-AIEE BRANCH

"Electronic Computers in the Field of Guided Missiles," by Prof. E. Q. Smith, Faculty, Virginia Polytechnic Institute; March 9, 1954.

WASHINGTON UNIVERSITY, IRE-AIEE BRANCH

Film: "Bob White Through the Years"; February 25, 1954.

General meeting; March 5, 1954.

"Magnetic Amplifiers," by C. W. Lufcy, Naval Ordnance Labs.; March 11, 1954.

WAYNE UNIVERSITY, IRE-AIEE BRANCH

"Wayne University's Large Digital Computer," by Dr. E. P. Little, Wayne University; February 16, 1954.

UNIVERSITY OF WISCONSIN, IRE-AIEE BRANCH

"Transistors"—demonstration with apparatus by Mr. Burke, Bell Telephone Company; March 10, 1954.

UNIVERSITY OF WYOMING, IRE-AIEE BRANCH

General meeting; March 8, 1954.

The following transfers and admissions were approved to be effective as of April 1, 1954:

Transfer to Senior Member

Aaron, B. D., 3157 Sepulveda Blvd., Los Angeles 34, Calif.
Ainlay, A., 987 Wilfred St., Newtonbrook, (Toronto) Ont., Canada
Armstrong, R. W., Box 95, Neptune, N. J.
Blackstone, H., 2020 Jericho Tnpk., New Hyde Park, L. I., N. Y.
Broersma, C. B., 149 DeLairessestraat, Amsterdam Z., Netherlands
Bullock, M. W., 6805 Northwood Rd., Dallas 25, Tex.
Cappels, J. L., J. L. Cappels and Associates, 1214 W. Madison, Chicago 7, Ill.
Carlstrom, T. H., Sylvania Electric Products, Inc., Emporium, Pa.
Cheatham, T. P., Jr., Hosmer St., Marlborough, Mass.
Danner, R. V., 2512 Parkview Dr., Baltimore 7, Md.
Faran, J. J., Jr., 195 Marsh St., Belmont 78, Mass.
Flyer, I. N., 1419 Kanawha St., Hyattsville, Md.
Gayer, J. H, I.F.R.B. Palais Wilson, Geneva, Switzerland
Gershon, J. J., 2533 N. Ashland Ave., Chicago 14, Ill.
Goodstein, L. P., 98 Thayer St., New York 34, N. Y
Grebe, J. J., The Dow Chemical Co., Midland, Mich.
Gregory, C. A., 132 N. Wayne St., Arlington 1, Va.
Harrison, A. E., 5318 N. Zadell Ave., Temple City, Calif

(Continued on page 114A)

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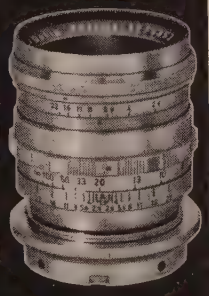

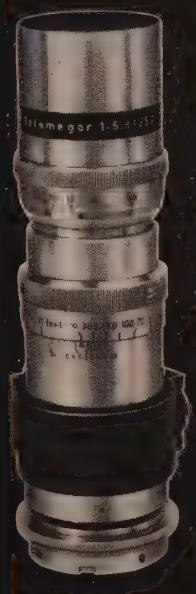

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4. Determine frequencies and frequency ratios, from 0 cps to 1 megacycle.
5. Provide a secondary frequency standard (stability, 1 part in 10^9).
6. Operate directly into (a) the new Berkeley Model 1452 single-unit printer, (b) Berkeley digital-to-analog converter, or (c) Berkeley data processer driving IBM card punches, electric typewriters, or teletype systems.

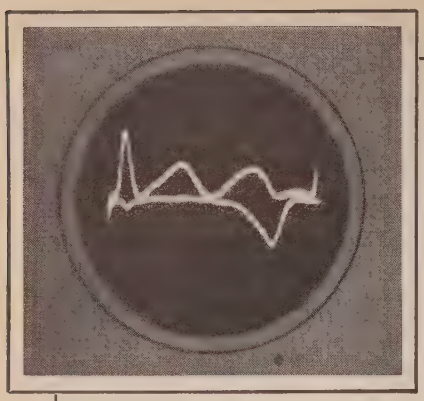
CONDENSED SPECIFICATIONS

Input Sensitivity: 0.2 v. rms (Freq. meas.); 1.0 v. peak to peak (other functions)
Input Impedance: 10 megohms shunted by 35 mmf.
Time Bases: 1 mc; 100, 10, and 1 kc; 100, 10 and 1 cps.
Gate Times: .00001, .0001, .001, .01, 0.1, 1.0 and 10 seconds
Crystal Stability: 1 part in 10^6 (temp. controlled)
Display Time: 0.2 to 5 seconds
Accuracy: ± 1 count, ± crystal stability
Power Requirements: 117 v. (± 10%), 50-60 cycles, 400 watts
Dimensions: 20¾" wide x 10½" high x 15" deep; panel, 8¾" x 19"
Price: Model 5510, $1,100.00 (f.o.b. factory).

Available for prompt delivery. Wire or write for technical bulletin, application data; please address dept. N5.

M32

Berkeley *division*
BECKMAN INSTRUMENTS INC.
2200 WRIGHT AVE., RICHMOND, CALIF.



MICROWAVE DEVELOPMENTS

Wheeler Laboratories is an engineering organization offering consulting and engineering services in the fields of radio and radar.

One test equipment developed by the Laboratories contains an oscillator sweeping over the 12% X-band at a rate of 12 round trips per second. As pictured above, an alternating pair of traces simultaneously display two characteristics of a network—in this case, the reflection coefficient and the differential phase angle in transmission; these are utilized while adjusting a multiple-tuned filter in a waveguide.

At present, Wheeler Laboratories includes a staff of twenty engineers under the personal direction of Harold A. Wheeler, a group of designers, and a model shop. Regular additions to the staff are continuing in order to keep pace with our expanding program of work.

Inquiries are welcomed regarding specialized problems in microwave design; a brief summary of our work is available on request.

Wheeler Laboratories, Inc.

122 Cutter Mill Road, Great Neck, N.Y.
HUnter 2-7876

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(Continued from page 112A)

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Munster, A. C., 175 Earl La., Hatboro, Pa.
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Patterson, G. W., 312 Dartmouth Ave., Swarthmore, Pa.
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Walsh, L., Box 491, Elizabeth 2, N. J.
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Zierdt, C. H., Jr., 554 Nottingham Rd., Syracuse 10, N. Y.

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Johnson, G. L., Acorn Hill Farm, Lyme, N. H.
Kay, L. M., 845 Riverside Dr., New York 32, N. Y.
Kenyon, D. E., 12 Monfort Dr., Huntington, L. I., N. Y.
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Kuchinsky, S., 503 City Line Ave., Phoenixville, Pa.
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Nieset, R. T., Biophysics Laboratory, Tulane University, New Orleans 18, La.
Reese, R. F., 909—21 St., Apt. 4, Santa Monica, Calif.
Robbins, R. W., 1341 South St., Geneva, Ill.
Rogers, W. E., Electrical Engineering Department, University of Washington, Seattle 5, Wash.
Sachs, R., 287 Pine St., Lowell, Mass.
Schramm, C. W., Bell Telephone Laboratories, Inc., Murray Hill, N. J.
Schweers, O. H., 5 Colony Cir., Camillus, N. Y.
Shekels, H. D., 4556 Pennlyn Ave., Dayton, Ohio
Shepard, D. H., Intelligent Machines Research Corp., 134 S. Wayne St., Arlington, Va.
Underhill, E. M., 274 Phelps Rd., Ridgewood, N. J.
Vadasz, A. J., Fayette Manor, Fayetteville, N. Y.
Yarbrough, S. L., Philco Corp., "C" and Tioga Sts. Philadelphia 34, Pa.

(Continued on page 119A)

AT LAST A TRULY FLEXIBLE AIR DIELECTRIC CABLE

ANDREW HELIAX

an important "for the first time" story...

There has long been cable easy to install.

There has long been highly efficient cable.

HELIAX is the first cable to deliver both characteristics. It is as flexible in application as solid dielectric cable, but has the sam
efficiency as copper air dielectric. HELIAX is superior in design, in efficiency and in electrical performance at microwave an
all lower frequencies, yet it is comparable in cost to lower frequency cables.

Ease of installation (HELIAX can be pulled through conduit and bent repeatedly without changing its characteristics) means substantial savings in installation costs.

HELIAX is crush proof, may be removed from one installation, coiled and reinstalled. Now available in 7/8" size in continuous lengths. Soon available in larger sizes. Send the coupon for detailed specifications.

ANDREW CORPORATION • 363 EAST 75th ST. • CHICAGO 19, ILL

Gentlemen:

Please send bulletin 70-A, giving technical details and specifications o
7/8" diameter flexible HELIAX cable (Type HX-0).

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POSITION_____

COMPANY_____

ADDRESS_____

when the job calls for controlled quality...

the call is for *Electra* deposited carbon resistors

Whether it is for hearing aids, guided missiles, or other electronic precision needs—the uniform quality of Electra carbon coat resistors is an important asset. Electra manufactures only one quality and it is the highest that can be humanly and scientifically produced.

Regardless of the carbon coat resistor need—we at Electra believe that only the highest grade resistor is safest, lowest cost to use. That's why Electra specializes in control of quality and exacting uniformity in every production detail.

This means Electra customers actually get more for their money—a more reliable component part for their product—a resistor whose rejection rate is practically nil. If you manufacture a quality product requiring a deposited carbon resistor, then—be sure—specify Electra.

8 SIZES: 1/8 watt to 2 watts and in two types—coated as well as hermetically sealed. MANUFACTURED TO SPECIFICATION MIL-R-10509A.

Electra carbon-coat PRECISION RESISTORS

Write for complete specifications

Electra **Manufacturing Co.**
2537 Madison Avenue
KANSAS CITY 8, MISSOURI

News—New Products

These manufacturers have invited PROCEEDINGS readers to write for literature and further technical information. Please mention your I.R.E. affiliation.

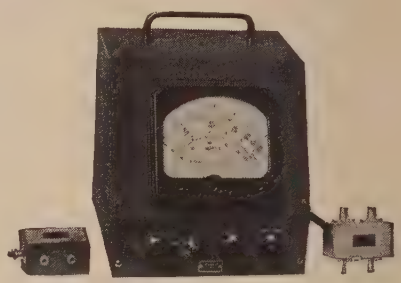
(Continued from page 78A)

position record equalizer that provides flat playback or equalization for the NARTB or AES recording curves. Speaker output impedances are 8 and 16 ohms.

Price is $79.50 net f.o.b. Chicago.

Phase & Victor Voltage Meter

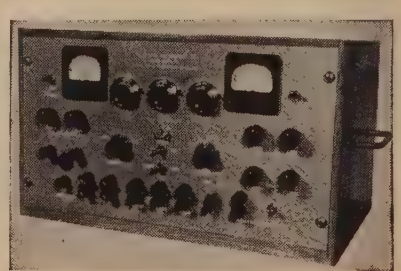
Type 213, developed by **Advance Electronic Co.,** P.O. Box 394, Passaic, N. J., is used to measure antenna systems, filters, amplifiers, and other high frequency transmission networks.



The frequency range of this instrument is 10 kc to 300 mc with jack-input probe and 100 kc to 1000 mc with coaxial probe. Both probes are supplied with the instrument. The lower limit may be extended to 1 kc with two additional coupling capacitors. The upper limit may be extended with relaxing accuracy. The input impedance is 2.5 $\mu\mu$f shunted with 100,000 ohms with jack-input probe. The coaxial input probe is normally designed for 50 ohms characteristic impedance. Differential values of characteristic impedance may be obtained upon request. The voltage range is 0.6, 60, and 6 volts full scale. The sensitivity can be lowered to 0.3 volt full scale through a simple adjustment. The phase angle range is 0–180 and 180–360° full scale. Ranges with better angular sensitivity may be arranged. The accuracy is ± 5 per cent nominal. The price is $228.00, F.O.B., Passaic.

Transistor Analyzer

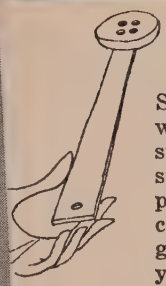
Polyphase Instrument Co., Bryn Mawr, Pa., announces its Model TA-2 Transistor Analyzer, a negative resistance and characteristic curve tracer which has been designed for use with a laboratory type oscilloscope. It will trace all negative resistance curves of both N-type and P-type point contact transistors.



(Continued on page 118A)

How to whip your relay problems!

Just call **ADVANCE...** *we're ready and willing to make the relays you need.*

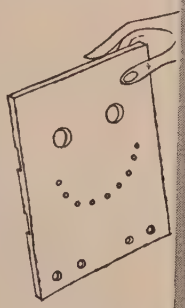


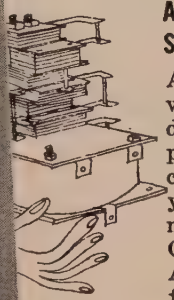
CONTACTS, *for example*

Say the word and we'll up the sizes... switch from fine silver to tungsten, palladium, silver-cadmium oxide, gold contacts, or... you name it!

INSULATION

Laminated phenolic, silicone glass, ceramic steatite, linen-base bakelite, or... you name it!





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We build a vast variety of sensitive, reliable relays...create superior new designs quickly. You'll find our quality control is now the tightest in the industry.

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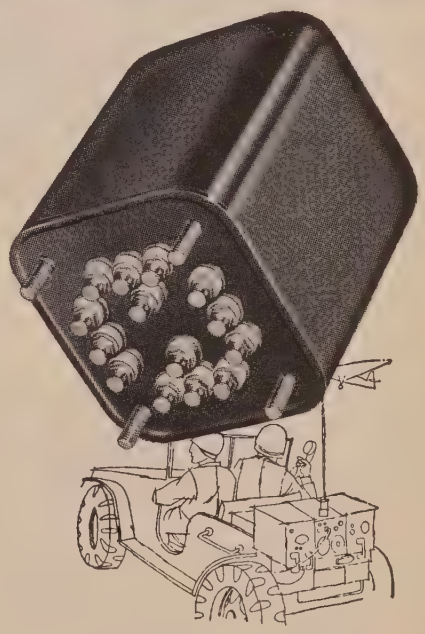
BURBANK, CALIFORNIA

Another First for CHICAGO

The complete MS (Military Standard) line of Hermetically-Sealed Power & Filament Transformers

CHICAGO TRANSFORMER now offers all units in the Military Standard (MS) line, as established jointly by the three armed forces (Army Signal Corps, Navy Bureau of Ships, and Air Force) working through ASESA (Armed Services Electronic Standards Agency) and in cooperation with the transformer industry. The complete line is housed in CHICAGO's one-piece drawn-steel cases. Outside case dimensions and mounting dimensions are within the tolerances of the Military Standard specification. Terminal arrangements and markings are also in accordance with the same specification. Tests conducted in the CHICAGO TRANSFORMER laboratories indicate that all units will meet the requirements of Grade 1, MIL-T-27 specifications for Class A operation. The Military Standard line should find wide usage in military airborne, marine, and ground communication equipment, and particularly for research and development applications, pilot runs and pre-production models.

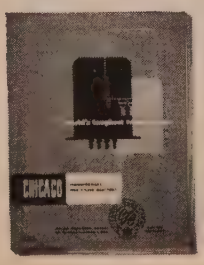


POWER TRANSFORMERS—INPUT REACTOR SYSTEMS (PRIMARY—105/115/125 V.—Frequency 54-66 cycles)

CATALOG NUMBER	MIL-T-27 PART NO.	HIGH VOLTAGE SECONDARY A-C Volts	D-C MA.	D-C V OUTPUT	RECT. FIL. Volts	Amps.	FIL. NO. 2 Volts	Amps.	WT. LBS.
PMS-70	MS-90026	200-100-0-100-200	70	385	6.3/5	2	6.3	3	4
PMS-70A	MS-90027	325-0-325	70	260	6.3/5	2	6.3	4	5
PMS-150	MS-90028	325-0-325	150	245	6.3	5	5	3	7¼
PMS-175	MS-90029	400-0-400	175	318	5	3	6.3	8	10
PMS-250	MS-90030	450-0-450	250	345	5	3	6.3	8	13
PMS-350	MS-90031	350-0-350	250	255					7½
PMS-550	MS-90032	550-0-550	250	419					11
PMS-800	MS-90036	800-0-800	250	640					16½

FILAMENT TRANSFORMERS (PRIMARY:—105/115/125 V.—Frequency 54-66 cycles)

CATALOG NUMBER	MIL-T-27 PART NO.	SECONDARY Volts	Amps	INSULATION VOLTS RMS	WT. LBS.
FMS-23	MS-90016	2.5	3.0	2500	1½
FMS-210	MS-90017	2.5	10	2500	2½
FMS-53	MS-90018	5.0	3.0	2500	1¾
FMS-510	MS-90019	5.0	10	2500	4
FMS-62	MS-90020	6.3	2.0	2500	1¾
FMS-65	MS-90021	6.3	5.0	2500	2¾
FMS-610	MS-90022	6.3 CT	10	2500	5
FMS-620	MS-90023	6.3	20	2500	8
FMS-210H	MS-90024	2.5	10	10000	4¾
FMS-510H	MS-90025	5.0	10	10000	7



Free "New Equipment" Catalog

You'll also want the full details on CHICAGO'S *New Equipment Line* of famous "Sealed-in-Steel" Transformers. Write for Free Catalog CT-153 today, or get it from your electronic parts distributor.

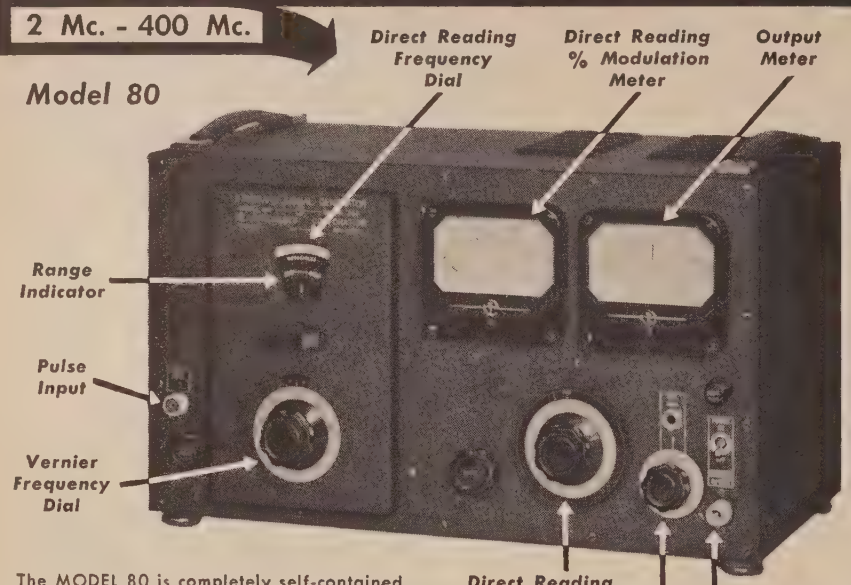
CHICAGO STANDARD TRANSFORMER CORP.

3501 W. ADDISON ST., CHICAGO 18, ILL.

STANDARD SIGNAL GENERATOR

2 Mc. - 400 Mc.

Model 80

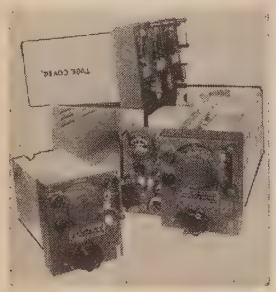


The MODEL 80 is completely self-contained, with built-in power supply and modulator. An extremely stable and accurate instrument for laboratory use, production testing or servicing.

Complete data on request

MEASUREMENTS CORPORATION

BOONTON NEW JERSEY



AN/APR-4 LABORATORY RECEIVERS

Complete with all five Tuning Units, covering the range 38 to 4,000 Mc.; wideband discone and other antennas, wavetraps, mobile accessories, 100 page technical manual, etc. Versatile, accurate, compact—the aristocrat of lab receivers in this range. Write for data sheet and quotations.

We have a large variety of other hard-to-get equipment, including microwave, aircraft, communications, radar; and laboratory electronics of all kinds. Quality standards maintained. Get our quotations!

NEW TS-13/AP X-BAND SIGNAL GENERATORS, with manual, $850.00 . . . TS-175/U Frequency Meters, 85-1,000 Mc., $625.00 . . T-47A/ART-13 Transmitters, $450.00 . . . and many more!

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Cosmic

CONDENSER SPECIALISTS SINCE 1923

ELECTROLYTIC and PAPER TUBULAR CONDENSERS FOR A.C. D.C. SETS PHONOGAPHS etc.

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Phone LUdlow 9-3360

News—New Products

These manufacturers have invited PROCEEDINGS readers to write for literature and further technical information. Please mention your I.R.E. affiliation.

(Continued from page 116A)

Since all circuit parameters controlling the negative resistance curves are available as metered variables on the front panel, the Model TA-2 enables the user to visually design any negative resistance circuit around a given transistor in a matter of minutes. The TA-2 will also determine the applicability of any transistor in a given negative resistance circuit.

In addition, the TA-2 will trace the collector characteristics, R_{22}, for both grounded emitter and grounded base connections, and transfer characteristic, R_{12}, of N-type and P-type point contact transistors and NPN and PNP junction transistors.

Transistor-Magnetic Microphone

By using transistors to make a diminutive pre-amplifier coupled with a high quality magnetic microphone, **Remler Company Ltd.,** 2101 Bryant St., San Francisco 10, Calif., has developed a new combination unit yielding improved speech intelligibility and reduced noise in radio and aircraft applications.

The transistors are built into the microphone unit in both straight microphone and handset applications. Units plug directly into existing equipment previously using carbon button microphones.

The amplifier, plastic seal-coated by a Remler-developed process, has been subjected to thousands of hours of life tests, humidity, hot and cold, and high altitude tests. It has also been flight-tested by a major airline. The microphone itself is the Remler ruggedized unit formerly supplied to the U. S. Navy and merchant marine service. Low noise and high intelligibility are attained without increasing the complexity and weight of the radio transmitter or announcing equipment. This new microphone is now in production at the San Francisco plant.

Variable Air Capacitors

Johanson Manufacturing Corp., Boonton, N. J., announces a new line of variable air capacitors. Small in size and of

(Continued on page 120A)

(Continued from page 114A)

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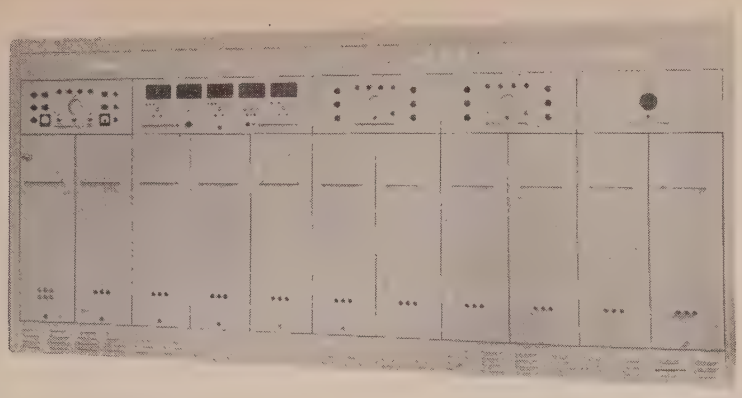
Birx, D. L., The Franklin Institute Laboratories, Philadelphia 3, Pa.
Bisso, R. J., 99 Meadow Rd., Emporium, Pa.
Bletcher, A. L., 4433 Silsby Rd., Cleveland 18, Ohio
Bogner, R. D., 67 Cheshire Rd., Bethpage, L. I., N. Y.
Boisvert, M., 1340 DuBuisson, Sillery, Que., Canada
Boxx, J. P., 2112 Westview Dr., Owensboro, Ky.
Brawner, C. A., Jr., 7207 Fisher Rd., R.F.D. 4, Dallas 14, Tex.
Brown, M. A., 308 Second Ave., Bradley Beach, N. J.
Buie, J. L., 11156 Valerio St., Sun Valley, Calif.
Cartwright, J. R., 52, Lytton Ave., Letchworth, Herts., England
Cerino, C., 53-10 Overbrook St., Little Neck 63, L. I., N. Y.
Couzens, D. T., S Tel O, RCAF AMCHQ, No. 8 Temporary Bldg., Ottawa, Ont., Canada
Davidson, R., Dynamic Electronics-New York, Inc., 73-39 Woodhaven Blvd., Forest Hills, L. I., N. Y.
Doe, U. P., Chief Engineer, Dept. of Broadcasting, 28 Windermere Crescent, Rangoon, Burma
Flynn, M. W., 24 Fairbank Ave., Toronto 10, Ont., Canada
Gilbert, E. M., 7512 Buena Vista Del Val, Burbank, Calif.
Gunderson, A. L., 3959 Highland Dr., Salt Lake City 7, Utah
Hansen, E., 2945 Alton Ave., Allentown, Pa.
Hinds, J. E., Jr., 6 Brompton Rd., Garden City, L. I., N. Y
Hix, E. W., 729 Noah, Akron, Ohio
Jahoda, J. R., 4321—69 St., Woodside 77, L. I., N. Y.
Khoury, N., 804 Brunswick Rd., Baltimore 21, Md.
Koppel, A. V., 60-01D—194 St., Fresh Meadows 65, L. I., N. Y.
Kunik, I. J., 521 Fifth Ave., New York 17, N. Y.
LaPointe, J. C., 5208—42 Pl., Hyattsville, Md.
Leadabrand, R. L., Radio Propagation Laboratory, Stanford University, Stanford, Calif.
Leskinen, J. I., 33-45—172 St., Flushing 58, L. I., N. Y.
MacDonald, C. F., 182 Duchess Ave., London, Ont., Canada
Magestro, J. C., 1783 Wacker Dr., Lancaster, Ohio
McDonough, V. T., 31-40—30 St., Apt. C4, Astoria, L. I., N. Y.
McKnight, C. J., c/o The American Legation, Tangier, Morocco
Murray, H. J., Jr., 703—37 Ave., Oakland 1, Calif.
Nail, P. L., 3028 E. 53 St., N. Kansas City 16, Mo.
Paulos, P. J., 11250 Playa St., Culver City, Calif.
Petersen, M. C., Box 605, Fair Oaks, Calif.
Reinish, G. B., 2709 Ocean Ave., Brooklyn 29, N. Y.
Rothauge, C. H., U. S. Naval Postgraduate School, Monterey, Calif.
Ruehl, G. C., Jr., 31 E. 21 St., Baltimore 18, Md.
Saxon, W. R., 1931 W. Vermont Ave., Phoenix, Ariz.
Smith, L. E., 12066 Milton St., Silver Spring, Md.
Spencer, N. A., 1 Terrace Cir., Great Neck, L. I., N. Y.
Velasco, F. C., 7835 S. Harvard, Los Angeles, Calif.
Winick, A. B., 1415 Oakridge Rd., Falls Church, Va.

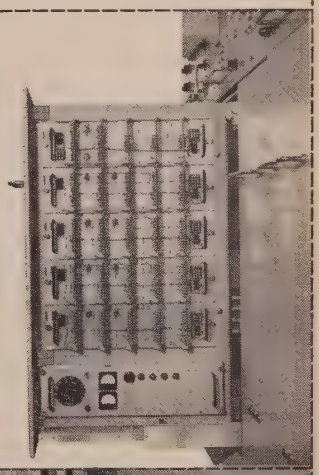
Admission to Member

Aasgaard, P. U., 1360 Ouimet St., Apt., 41, Ville St. Laurent, Que., Canada

(Continued on page 122A)



This compact equipment is accessible, too, because of this



This Time Signal Generator was built for Air Force ballistics testing by Vitro Laboratories, a division of Vitro Corporation of America. It employs Grant Industrial Slides.

We asked Charles K. Raynsford, project group leader, why Grant Slides were used here. His answer:

"Primarily for the convenience of the service technician. Each of those eleven sections contains approximately 150 vacuum tubes. Even with the low tube failure rate of 2% per 1000 hours, fast serving for preventive maintenance becomes quite important!

"In addition, this compactness would have been impossible without the Slides. In a conventional arrangement, the unit would have been twice as large."

"May we quote you on that?" we asked.

"Well," he answered, "say 'appreciably larger'. That would have increased the wiring capacitance, which, in turn, would have required more power to get the same band width."

"All in all, we're proud of this design," he added.

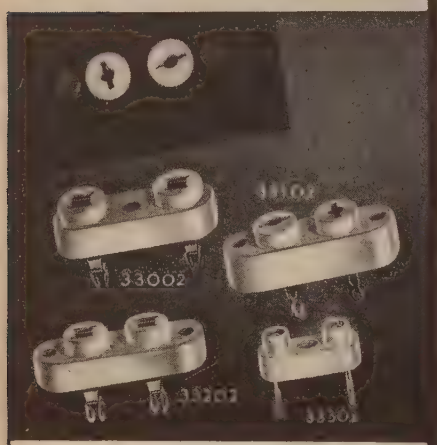
That makes it even. We're proud of our Slides.

Grant Industrial Slides

A product of Grant Pulley and Hardware Corporation
31-75 Whitestone Parkway, Flushing, New York

Designed for

Application



Crystal Holder Sockets
33002, 33102, and 33202
Plus new 33302 for CR7

In addition to the original 33002, 33102 and 33202 exclusive Millen "Designed for Application" steatite crystal holder sockets, there is now also available the new 33302 for the new CR7 holder. Essential data:

Type	Pin Dia.	Pin Spacing
33002.....	.125	.750
33102.....	.095	.500
33202.....	.125	.500
33302.....	.050	.500

JAMES MILLEN MFG. CO., INC.

MAIN OFFICE AND FACTORY
MALDEN
MASSACHUSETTS

News—New Products

These manufacturers have invited PROCEEDINGS readers to write for literature and further technical information. Please mention your I.R.E. affiliation.

(Continued from page 118A)

instrument quality, the new units feature simplicity of design for all uhf and rf applications. Single to four gang sections are standard units of the new line.

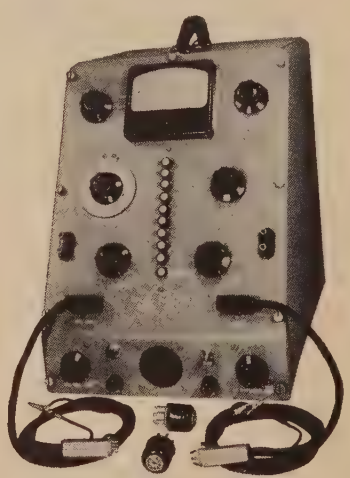
Each of the new capacitors has a frame which is swaged and soldered. Each has a three point mounting. The units are provided with hardened stainless steel ball bearings, and all rotors and stators of the units are soldered to further insure permanency. The wipers are of hardened, silver-plated, beryllium copper. The units are constructed of silver-plated brass (or invar) for low surface resistance.

The compression-loaded ceramic rod stator suspension gives it long dielectric path and mechanical reliability. A variable air trimmer is an integral part of each section in the unit.

For additional information, write Johanson directly.

Circuit Comparison Bridge

New general utility comparison bridge, trade marked "Circuit Matcher," improves efficiency in production testing of equipments, is available from **McShan Development Corp.**, 71 Murray St., New York 7, N. Y. Comparison of unknown circuits is made against those of a prototype of known performance.

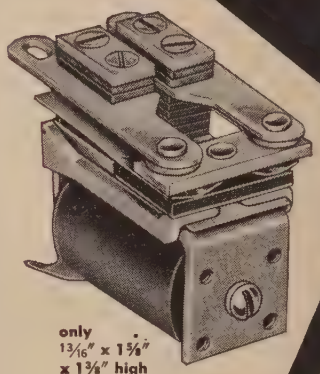


The instrument is equipped with two 9-conductor shielded cables terminated in noval plugs and mating adapters for other socket types. A pushbutton selector provides rapid comparison of similar circuits by ac or dc bridge tests. A point-to-point ohmmeter with flexible selector means is provided as a servicing aid. Two pair of panel terminals are provided for component comparison. A single balance control indicates circuit deviation from 75 to 30 per cent of normal impedance with 1 per cent accuracy over a wide range. Control simplicity enables non-skilled operation.

(Continued on page 128A)

Potter & Brumfield's MB3D miniature DC power contactor

for very high current applications

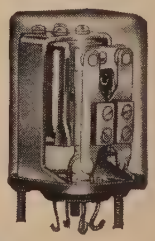


only 1 3/16" x 1 5/8" x 1 3/8" high

1. Massive, solid silver shorting bar works against ¼" dia. pure silver contacts.
2. Standard contact rating 60 amperes 28 volt DC non-inductive load.
3. Heavy brass contact arms with large, tinned solder terminals.
4. Over-travel spring of nickel silver spring temper for maximum contact pressure, long life.
5. Magnetic structure utilizes high permeability relay steel.

● Ideal for dynamotor starting applications in mobile, aircraft or marine equipment, and for other high current applications where space and weight are at a premium. Samples, recommendations and quotations on request.

ALSO AVAILABLE HERMETICALLY SEALED



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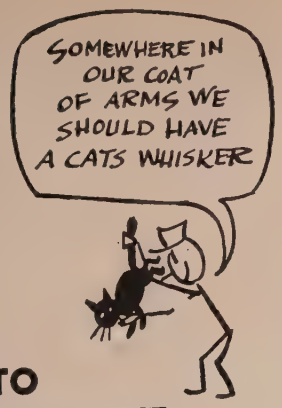

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- NO CARBONIZATION
- ZERO COLD FLOW
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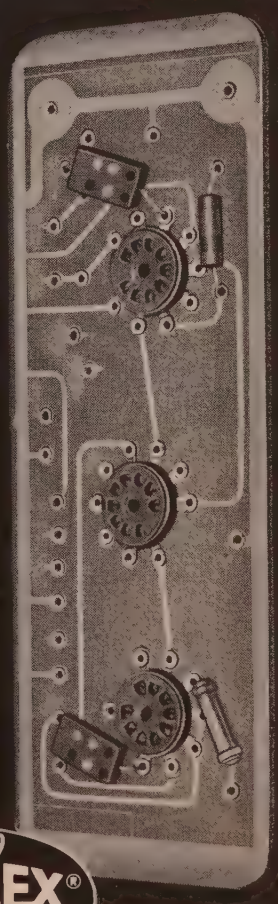

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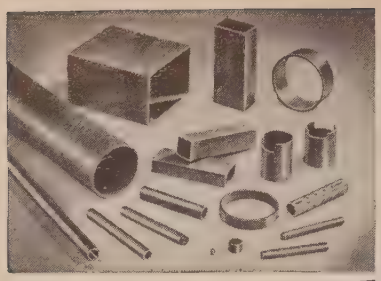

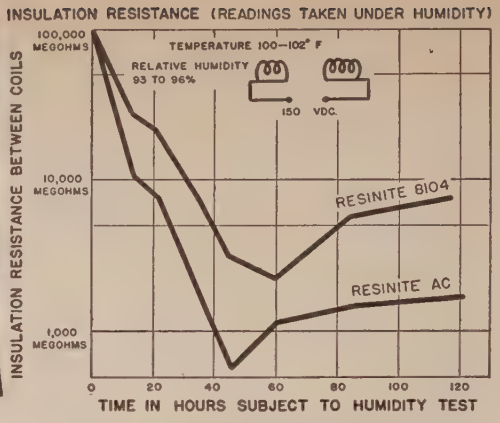

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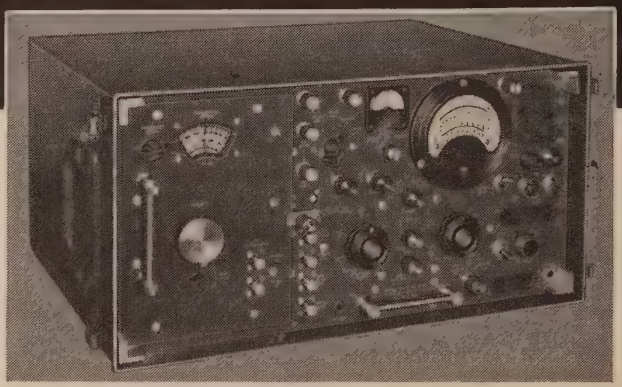

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Fletcher, M. W., 490 Ontario St., Toronto, Ont., Canada
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Gunst, M. A., Jr., 5851 Romany Rd., Oakland 18, Calif.
Haase-Dubosc, A. A., 135 Broadway, New York 6, N. Y.
Hansen, G. K. F., 81 Tornerosevej, Herlev, Denmark
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Hart, C. O., 803 S.W. 49 St., Oklahoma City 9, Okla.
Hawkins, W. E., James Millar Associates, Box 7116, Station C, Atlanta, Ga.
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Kaylin, R. R., 16-02 Everett Ter., Fair Lawn, N. J.
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(Continued on page 124A)

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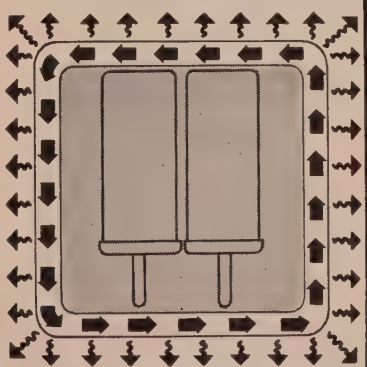
Thru "Thermaflow" Design*

Temperature, like water, seeks its own level. Instead of trying to "dam up" heat within the oven, by use of massive heat retaining elements, the JKO9 oven is designed to permit a uniform loss and uniform replacement of heat. Heat is simply replaced as it is lost from the low mass, high conductivity shell. And within this shell the crystal unit remains wrapped in a blanket of warm air. Because sufficient heat is always lost by the shell none need be yielded by the crystal.



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Symbol of Service

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AVAILABILITY

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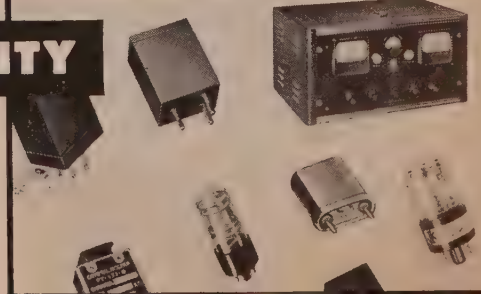
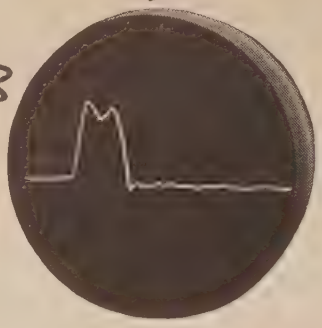
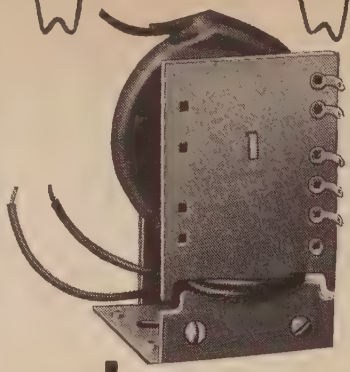


* A James Knights Patent.





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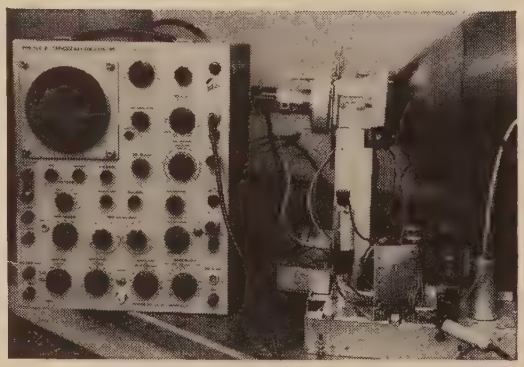

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(Continued from page 122A)

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Ponemon, R. D., Pyroferric Bldg., New York 67. N. Y.
Poucher, R. A., Jr., 3021 Mountain View Ave., Sacramento 21, Calif.
Powell, J. A., Frenchtown Porcelain Co., Frenchtown, N. J.
Ramsay, J. H., 219 Byberry Rd., Hatboro, Pa.
Rapaport, H., 315 E. Fifth St., New York 3, N. Y.
Richart, D. H., 1941 Fry Ave., Williamsport, Pa.
Robertson, H. E., 120 N. Broadway, Irvington, N. Y.
Roman, A. F., Electronic Tube Corp., 1200 E. Mermaid La., Philadelphia 18, Pa.
Rosencrantz, M. E., 4744 Overland Pkwy., Apt. 204, Toledo, Ohio
Samuel, R., 800 Riverside Dr., Apt. 6F, New York 32, N. Y.
Silverman, R. A., 15 Bryn Mawr Rd., Wellesley 81, Mass.
Steele, R. W., 401 Fourth Ave., White Sands Proving Ground, Las Cruces, N. Mex.
Stewart, A. C., 1200 E. Washington Ave., Council Bluffs, Iowa
Stewart, T. A., Federal Electric Mfg. Co., Ltd., 9600 St. Lawrence Blvd., Montreal 14, Que., Canada
Stone, N., 64 Wadsworth Ter., New York 33, N. Y.
Strongin, J., 86 Beekman Ave., Mount Vernon, N. Y.
Swanson, C. R., 3944 N. Greenview Ave., Chicago 13, Ill.
Tomlinson, C. C., 1721 Lansing Rd., Glen Burnie, Md.
Tullius, R. F., 5503 W. Lovers La., Dallas 9, Tex.
Tyndall, R. M., 423 W. Main St., Huntington, N. Y
White, C. E., 818 N. Harrison St., Arlington 5, Va.
Wild, T. A., York Rd., Timonium, Md.
Williams, R. O., 1515 Boutz, Las Cruces, N. Mex
Woerner, J. J., Box 185, Kentfield, Calif.
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Acker, A. E., 1244 Morningside Dr., Sunnyvale, Calif.
Acker, R. E., 2017 Parkview Ter., Spring Lake Heights, N. J.
Adamson, H. D., 6013 First Ave., N., Birmingham 6, Ala.
Ahmed, G., Lady Cunningham Maternity and Child Centre, Abbottabad, Pakistan
Akerberg, E. D., 181 Signal Depot Co., San Francisco, Calif.
Allgeyer, G. O., 7818 Flight Ave., Los Angeles 45, Calif.
Althaus, E. J., Hughes Aircraft Co., Florence & Teale Sts., Culver City, Calif.
Amey, G. B., 968 N. Market St., Williamsport, Pa.

(Continued on page 127A)

(Continued from page 124A)

Appleberry, C. M., Jr., Thomas Electronics, Inc., 118 Ninth St., Passaic, N. J.
Arkwright, L. P., 507 Burgess Dr., Rochelle Park, N. J.
Asher, J. W., Box 382, Corning, N. Y.
Baker, R. H., 375 Tavistock Blvd., Haddonfield, N. J.
Bartels, H. C., 2004 Hilltop Rd., Flourtown, Pa.
Bates, J. F., 1615 Annandale Rd., Falls Church, Va.
Beatty, W.H.,Valley Rd., R.F.D. 1, Millington, N.J.
Belcher, H. P., 2021 S. Fillmore St. Arlington 4, Va.
Bennett, H. L., Box 460, Independence, Kans.
Biss, M. W., 15 Transit Dr., Greenock Heights, McKeesport, Pa.
Black, V. G., 2619 Post Oak Rd., Houston, Tex.
Blocki, W. G., Box 609, Matteson, Ill.
Bolinger, N. C., 2613 Indiana, N.E., Albuquerque, N. Mex.
Bost, M. E., 738 Harvard St., Houston 7, Tex.
Bowman, J., Jr., Technical Representation, 25 S. Easton Rd., Glenside, Pa.
Boyd, J. E., 146 Wayne Pl., S.E., Washington 20, D. C.
Bradley, F. B., 1507 Auline La., Houston 24, Tex.
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Bratt, M. J., 7462 Werner Ave., Cincinnati 31, Ohio
Breese, H. R., 139 Twin La., N., Wantagh, L.I., N. Y.
Brown, M. H., 12521 S. 68 Ct., Palos Heights, Ill.
Brush, J. W., Grubb Rd., Chevy Chase, Washington, D. C.
Bryant, B. L., Jr., 263-23—73 Ave., Glen Oaks, L. I., N. Y.
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Carpenter, L. B., 1624 Kingsway Rd., Baltimore 18, Md.
Carter G. D., 1300 Oakland Rd., Apt. 1210, Cedar Rapids 1, Iowa
Cass, J., 2032—61 St., Brooklyn 4, N. Y.
Chappell, I. C., 69 Page St., Avon, Mass.
Checketts, E. L., 9300 Claremont Ave., Albuquerque, N. Mex.
Chew, Y. P., 361 King St., E., Toronto, Ont., Canada
Chisholm, T. L., Jr., 110 W. 75 St., New York 23, N. Y.
Chu, T., 904 N.W. Eighth St., Oklahoma City, Okla.
Church, A. K., 800 S. 31 St., Lincoln 8, Nebr.
Clayton, M. W., Earle Hotel, Rm. 304, Biloxi, Miss.
Cleary, R. E., Commercial Bank Bldg., Berea, Ohio
Cooper, C. B., 648 Second Ave., S., St. Petersburg 5, Fla.
Corning, J. J., 95 High St., Newburyport, Mass.
Couch, L. J., Box 406, Emporium, Pa.
Crowley, J. F., 32 Bay Ave., Bloomfield, N. J.
Damme, M. C., 545 N. Grandview Ave., Papillion, Nebr.
DasGupta, A. K., Solar Manufacturing Corp., Seville & E. 46 St., Los Angeles, Calif.
Davis, B. L., Division 12.5, National Bureau of Standards, Washington 25, D. C.
Davis, R. B., Jr., A. D. Cardwell Electronics Products Corp., Plainville, Conn.
Davisson, C. O. C., 616 W. Cliveden St., Philadelphia 19, Pa.
DeLisle, A. B. C., Box 3216, USAF, Wright-Patterson AFB, Ohio

(Continued on page 150A)

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- Encapsulated components—shock mounted meters recessed behind plastic windows.
- No tubes—lamps—carbon piles—commutators or moving contacts and no radio interference.
- STAVOLT Rectifiers are rugged and can take heavy intermittent overloads. They meet MIL-E-7894 with both range and characteristics.

Engineers will be interested in learning more about the fine qualities of STAVOLT Rectifiers. Write for catalog with detail specifications.

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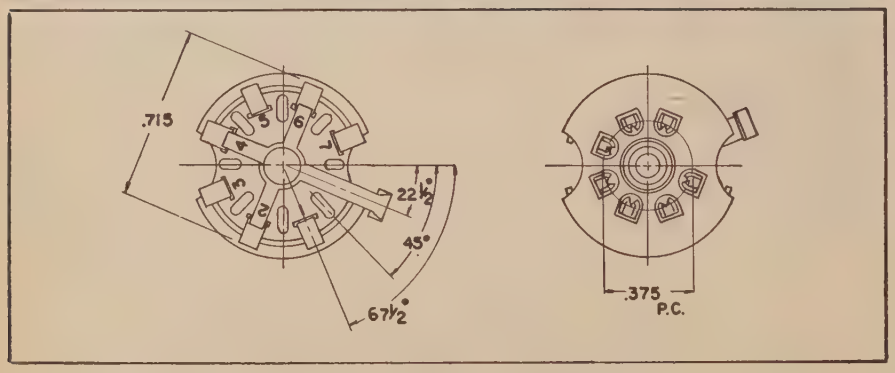
25th Year of Rectifier Manufacturing

EBY PRINTED CIRCUIT SOCKETS

7 pin miniature for printed circuit use—the first of a series of high quality sockets for this service—by EBY



FEATURING

1. ***Snap action contacts fit snugly in chassis and hold socket in place for subsequent operations.***
2. ***Spoke-type (non-rotating) center ground permits positive grounding of any or all contacts as specified.—also provided with long strap for grounding tube shield.***
3. ***Available in low-loss or general purpose phenolic.***
4. ***Contacts of standard miniature type design, with two lengths of tail—long for top side soldering.***

Write for complete details, and samples.

HUGH H. EBY INC.

4702 Stenton Ave., Phila. 44, Pa.

News—New Products

These manufacturers have invited PROCEEDINGS readers to write for literature and further technical information. Please mention your I.R.E. affiliation.

(Continued from page 120A)

Thomas to Produce Lawrence Color Tube

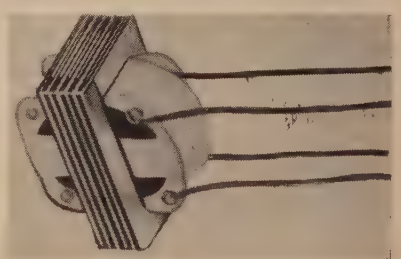
Thomas Electronics' President, Thomas L. Clinton, and Richard Hodgson, President of **Chromatic Television Laboratories, Inc.,** signed an agreement recently which made Thomas the first licensed producer in the East of the Lawrence single-gun color TV picture tube.

The single gun principle, in addition to permitting a large picture, is both cheaper to produce and simpler to operate than the three-gun color tube previously experimented with by other companies. Hodgson stated that the Lawrence color tube gives a higher quality black-and-white reproduction than the standard telecasting of today. Thomas Electronics said that the Lawrence tube is one-third larger than present color tubes but that its design will permit its use in cabinets much smaller than those used previously for color sets.

Pioneers in making TV tubes, Thomas now uses 225,000 square feet of space in its Passaic, N. J. plant. The company turns out about 30 different types of black-and-white tubes.

Transistor Transformers

Microtran Co., Div., Crest Laboratories, Inc., announces the introduction of a line of miniature transformers, designated as Veri-Miniature transformers.



Designed for use in transistorized circuitry, such as hearing aids and radio paging, they are available for input, interstage, output, and choke applications. Their size is: $\frac{7}{16} \times \frac{1}{2} \times \frac{7}{16}$ inches. Frequency response is flat from 300 to 10,000 cps.

In addition to standard items, special units can be constructed upon receipt of customer's specifications. Full details are available upon request.

(Continued on page 153A)

electronic engineers

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2. Circuit design
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5. Field Engineers

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News—New Products

These manufacturers have invited PROCEEDINGS readers to write for literature and further technical information. Please mention your I.R.E. affiliation.

(Continued from page 128A)

Selenium Rectifier Handbook

The design, application, specification and circuitry of selenium rectifiers are described in comprehensive fashion in the second edition of *Federal's Selenium Rectifier Handbook*, published by the **Federal Telephone and Radio Co.**, Clifton, N. J., a division of the International Telephone and Telegraph Corp.

This enlarged edition of 80 pages treats in detail a representative listing of the selenium rectifiers manufactured by Federal for radio and television use. It also covers rectifiers designs and power-supply circuits for such applications as phonographs, audio amplifiers, amateur radios, mobile radios, photocell amplifiers, intercommunication systems and other dc power supply requirements.

In addition, servicing information for the selenium rectifier is presented in simple, easy to follow, form. Trouble-shooting tables help the serviceman or technician follow proper procedures in checking a rectifier when specific circuit conditions exist, and diagrams are provided for setting up forward and reverse current tests in the repair shop.

Available through distributors or from Federal direct, price $.50.

(Continued on page 144A)

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Today's horizons in electronic engineering are limited only by the vision of the individual himself. To those qualified men who desire to stand on the constantly changing frontiers of electronic development, we offer a chance to pioneer and grow with a soundly-established, yet young and progressive company.

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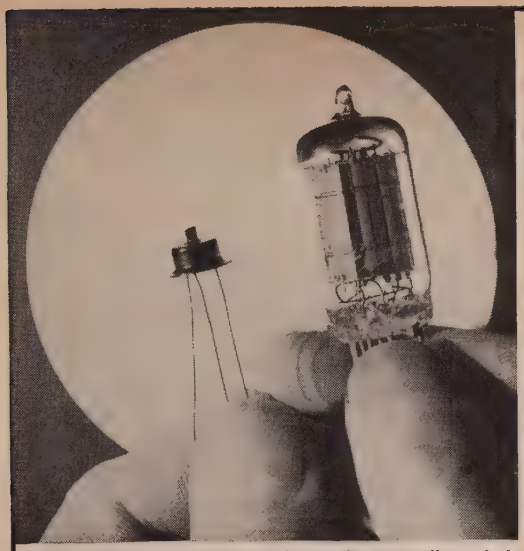

General Electric hermetically-sealed, evacuated junction transistor (left), contrasts sharply with standard miniature vacuum tube.

WORK ON THE NEW DEVELOPMENTS IN ELECTRONICS AS AN ENGINEER WITH GENERAL ELECTRIC

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For theoretical development of wave propagation in the region of dielectric and metallic boundaries, including methods of approximation in boundary value problems, investigation of propagation through ferrites, development of specialized microwave components and application to systems. Complete digital and analogue computing facilities are available. Applicant must also be capable of directing an experimental program for verification of theoretical predictions.

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To conduct basic experiments on microwave circuits and transmission through dielectric media, and to aid in the design of basic test equipment involving novel principles. A good understanding of electromagnetic wave phenomena is essential. Recent graduates with outstanding records will be considered.

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HIGH VACUUM SYMPOSIUM

Of particular interest to individuals and organizations concerned with applications of high vacuum technology is the High Vacuum Symposium which will be held June 16, 17, and 18, 1954, at the Berkeley-Carteret Hotel at Asbury Park, N.J.

The program, comprising approximately 25 technical papers, will include discussion of Nomenclature and Standards; New Equipment and Instruments; Fundamental Developments in Vacuum Technology; Methods and Techniques; Applications and Processes. In planning the program, special care has been taken to include subjects of practical as well as theoretical importance.

Persons and organizations interested in attending the Symposium are invited to write The Committee On Vacuum Techniques, Box 1282, Boston 9, Mass.

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If you qualify for one of the present five openings in the Gilfillan air defense program, please send a brief resume to R. E. Bell, Gilfillan Bros., 1815 Venice Blvd., Dept. 524, Los Angeles, Calif. Conference covering details will be arranged at a convenient location.

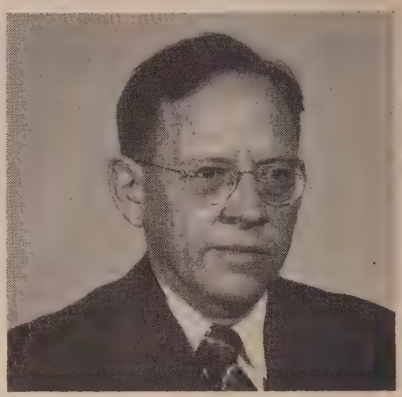
Gilfillan
LOS ANGELES

News—New Products

These manufacturers have invited PROCEEDINGS readers to write for literature and further technical information. Please mention your I.R.E. affiliation.

(Continued from page 142A)

Ford Appoints Slawson



Kenneth Slawson has been appointed Assistant to the President of **Ford Instrument Co., Div. of The Sperry Corp.** Mr. Slawson started with Ford in 1928. In 1946 he was appointed Manager of E. G. Staude Manufacturing, a new Sperry Corp. Division. As a result of the recent sale of Staude to the Bryant Chucking Grinder Corp., Mr. Slawson returned to Ford to his present position.

(Continued on page 146A)

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YORK Division . . . . .

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At Airtron, Inc., you do stimulating design and development work on high frequency components in company with like-minded engineers and professional men, in a firm geared to an engineer's point of view. You work in Linden, N. J., in the Metropolitan New York area—surrounded by some of the finest suburban and country living facilities—only a short half-hour from the New Jersey shore, and the same distance from New York City.

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to design and develop high-speed electronic digital equipment involving vacuum-tube and magnetic-core circuitry.

Openings are also available for outstanding candidates to work on use of high-speed digital computers to control large physical systems. This involves study of control requirements of the whole system, reduction of them to a simple pattern of control instructions, and translation of the pattern into computer code. Position requires appreciation for physical systems, ingenuity, and imagination. Training in computer principles provided as necessary.

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News—New Products

These manufacturers have invited PROCEEDINGS readers to write for literature and further technical information. Please mention your I.R.E. affiliation.

(Continued from page 144A)

New Switch

A new, precision, low loss 18-position Brown-Hill ceramic wafer-type, selector switch is now being produced by **R-F Electronics, Inc.,** 291 N.E. 61st St., Miami, Fla.

These Brown-Hill switches have a voltage break-down of 4,000 volts ac (peak at 60 cps). They have a current carrying capacity of 30 amperes. at 60 cps, non-inductive load.

Available in 1 to 6 gangs, this switch features a new, 20-degree positive detent mechanism with adjustable stops. Its rotor blades and stator contacts are of solid pure silver, providing low resistance, high current capacity, thereby entirely eliminating receiver switching noise. The rotor contact segments do not depend on rivets to carry current. A hole and cutouts are provided in the contact for mechanical attachment of wires.

Dow-Corning-200 impregnated steatite (Grade L-5 or better), is used throughout. These switches are available either bushing or stud mounted. A life test of one million revolutions shows a minimum of wear. Additional information about the switches can be secured from the manufacturer.

(Continued on page 153A)

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Minimum 3 to 5 years experience in X band and K band microwave antenna systems for precision fire control and missile radar.

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These are top level positions for men who wish to apply their knowledge and abilities in a company organized, owned and run by engineers, where the only limitations are the individual's ability and ambition.

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Please forward complete resume to attention of Mr. E. Eddy

Guided Missiles Division
99 Church Street
New York 7, New York

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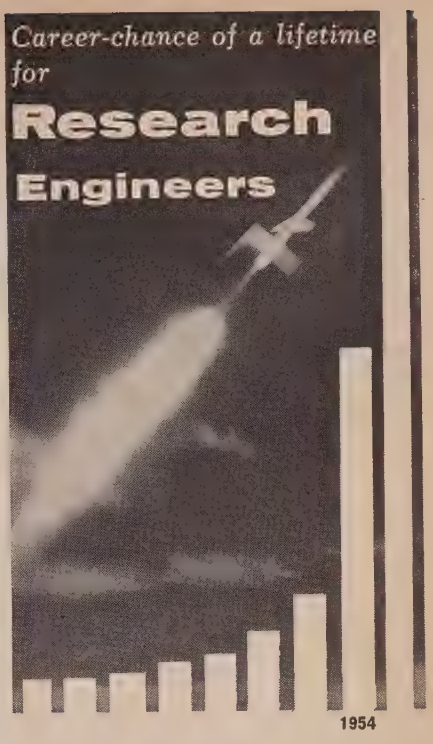

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my field of engineering

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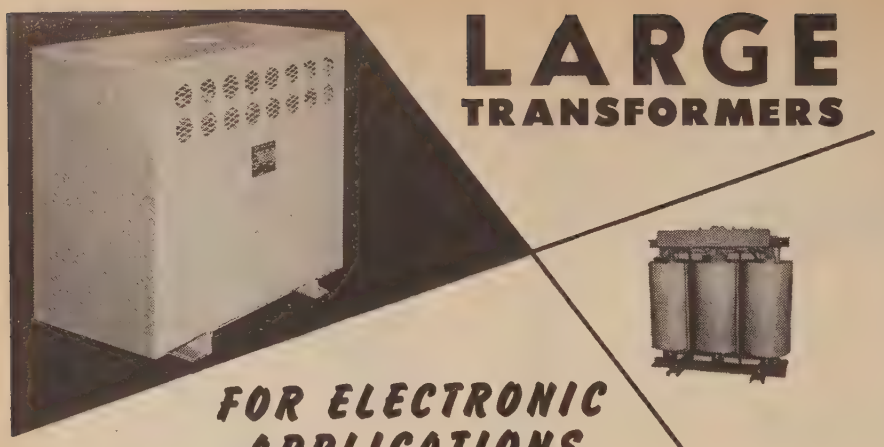

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is an operating unit of
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(Continued from page 127A)

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DeRosa, G. D., 2801 Mocking Bird La., Midwest City, Okla.
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Doll, C. L., 8012 State St., Ralston, Nebr.
Douglas, J. E., 7003 Housman St., Houston, Tex.
Drake, E. S., 117 Highview St., Mamoroneck, N. Y.
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Fenton, A. P., Box 421, San Leandro, Calif.
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Freidner, G. E., 590 Buffalo St., Conneaut, Ohio
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Gebhard, R. F., A.M.R.L., Fort Knox, Ky.
Geismar, G. G., 99-05—63 Dr., Rego Park, L. I., N. Y.
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Gerlach, R. K., 227A Princeton Rd., Haddonfield, N. J.
Gibson, J. J., Huskvarnavagen 122, Vattersnas, Sweden
Gilmore, E. H., 2nd AACS 1 M Sqdn., APO 942, c/o Postmaster, Seattle, Wash.
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Gunn, D. L., 828 S. Fairfield Ave., Lombard, Ill.
Gyllenhaal, P. R., Woodland Rd., Bryn Athyn, Pa.
Hamlin, W. O., R.F.D. 2, Emporium, Pa.
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Harris, C. F., 111 E. Rische, San Antonio 4, Tex.
Harrison, H. W., 3999 McKenzie Ave., Mt. View, Calif.
Hartnett, J. D., 113 Osborn Rd., Aberdeen, Md.

(Continued on page 151A)

(Continued from page 150A)

Hayes, P. J., RCA International Division, 30 Rockefeller Plaza, New York 20, N. Y.
Henkel, G. F., 3958 W. 157 St., Cleveland 11, Ohio
Herscher, M. B., 5832 Pemberton St., Philadelphia 43, Pa.
Hicks, F. M., Box 599, Houston 1, Tex.
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Mahoney, T. J., Jr., 3323—71 St., Jackson Heights 72, L. I., N. Y.
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Mann, H., 141-11—78 Rd., Flushing, L. I., N. Y.
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Martin, H. F., 20 Hensley Ave., San Jose 12, Calif.
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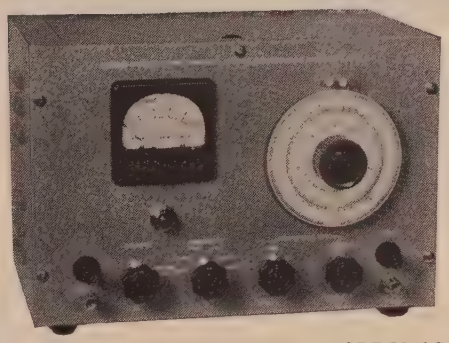
(Continued on page 152A)

NEW

PHASE DETECTOR
0.1° ACCURACY
10 KC to 10 MC



TYPE 205

Type 205 Phase Detector was developed to meet the increasing need for detecting phase angle with error of less than 0.1 degrees in high frequency communications systems such as color television.

PRICE $445

SPECIFICATIONS

ACCURACY: ±0.1 degree in phase reading or ±1 % of the time delay indicated on the dial of the continuously variable delay line.

RESOLUTION TIME: 5 x 10^{-10} seconds or smaller; the smallest phase angle can be read on the dial is approximately equal to 5 x 10^{-10} x 360 x frequency in degrees.

TIME DELAY: Three continuously variable delay lines are supplied with the unit, 0 to 0.45 microsecond, 0 to 0.25 microsecond and 0 to 0.05 microsecond.

PHASE RANGE: The maximum phase range is equal to the total time delay of the continuously variable delay line multiplied by the frequency of the signals and 360.

INDICATOR SENSITIVITY: Approximately 0.02 volts full scale maximum.

INPUT IMPEDANCE: Two low capacity probes with input capacitance less than 4 mmf are supplied with the unit. The panel binding posts have about 1 megohm shunted with 12 mmf on both input channels.

CONTINUOUSLY VARIABLE TIME DELAY
PASSIVE NETWORK
with 5x10-10 SEC. RESOLUTION TIME

TYPE 2011

Type 2011 is well suited for variable time delay for pulses, phase angle measurement, or time measurement of various signals.

- Faithful reproduction of the wave form of the input signal.
- No time jitter.
- No limit on repetition rate.
- Millimicrosecond accuracy.

PRICE $245

SPECIFICATIONS

TIME DELAY: Continuously variable from 0 to 11 microseconds.
RESOLUTION TIME: Less than 5x10^{-10} seconds from 0 up to 11 microseconds.
BANDWIDTH: 10 cps to 15 megacycles when the step variable delay line is at its off position (time delay from 0 cps to 0.5 usec), otherwise, the upper limit decreases to 3.3 megacycles.
RISE TIME: Less than 10% of the time delay at any point.
INPUT IMPEDANCE: 20 uuf shunted with 1 megohm direct.
OUTPUT IMPEDANCE: 300 ohms nominal.
ACCURACY OF TIME DELAY: Maximum error less than 10^{-9} second or 0.1% of the time delay at any point with correction curve, otherwise, one percent of the time delay at any point.

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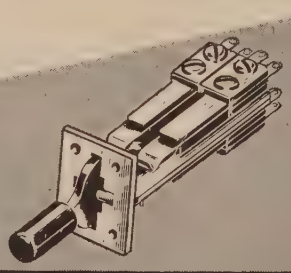

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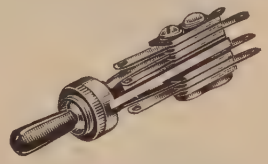

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News—New Products

These manufacturers have invited PROCEEDINGS readers to write for literature and further technical information. Please mention your I.R.E. affiliation.

(Continued from page 146A)

New Helipot Plant

Helipot Corp., precision potentiometer manufacturer, opened a new eastern plant in Mountainside, N. J. on December 3. Helipot, with headquarters in South Pasadena, Calif., is a division of Beckman Instruments, Inc.

IF Transformer

J. W. Miller Co., 5917 S. Main St., Los Angeles 3, Calif., is producing subminiature k-tran IF transformers under a patent arrangement with Automatic Mfg. Corp.

The component supplies 455 kc. Through the use of a ferrite shell core material, these sub-miniatures claim the gain and bandwidth characteristics previously obtained in larger assemblies. Dimensions ½ inch square×1½ inches high. Price $2.50.

(Continued on page 159A)

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Roth, J. F., 4610—31 Rd., S., Arlington, Va.
Rupp, V. T., 2230 W. 11 St., Los Angels 6, Calif.
Ruscus, P. V., 141 Jacques St., Elizabeth 4, N. J.
Ryan, C. C., Westinghouse Electric Corp., Baltimore, Md.
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Schafer, E. L., 5312 Holly St., Bellaire, Tex.
Shearer, W. G., Box 3116, MCLI, Wright-Patterson AFB, Dayton, Ohio
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Schloemer, R. D., 1024 McMorris Dr., Honolulu 18, Hawaii
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Simpson, R. H., 2209 David Dr., Fort Worth 11, Tex.
Skaritka, R., 1919 S. Shelby Ct., Chicago 8, Ill.
Skinner, A. M., Computation Laboratory, Wayne University, Detroit 1, Mich.
Smith, R. W., 2220 Savannah Ter., S.E., Apt. 33, Washington 20, D. C.
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(Continued on page 156A)

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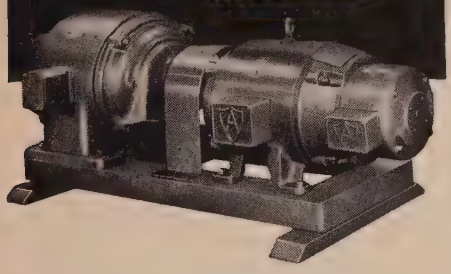

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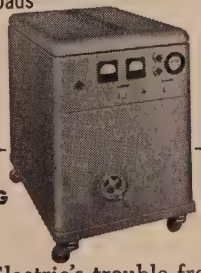

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Employs American Electric's trouble-free, direct-connected, high-cycle Inductor Alternator—which has no commutator, slip rings, brushes, springs etc. Output is rectified and exciter-regulated by either electronic or magnetic amplifier means then returned to the rotating field of the main alternator.

Write for details and quotations outlining requirements.

Also Manufacturers of High Frequency Inductor Alternators, Miniature Electric Motors, A.C. Industrial Motors, Motor Driven Blowers and Fans.

AMERICAN ELECTRIC MOTORS TRADE MARK

4811 Telegraph Road, Los Angeles 22, California

FIELD ENGINEERING REPRESENTATIVES:

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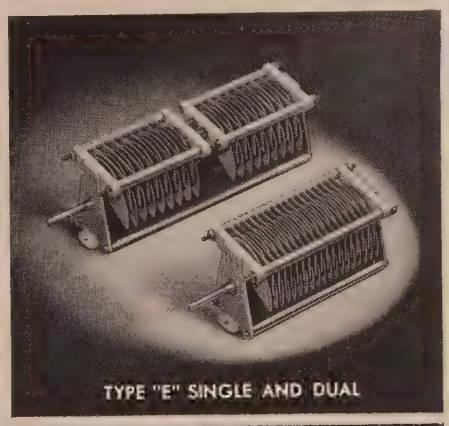
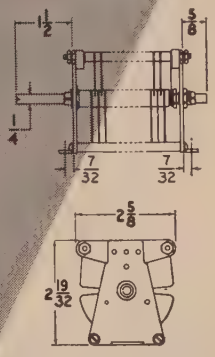
SPACE TIGHT?

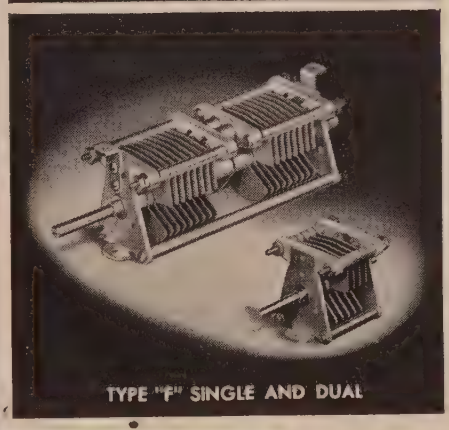
USE JOHNSON TYPE "E" AND "F" CAPACITORS



TYPE "E" SINGLE AND DUAL

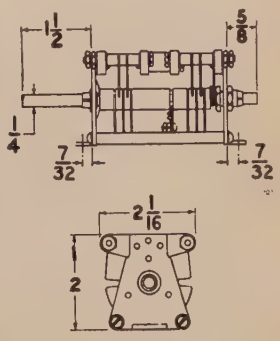
TYPE E





TYPE "F" SINGLE AND DUAL

TYPE F



LESS PANEL AREA—MORE CAPACITY PER CUBIC INCH!

Rugged and compact, JOHNSON Type "E" and "F" Capacitors for medium and low power transmitters have more capacity per cubic inch and occupy less panel space for their ratings than any other capacitor on the market.

In addition, their simple, functional design provides extremely good mechanical rigidity and exceptional electrical performance. Points of superiority include steatite insulation, heavy aluminum plates with rounded edges, ¼" tie rods, efficient cadmium plated phosphor bronze contacts and stainless steel shafts. Designed for chassis mounting with stator up to reduce capacity to ground, they have unusually low minimum capacity ratings.

Available in 40 different models, single and dual sections with plate spacings from .045" to .125" and maximum capacities from 35 to 500 mmfd.

Quotations on standard models or "specials" to meet your requirements are available on request. Investigate Johnson's "E" and "F" Capacitors for your application today!

Write for the **NEW** *Johnson General Products Catalog #975.*

E. F. JOHNSON COMPANY

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At the I.R.E. Show, thousands discovered

the pioneer is the leader

HOW TO SAVE MONEY BY SAVING TIME with PANORAMIC EQUIPMENT

NEW

Panoramic Model

SPA-1

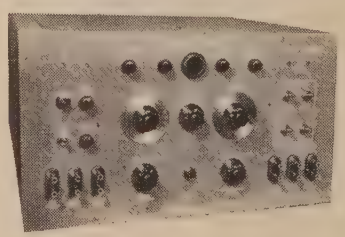
A Convenient Single Package VHF-UHF Spectrum Analyzer 50MC-950MC with Two RF Tuning Heads

The Model SPA-1 Panoramic Spectrum Analyzer incorporates a superior panoramic indicator, power supply and optional tuning heads, RF-2 and RF-3, which cover the ranges between 50MC to 250MC and 220MC to 950-MC, respectively. The SPA-1 is exceptionally simple to operate, lending itself to production tests as well as laboratory usage.

- Direct Frequency Reading, 50MC-950MC
- 10MC scanning width, continuously reducible to almost 0KC
- Variable resolution 9KC-100KC
- Sweep Rates: 1cps, 5cps, 30cps and 25-35cps, variable
- High inherent stability
- Low cost

Model SG-1 Panoramic's NEW Sweep Generator

for Accurate Inspection of Responses of Sonic and Ultrasonic Systems and Devices



- Direct Frequency and Amplitude Reading Screen for Slave Scopes
- Frequency Range: 40cps-20KC, 400cps-200 KC, selectable
- Frequency scales: logarithmic or linear, selectable
- Scan rate 1cps internal; 60-0.04cps external with Model TW-1 Triangular Wave Generator
- Amplitude scales: linear or 2 decade logarithmic, selectable
- Variable linear sweep range
- Internal frequency markers

Panoramic's NEW

Model FM-1 FM Monitor

A Low Cost Portable Package for Rapid Visual Measurement of *Actual Bandwidth* of Mobile FM Transmissions

- Instantaneous panoramic presentations of carrier and sidebands of voice transmissions
- Helps prevent channel spillover
- Indicates modulation symmetry
- Accurately measures deviation by constant tones
- Has band limit markers
- Simple to operate

PANORAMIC RADIO PRODUCTS, INC.

If you missed the I.R.E. Show—and the PANORAMIC demonstration of these new and important instruments at work on actual problems—write today for full information.

Inquiries invited on Panoramic Spectrum Analyzers for Special Problems.

12 South Second Avenue, Mount Vernon, N.Y.
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(Continued from page 154A)

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Thomson, W. E., 5803 Bellaire Blvd., Bellaire, Tex.
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Underwood, M. D., River Rd., Chatham, N. J.
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Wedge, T. E., 3500 Dudley Ave., Baltimore 13, Md.
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Wright, R. H., 6903 Schiller, Houston, Tex.
Zapalac, R., 2227 Mimosa La., Houston, Tex.
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Ziller, M. A., 4 Holland Ave., Long Branch, N. J.
Zimmerman, D. N., 306 Comm Recon Bn., Fort Bragg, N. C.
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News—New Products

These manufacturers have invited PROCEEDINGS readers to write for literature and further technical information. Please mention your I.R.E. affiliation.

(Continued from page 153A)

New Book

N. V. Philips Gloeil Ampenfabriken, Eindhoven, the Netherlands, has published and released a new book through their Technical and Scientific Literature Dept.

The work, **Industrial Electronics,** has 250 pages, 266 illustrations, and was written by Dr. R. Kretzmann.

The book is divided into two basic sections: Part I which lists 9 types of tubes and their basic circuits, and Part II which lists 10 electronic devices for industrial purposes.

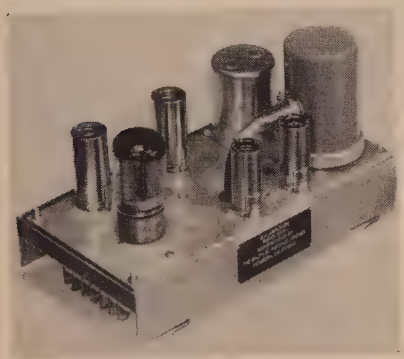
This manual is distinguished from other books in the field in that it is largely devoted to a detailed descriptive study of successful modern devices of many types employing electronics, with numerous circuit diagrams and photographs.

Distributors: Messrs. Elsevier Press, Inc., 402 Lovett Boulevard, Houston 6, Texas, and 155 East 82 St., New York 28, N. Y.

Price: U.S.A. $5.50.

DC Amplifier

The Model 3501 DC Amplifier, manufactured by **The Ralph M. Parsons Co.,** 135 W. Dayton St., Pasadena, Calif., has been designed for use in special purpose computer circuits. Dc drift, noise, and nonlinearities have been held to an absolute minimum, the firm claims.

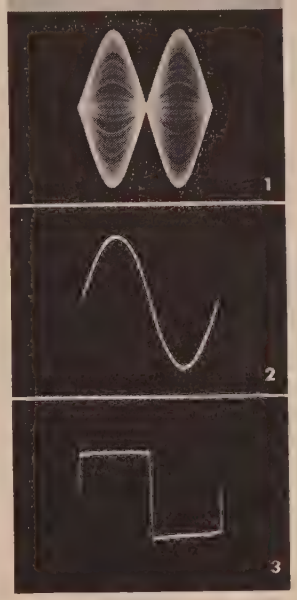


Contact stabilized with a 60 cps chopper, the output is linear with ±0.01 per cent of maximum output up to ±140 volts for loads as low as 10 k and up to ±100 volts for loads as low as 5 k. Output ripple is less than 0.50 millivolts rms and the normal 24-hour drift for a temperature rise of from 10 degrees to 60 degrees C is less than 0.5 millivolt after 15 minute warmup.

The normal open-loop gain exceeds 10^5; fixed gains may be obtained accurate to ±0.01 per cent. The Cd plated plug-in steel chassis may be mounted wherever convenient since there are no operating controls on the chassis.

(Continued on page 165A)

"Wanna Nyquist Diagrammed? A Function Transferred?"



Output wave forms of Servoscope displayed against internal linear sweep generator, frequency ½ cycle.

In one convenient instrument, here is test equipment for determining, in design or production phases, the dynamic performance of regulators, governors, process controls, positioning servomechanisms.

Only the

SERVOSCOPE®

has all these features:

- Applicable to both AC carrier and DC servo systems.
- Generates: **1. Sine-wave modulated carrier 2. Low frequency sine wave 3. Low frequency square wave**
- Built-in electronic sweep with no sweep potentiometer to wear out and require replacement.
- Dynamic frequency control range of 200 to 1.

For detailed specifications write Dept. IRE-5.

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New Hyde Park, New York

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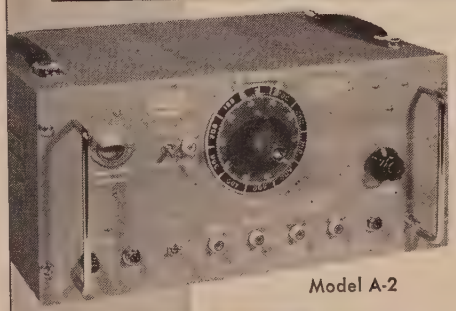
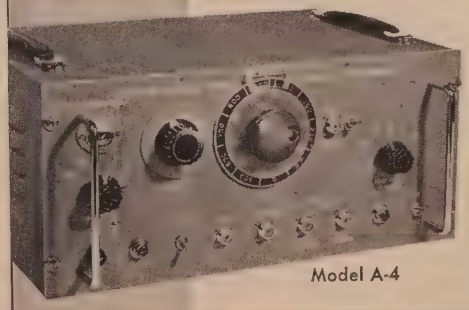
PRECISION TIMING INSTRUMENTS



Model A-2



Model A-4

Our TIME DELAY GENERATORS:

Each provides accurate and variable time intervals in five ranges. They feature low jitter (.008%), linear scales, built-in calibration indicator, 1,000-division dial, small repetition rate effects, blocking oscillator output and wide pulse output.

A-2 — Range: .8 μs to 100,000 μs
Get complete data: our Bulletin I-A-2

A-4 — Range: .00001 to 10 secs.
Get complete data: our Bulletin I-A-4

Rutherford ELECTRONICS CO.

Telephone: TExas 0-4362
3707 S. ROBERTSON BLVD.
CULVER CITY, CALIFORNIA

COMMUNICATIONS EQUIPMENT CO.

MICROWAVE COMPONENTS

"S Band," RG48/U Waveguide

POWER SPLITTER for use with type 726 or any 10 CM Shepherd Klystron. Energy is fed from Klystron antenna through dual pick-up system to 2 type "N" connectors..$22.50 EACH

DIRECTIONAL COUPLER. Broadband type "N" Coupling, 20 db. with std. flanges. Navy #CAEW47AAN-2 (as shown) $37.50

LHTR. LIGHTHOUSE ASSEMBLY. Part of RT39 APG 5 & APG 15. Receiver and Trans. Cavities w/assoc Tr. Cavity and Type N CPLG. To Recvr. Uses 2C40, 2C43, 1B27, Tunable APX 2400 2700 MCS. Silver Plated $49.50

BEACON LIGHTHOUSE cavity p/o UPN-2 Beacon 10 cm. Mfg. Bernard Rice, each $47.50

MAGNETRON TO WAVEGUIDE Coupler with 721-A Duplexer Cavity, gold plated $45.00

RT-30 APG-5 10 cm. lighthouse RF head c/o Xmtr.-Recvr-Tlt, cavity compl. recvr. & 30 MC IF strip using 60K5 (2C40, 2C43, 1B27 lineup) w/Tubes.

721A TR BOX complete with tube and tuning plungers $12.50

McNALLY KLYSTRON CAVITIES for 707B or 2K28 .....$4.00

WAVEGUIDE TO ⅞" RIGID COAX "DOOR-KNOB" ADAPTER CHOKE FLANGE, SILVER PLATED BROAD BAND $32.50

AS14A AP-10 CM Pick up Dipole with "N" Cables $4.50

OAJ ECHO BOX. 10 CM TUNABLE $22.50

HOLMDELL-TO-TYPE "N" Male Adapters, W.E. #D167284 .....$2.75

I.F. AMP. STRIP: 30 MC, 30 d.b. gain, 4 MC Bandwidth, uses 6AC7's—with video detector. A.F.C. less tubes $24.50

BEACON ANTENNA, AS31/APN-7 in Lucite Ball. Type "N" feed $22.50

ANTENNA, AT49A/APR: Broadband Conical, 300-3300 MC Type "N" Feed $12.50

"E" or "H" PLANE BENDS, 90 deg. less flanges $7.50

X Band— RG 52/U WAVEGUIDE

VSWR Measuring Section. Consisting of 6" straight section, with 2 pick-up, Type "N" Output Jacks. Mounted ½ Wave apart $8.50

1" x ½" waveguide in 5' lengths, UG 39 flanges to UG40 cover .....per length $7.50

Rotating-Joints supplied either with or without deck mounting. With UG40 flanges .each, $17.50

Bulkhead Feed-thru Assembly $15.00

Pressure Gauge Section 15 lb. gauge and press nipple $10.00

Pressure Gauge, 15 lbs. $2.50

Directional Coupler, UG-40/U Take off 20db $17.50

TR-ATR Duplexer section for above $8.50

Waveguide Section 12" long choke to cover 45 deg. twist & 2¼" radius, 90 deg. bend $4.50

Waveguide Section 3 ft. long silver plated with choke flange $5.75

Rotary joint choke to choke with deck mounting $17.50

90 degree elbows, "E" plane 2½" radius .$12.50

Microwave Receiver, 3 CM. Sensitivity: 10-13μ Watts. Complete with L.O. and AFC Mixer and Waveguide Input Circuits. 6 I.F. Stages give approximately 120 DV gain at a bandwidth of 1.7 MC. Video Bandwidth: 2 MC. Uses latest type AFC circuit. Complete with all tubes, including 723A/B Local Oscillator $175.00

ADAPTER, waveguide to type "N", UG 81/U, p/o TS 12, TS-13, Etc. $27.50

ADAPTER, UG-163/U round cover to special btl. Flange for TS-45, etc. .....$2.50 ea.

K Band—½" x ¼" W.G. 1.25 CM.

APS-34 Rotating joint $49.50

Right Angle Bend E or H Plane, specify combination of couplings desired $12.00

45° Bend E or H Plane, choke to cover $12.00

Mitered Elbow, cover to cover $4.00

TR-ATR-Section, Choke to cover $4.00

Flexible Section 1" choke to choke $5.00

"S" Curve Choke to cover $4.50

Adapter, round to square cover $5.00

Feedback to Parabola Horn with pressurized window $27.50

90° Twist $10.00

JAN WAVEGUIDE FLANGES

UG 39/U	$1.10	UG 51/U	$1.65
UG 40/U	$1.25	UG 52/U	$3.40
UG 40A/U	$1.65	UG 52A/U	$3.40

MAGNETRONS

Type	Freq. Range(MC)	Peak Power Out (KW)	Duty Ratio	Price
2J21A	9345-9405	50		$8.75
2J22	3267-3333	265		7.50
2J27	2965-2992	275	.002	19.95
2J31	2820-2860	285	.002	24.50
2J32	2780-2820	285	.002	28.50
2J38*	3249-3263	5		16.50
2J39*	3267-3333	8.7		24.50
2J48	9310-9320	50	.001	24.50
2J49	9000-9160	50	.001	59.50
2J56*	9215-9275	50	.001	132.50
2J61†	3000-3100	35	.002	34.50
2J62†	2914-3010	35	.002	34.50
700B	690-700	40	.002	22.50
700D	710-720	40	.002	39.75
706EY	3038-3069	200	.001	32.50
706CY	2976-3007	200	.001	32.50
725-A	9345-9405	50	.001	Write
730-A	9345-9405	50	.001	24.50
4J38	3550-3600	750	.001	169.45

*—Packaged with magnet.
†—Tunable over indicated range.

5J23	$49	QK-253	$249
4J34	$125	3J3J	$85

70 WATT C. W. MAGNETRONS*

QK-60 ....2800-3025 MC. QK-62 ....3150-3375 MC.
QK-61 ....2975-3200 MC. PRICE .....$85 Ea.
*Tunable, and packaged with magnet.

—IN STOCK—

AIA	ASH	SO-8
APA-9	BG	SG-1
APA-10	DAS†	SJ-1
APN-3	DBS†	TAJ
APN-7	APT-2	TBK
APN-9*	APT-4	TBL
APS-2	MKIV	SCR520*
APS-3	MKX	SCR521
APS-4	RC145	SCR518
APS-6	RC148	
ASD	SO-1	

* COMPONENTS. † LORAN EQUIPMENT

—TEST SETS—

TS-10	TS-12	TS-159
TS-36	TS-56	TS-268
TS-47	TS-34	TS-270

I. F. AMPLIFIER STRIPS

MODEL SO: 30 McGain figure is 120 db. Bandwidth: 2Mc. Uses 6 stages of 6AC7 Plus one video detector. Less Tubes $24.50

Model 15: 30 Mc center frequency, Bandwidth 2.5 Mc. gain figure: 65 db. Uses 5 stages of 6AC7's. Has D. C. Restorer and Video Detector. A.F.C. Strip included. Input impedance: 50 Ohms. Less tubes $27.50

Model APS-4: Miniature IF strip, using 6AK5's 60 Mc center Freq. Gain: 95db at Bandwidth of 2.7 Mc. Less tubes $45.00

THERMISTORS

D-164699 Bead Type DCR: 1525-2550 Ohms @ 75 Deg. F. Coefficient: 2% Per Deg. Fahr. Max. Current 25 MA AC/DC $2.50

D-167332 Bead Type, DCR is 1525-2550 Ohms. Rated 25 MA at .825-1.175 VDC $1.35

D-167613 Disk Type DCR: 355 Ohms @ 75 Deg. F.P.M. 2.5%, 1 Watt $1.35

D-166228 Disk Type 7120 Ohms @ 60°F. 4220 Ohms @ 80°F. 2590 Ohms @ 100°F.. 1640 Ohms @ 120°F. $1.35

PULSE NETWORKS

15A—1-400-50: 15 KV, "A" CKT. 1 microsec. 400 PPS, 50 ohms imp. $37.50

G.E. #3E (3-84-810) 8-2.24-405) 50P4T ; 3KV "E" CKT Dual Unit; Unit 1, 3 sections, 0.84 Microsec. 810 PPS, 50 ohms imp.; Unit 2, 8 Sections, 2.24 microsec. 405 PPS 50 ohms imp. $6.50

7-5E3-1-200-67P, 7.5 KV "E" Circuit, 1 microsec. 200 PPS, 67 ohms impedance 3 sections $7.50

7-5E4-16-60, 67P, 7.5 KV "E" Circuit, 4 sections 16 microsec, 60 PPS, 67 ohms impedance $15.00

7-5E3-3-200-67P, 7.5 KV, "E" Circuit, 3 microsec. 200 PPS, ohms imp. 3 sections $12.50

#755: 10KV. 2.2usec., 375 PPS, 50 ohms imp. ...$27.50

#754: 10KV. 0.85usec., 750 PPS, 50 ohms imp. .. $27.50

KS8865 CHARGING CHOKE: 115-150 H @ .02A, 32—40H @ .08A. 30,700V Corona Test. 21KV Test $37.50

G.E. 25E5-1-350-50 P2T. "E" SKT, 1 Microsec. Pulse @ 350 PPS. 50 OHMS Impedance $69.50

KS9623 CHARGING CHOKE: 16H @ 75 MA, 380 Ohms DCR. 9000 Vac test $14.95

G.E. 6E3-.5-2000 50 P2T: 6 KV., "E" Circuit 0.5 usec /2000 PPS/50 ohms/2 sections $7.50

PULSE EQUIPMENT

MIT. MOD. 3 HARD TUBE PULSER: Output Pulse Power 144 KW (12 KV at 12 Amp). Duty Ratio: .001 max. Pulse duration: 5, 1.0, 2.0 microsec. Input voltage: 115 v. 400 to 2400 cps. Uses: 1-71B, 4-89-R. 3-'72's, 1-'73. New .....Less Cover—$135

TPS-3 PULSE MODULATOR. Pk. power 50 amp. 24 KW (1200 KW pk): pulse rate 200 PPS. 1.5 microsec. pulse line impedance 50 ohms. Circuit series charging version of DC Resonance type. Uses two 705-A's as rectifiers, 115 v. 400 cycle input. New with all tubes $49.50

PULSE TRANSFORMERS

RAYTHEON WX 4298E: Primary 4KV., 1.0 USEC. SEC: 16KV-16 AMP DUTY RATIO: .001 400 CYCLE FIL. TRANS. "BUILT-IN" $42.50

WECO: KS 9948: Primary 700 ohms; Sec: 50 ohms. Plate Voltage: 18 KV. Part of APQ-13 .....$12.50

GE #K-2449A
Primary: 9.33 KV, 50 ohms Imp.
Secondary: 28 KV, 450 ohms.
Pulse length: 1.0/5 usec @ 635/120 PPS, Pk Power Out: 1,740 KW
Bifilar: 1.5 amps (as shown) ...$62.50

GE #K-2748-A, 0.5 usec @ 2000 Pps. Pk. Pwr. out is 32 KW impedance 40:100 ohm output. Pri. volts 2.3 KV Pk. Sec. volts 11.5 KV Pk. Bifilar rated at 1.3 Amp. Fitted with magnetron well $39.50

K-2745 Primary: 3.1/2.8 KV, 50 ohms Z. Secondary: 14/12.6 KV 1025 ohms Z. Pulse Length: 0.25/1.0 usec @ 600/600 PPS. Pk. Power 200/150 KW. Bifilar: 1.3 Amp. Has "built-in" magnetron well $42.50

K-2461-A. Primary: 3.1/2.6 KV—50 ohms (line). Secondary 14/11.5 KV—1000 ohms Z. Pulse Length: 1 usec @ 600 PPS. Pk. Power Out: 200/130 KW. Bifilar: 1.3 Amp. Fitted with magnetron well .$39.75

UTAH X-151T-1: Dual Transformer, 2 Wdgs. per section 1:1 Ratio per sec 13 MH inductance 30 ohms DCR $7.50

UTAH X-150T-1: Two sections, 3 Wdgs. per section. 1:1:1 Ratio, 3 MH, 6 ohms DCR per Wdg. $7.50

68G711: Ratio: 4:1 Pri: 200V. Sec. 53V, 1.0 usec Pulse @ 2000 PPS, 0.016 KVA $4.50

TR1049 Ratio: 2:1 Pri, 220 MH, 50 Ohms, sec. 0.75 H. DCR 100 Ohms $6.75

K-904695-501: Ratio 1:1, Pri. Imp. 40 Ohm, Sec. Imp. 40 Ohms. Passes pulse 0.6 usec with 0.05 usec rise $8.95

Ray UX 7896—Pulse Output Pri, 5v sec. 41v $7.50

Ray UX 8442—Pulse inversion—40v + 40v $7.50

PHILCO 352-7250. 352-7251. 352-7287

RAYTHEON: UX8693, UX5986 .....$5 ea.

W.E.: D-166310, D-166638, KS9800, KS9948.

UTAH #9262, with Cracked Beads, but will operate at full rated capacity $5.00

UX 8693 (SCS #2Z9627-54): 3 Wdgs, 32 turns #18 wire. DCR is: 362/.372/.4 ohms. Total voltage 2500 vdc. $5.00

D-166173: Input: 50 ohms Z. Output: 900 ohms 3 Wdgs. Freq range 10 kc-2mc. P/O AN/APQ*-13 $12.50

K-2450: Pulse-inversion auto-transformer: primary 13 kv, 4 usec. Output: 14 kv @ 100 kw peak $34.50

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AT49/APR—Broadband Conical. 300-3300 MC. Type N Feed $8.95

AS-31/APN-7: 10 cm Polyrod in Lucite Ball. Type N Fitting Coax Feed $22.50

Relay System Parabolic reflectors approx. range 2000 to 6000 Mc. Dimensions 4½" x 3'. New $100.00

Dipole for above $12.00

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News—New Products

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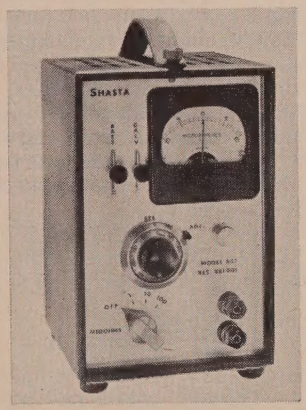
(Continued from page 159A)

High Resistance Bridge

Covering a range from 1,000,000 ohms to 100 megohms, the Model 601 Bridge developed by **Shasta, div. of Beckman Instruments, Inc.,** P.O. Box 296, Station A, Richmond, Calif., measures resistance to an accuracy of better than 0.25 per cent. The circuit employed is a simple Wheatstone Bridge network with the unknown and a 10-turn Helipot potentiometer in adjacent legs. Bridge unbalance voltage is amplified by a vacuum-tube voltmeter and fed to a null indicating panel microammeter. The circuit is such that the meter cannot be damaged by severe unbalance conditions. At balance, resistance values correspond to Helipot dial settings multiplied by appropriate factor of 10.



This instrument can be used in the laboratory for production testing of resistors, and for general-purpose resistance measurements. It is basically ac mains operated, but contains a 45 volts B battery for the bridge voltages for greater stability.

The Model 601 Bridge is housed in a miniature Shasta "A" cabinet.

Frankel New PA For CBS Columbia

M. A. Gardner, Vice President in charge of Purchases, of CBS-Columbia, the television receiver manufacturing division of Columbia Broadcasting System, announced today that Albert J. Frankel has been promoted to the position of Purchasing Agent.

(Continued on page 166A)

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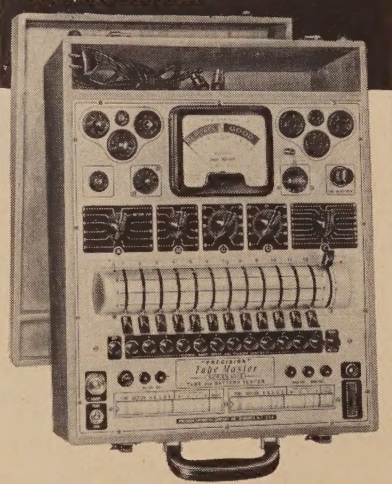

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News—New Products

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(Continued from page 165A)

Electron Tube Reliability

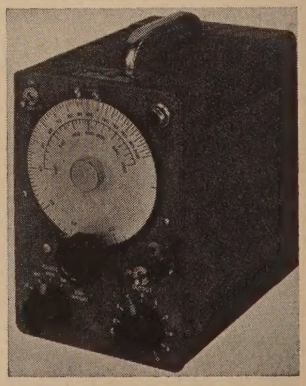
A 97-page report entitled "*Investigation of Electron Tube Reliability in Military Applications*" has just been published by **Aeronautical Radio, Inc.,** (ARINC), a non-profit organization currently conducting an extensive study of military electronic tube reliability for the Army, Navy, and Air Force under a Bureau of Ships contract.

The report covers the period from April 4, 1951 through March 31, 1953, and describes the scope of the ARINC surveillance program, the methods of tube collection and of engineering and statistical analysis, and the result of the program to date. It contains a discussion of the characteristic defect patterns of tube returns from each of the eight military bases under surveillance and an analysis of each of the 20 individual tube types ranking highest in number of tubes returned. There is also an evaluation of tube weaknesses, unreliability, such as environment, operating procedures, misapplications, and maintenance. Incorporated in the report are 65 tables, 24 charts, and photographs, and a statistical appendix.

The report is available for 50¢ from L. E. Davis, Aeronautical Radio, Inc., 1523 "L" St., N.W., Washington 5, D. C.

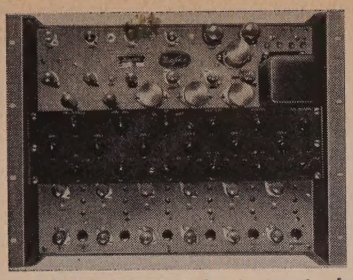
Audio Oscillator

The model 411 oscillator, manufactured by **The Clough-Brengle Co.,** 6041 Broadway, Dept. P, Chicago 40, Ill., is suitable for making measurements requiring a sine wave signal over the range from 20 cps to 1 mc.



A resistance-capacity tuned-type oscillator and a cathode follower in the output system are employed to provide uniform response. Frequency accuracy is maintained by the use of deposited carbon resistors in the frequency determining network. Low-level measurements are facilitated by a panel switch which reduces output voltage, distortion and hum output. Other features of the Model 411 are good case ventilation, a well spread dial calibration for ease in reading, and compact size with light weight. For additional information write to The Clough-Brengle Co.

(Continued on page 167A)



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News—New Products

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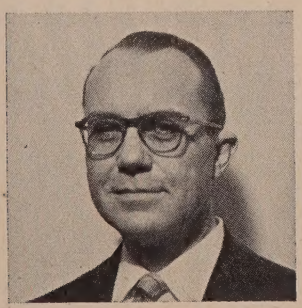
(Continued from page 166A)

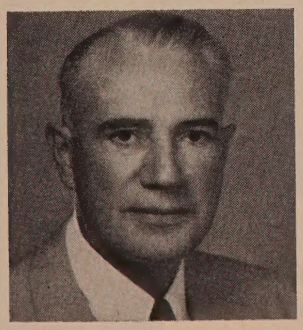
Miller New Field Engineer for General Plate

O. V. (Jack) Miller has been named field engineer for the Wisconsin-Minnesota area by the **General Plate Div. of Metals & Controls Corp.**, Attleboro, Mass.



His new duties will include application engineering of General Plate products: electrical contacts and contact materials, Truflex thermostat metals and fabricated parts, solid and composite precious metals, composite base metals, and other related products.

Tollaksen Moves to Remler



The appointment of Leslie B. Tollaksen as manager of technical products by **Remler Company, Ltd.**, San Francisco electronics manufacturer, has been announced by R. C. Gray, president. Tollaksen's headquarters will be the Remler office in the Wyatt Bldg., Washington, D. C.

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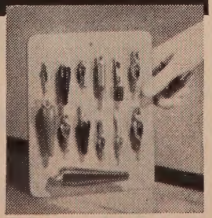

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